

Slip, Swim, Mix, Pack: Fluid Mechanics at the Micron Scale

Eric Lauga

*Division of Engineering and Applied Sciences
Harvard University*

Thesis advisors

Michael P. Brenner

Howard A. Stone

Current address: Department of Mathematics, MIT

Outline

SLIP

The no-slip boundary conditions in hydrodynamics

SWIM

Locomotion of swimming microorganisms near surfaces

MIX

3D flows in microchannels

PACK

Capillary-driven assembly of microparticles

Outline

SLIP

The no-slip boundary conditions in hydrodynamics

SWIM

Locomotion of swimming microorganisms near surfaces

MIX

3D flows in microchannels

PACK

Capillary-driven assembly of microparticles

Flow boundary conditions

What is the appropriate boundary condition for **Newtonian liquid** flow past a **solid** surface?

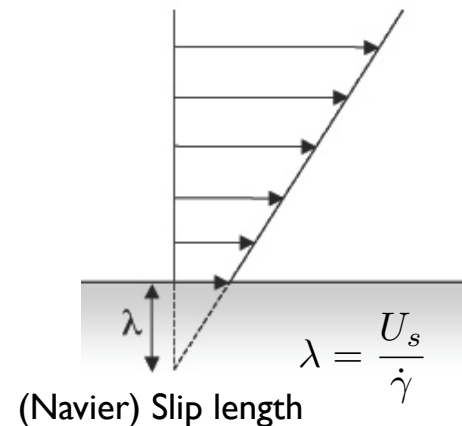
$$\rho \frac{D\mathbf{u}}{Dt} = -\nabla p + \mu \nabla^2 \mathbf{u}, \quad \nabla \cdot \mathbf{u} = 0$$

$$\mathbf{u} \cdot \mathbf{n} = 0 \qquad \mathbf{u} \cdot \mathbf{t} = ?$$

D. Bernoulli, Euler, Coulomb, Girard,
Navier, Poisson, Poiseuille, Stokes, Hagen,
Darcy, Helmholtz, Maxwell, Couette...

Today, the no-slip boundary condition is in all textbooks.

Recent series of experiments: **apparent breakdown** of no-slip.

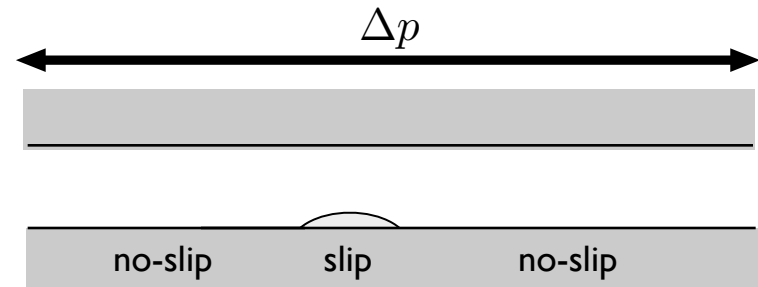


Apparent slip

Pressure drop vs. flow rate

Simple models for heterogeneous boundary conditions: distributed regions of (perfect) slip such as bubbles. Two parameters: length scale, surface coverage.

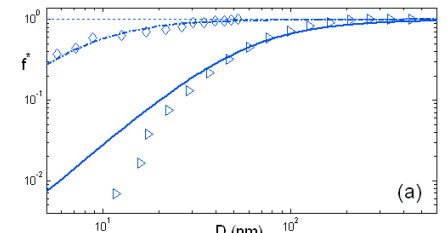
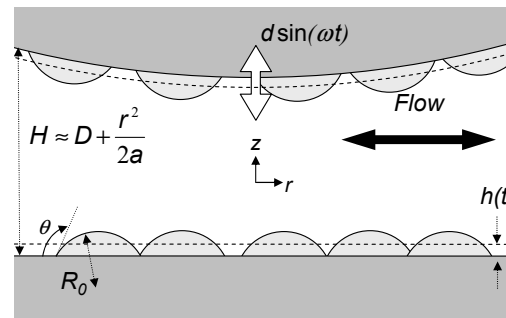
Lauga & Stone (2003) *J. Fluid Mech.* **489**, 55



Drainage (squeeze-flow)

Surface-attached bubbles in lead to shear-dependent apparent slip due to bubble diffusion and compression.

Lauga & Brenner (2004) *Phys. Rev. E.* **93**, 026311

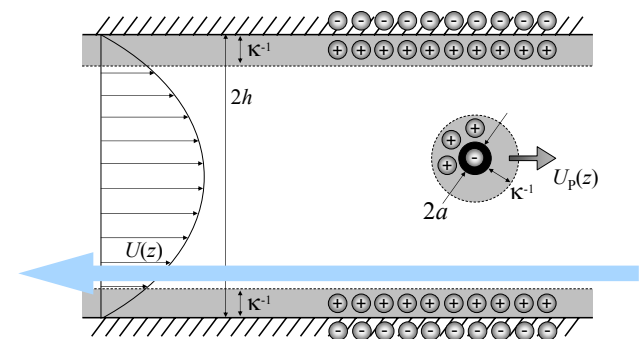


Zhu & Granick (2002) *Phys. Rev. Lett.* **88**, 106102

PIV

Flow of an electrolyte: If the tracer particles are charged, their velocities can include a fake (electrokinetic) slip component.

Lauga (2004) *Langmuir* **20**, 8924

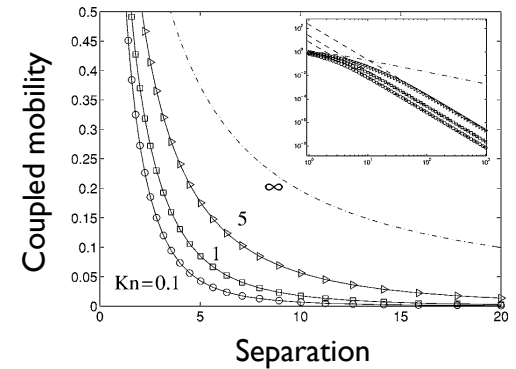
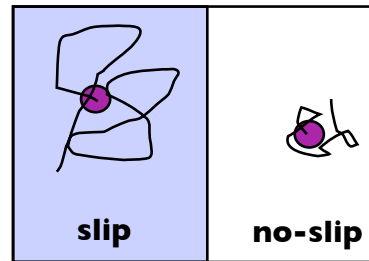


Other studies

A new method to measure slip

Diffusion of a colloidal probe near a slip surface can be used to infer slip lengths.

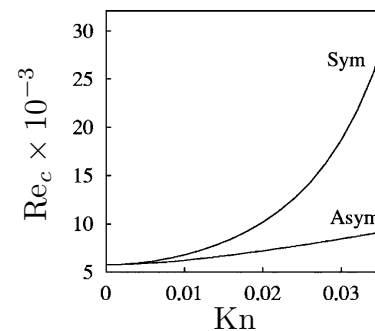
Lauga & Squires (2005) *Phys. Fluids* **17**, 103102



Influence of slip on flow stability

Slip has large stabilizing effect on normal modes but negligible effect on transient growth.

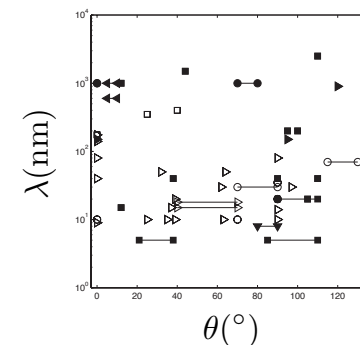
Lauga & Cossu (2005) *Phys. Fluids* **17**, 088106



Review

Discussion of theory, simulations, experiments.

Lauga, Brenner & Stone (2006) *Handbook of Experimental Fluid Dynamics* - In press



Outline

SLIP

The no-slip boundary conditions in hydrodynamics

SWIM

Locomotion of swimming microorganisms near surfaces

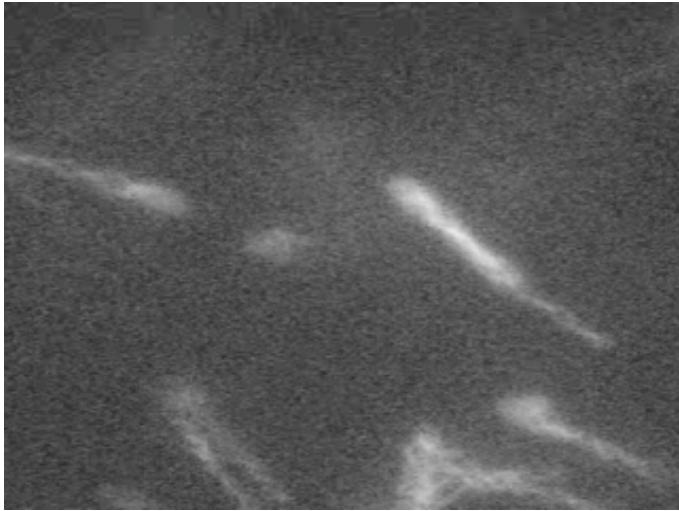
MIX

3D flows in microchannels

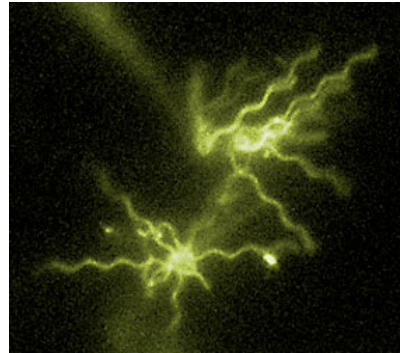
PACK

Capillary-driven assembly of microparticles

Swimming bacteria - *Escherichia coli*



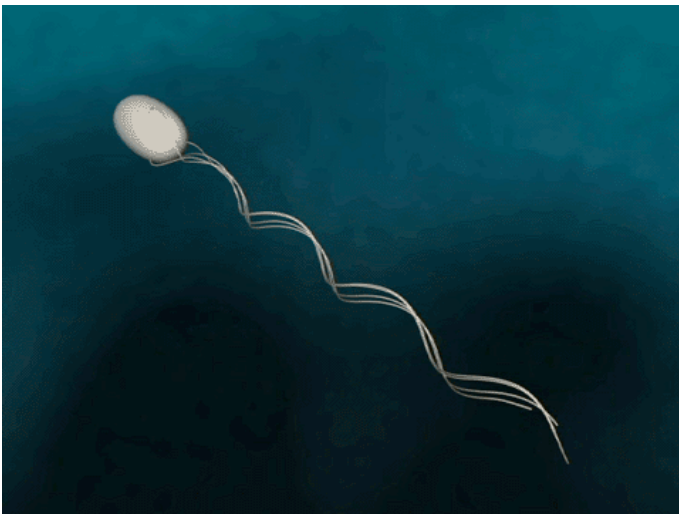
Berg Lab - Rowland Institute at Harvard



Berg Lab



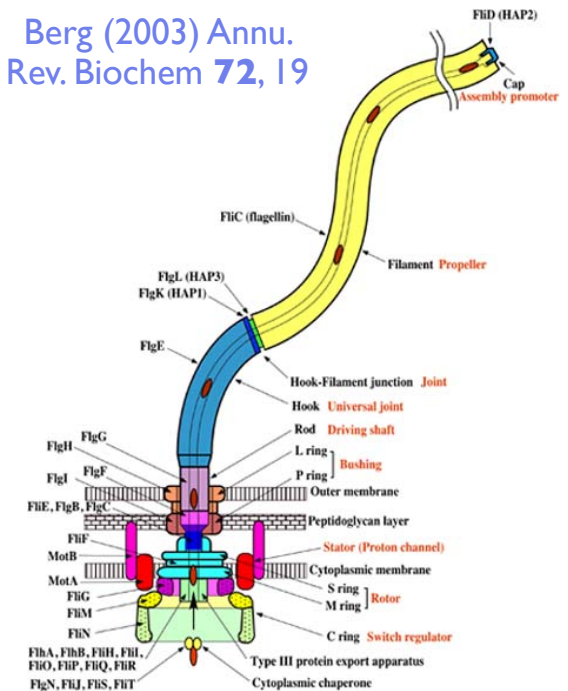
Berg (2003) Annu. Rev. Biochem **72**, 19



Protonic Nanomachine Project
<http://www.npn.jst.go.jp/>

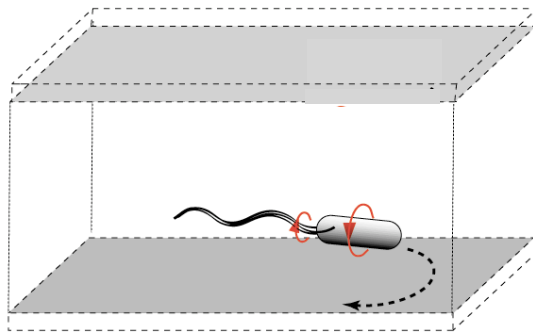


D. Kunkel Microscopy Inc.

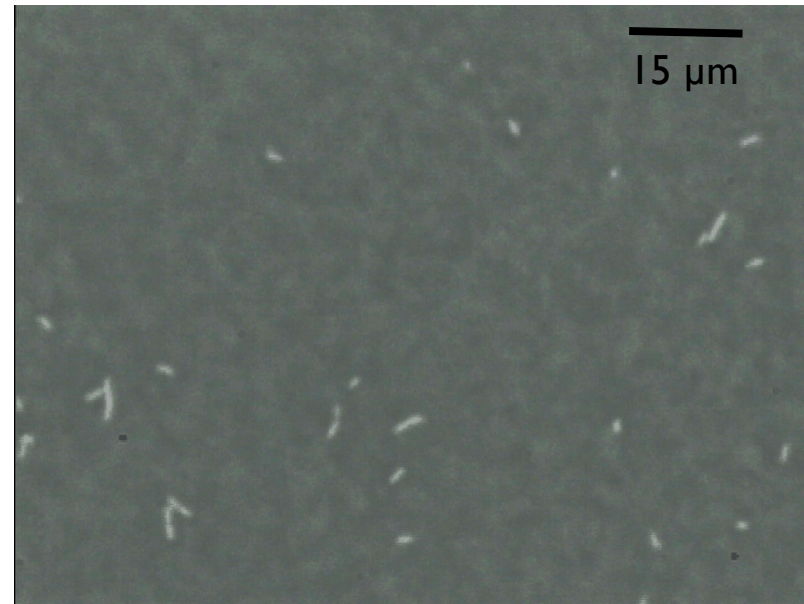


Protonic Nanomachine Project
<http://www.npn.jst.go.jp/>

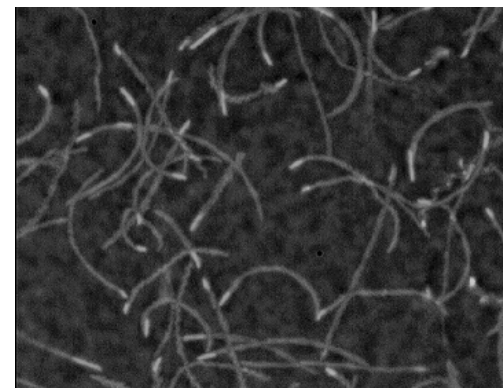
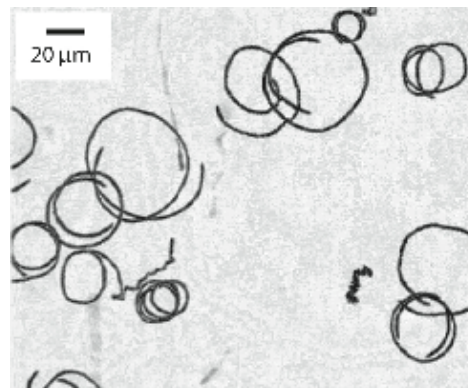
E. coli swimming near a surface



- Smooth-swimming *E. coli* in growth media
- Real time
- Velocity $\sim 30 \mu\text{m/s}$
- Microchannel height = $105 \mu\text{m}$
- Solid surface: PDMS

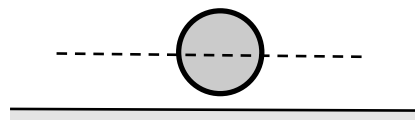
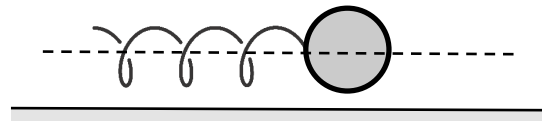


Experiment by Willow DiLuzio, DEAS, Harvard University

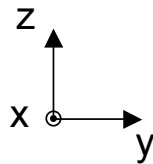


Superimposed images showing clockwise, circular paths

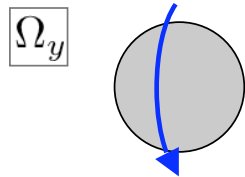
Why does the bacterium rotate?



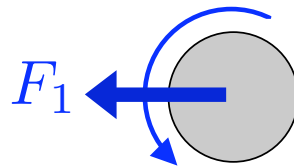
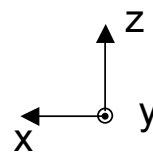
SIDE



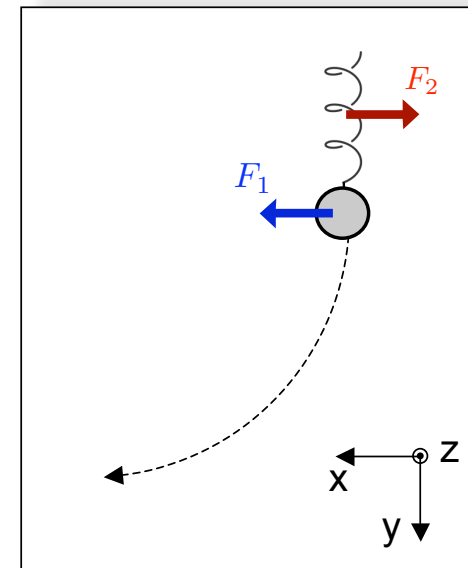
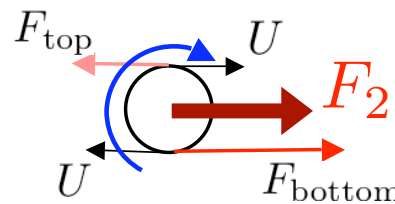
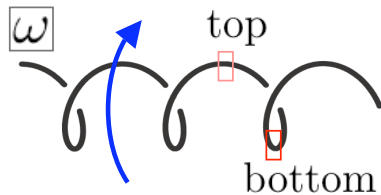
Head



FRONT



Flagella bundle

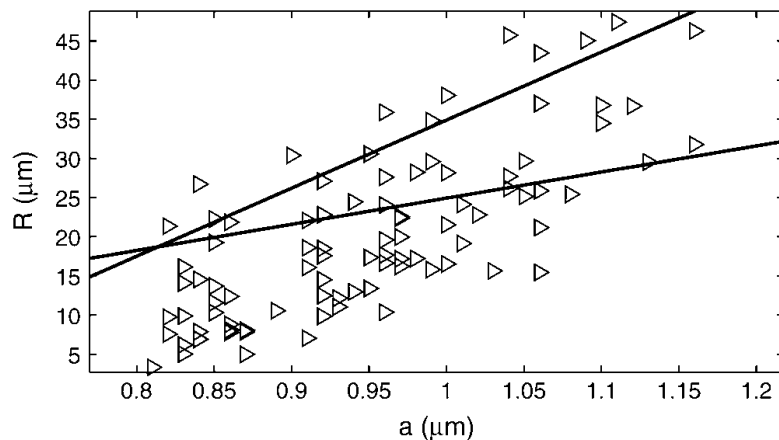
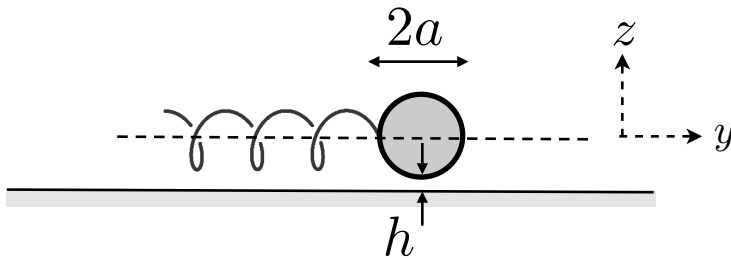


Hydrodynamic interactions between the free-swimming bacterium and the surface lead to an out-of-plane **torque**, and since the cell has to be **torque-free**, it will **rotate**.

Computing the circle radius

Axial propulsive force balances translational drag

Wall-induced torque balances rotational drag



$$\mathcal{M}_{yy}^{\text{F}\Omega}(\omega - \Omega_y) \approx (\mathcal{M}_{yy}^{\text{FU}} + \mathcal{N}_{yy}^{\text{FU}})U_y$$

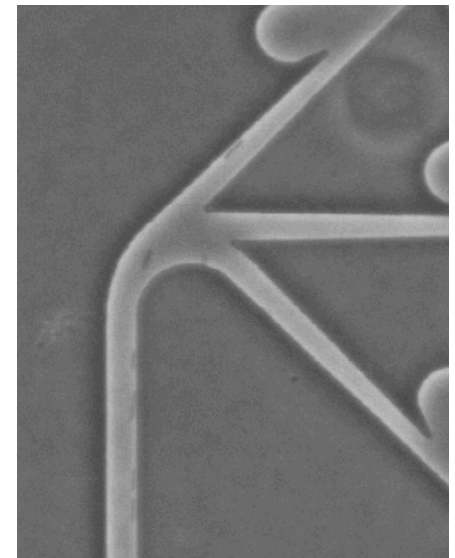
$$\mathcal{W}_{zy}^{\text{L}\Omega}(\omega - \Omega_y) \approx -\mathcal{M}_{zz}^{\text{L}\Omega}\Omega_z$$

$$\mathcal{R} \approx \frac{U_y}{|\Omega_z|} \approx \frac{\mathcal{M}_{zz}^{\text{L}\Omega} \mathcal{M}_{yy}^{\text{F}\Omega}}{\mathcal{W}_{zy}^{\text{L}\Omega} (\mathcal{M}_{yy}^{\text{FU}} + \mathcal{N}_{yy}^{\text{FU}})}$$

Radius of the circle

Exploiting the results

$$\frac{d\mathcal{R}}{da} > 0$$



Outline

SLIP

The no-slip boundary conditions in hydrodynamics

SWIM

Locomotion of swimming microorganisms near surfaces

MIX

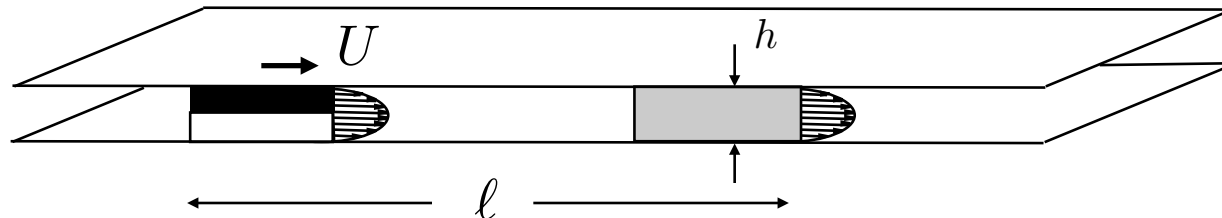
3D flows in microchannels

PACK

Capillary-driven assembly of microparticles

Mixing on (not too) small scales

Typically, **low Reynolds number** but **high Peclet number**.
How long downstream will molecular diffusion mix?

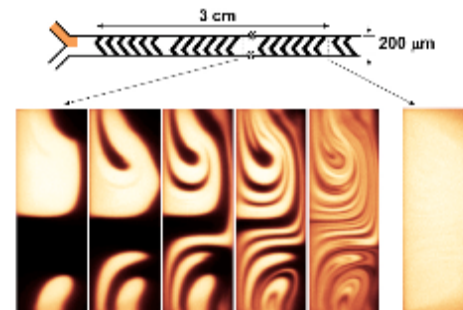
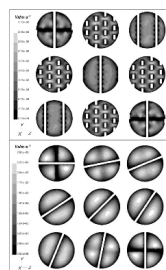
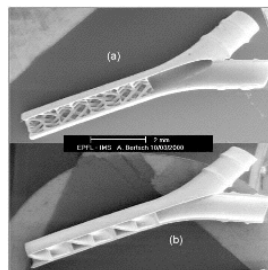


Diffusion time across the channel $\tau \sim h^2/D$

Distance traveled $l \sim U\tau$

$$\frac{l}{h} \sim \frac{Uh}{D} = \text{Pe} \gg 1$$

Solution: generate transverse flows to replace diffusive transport by convective transport

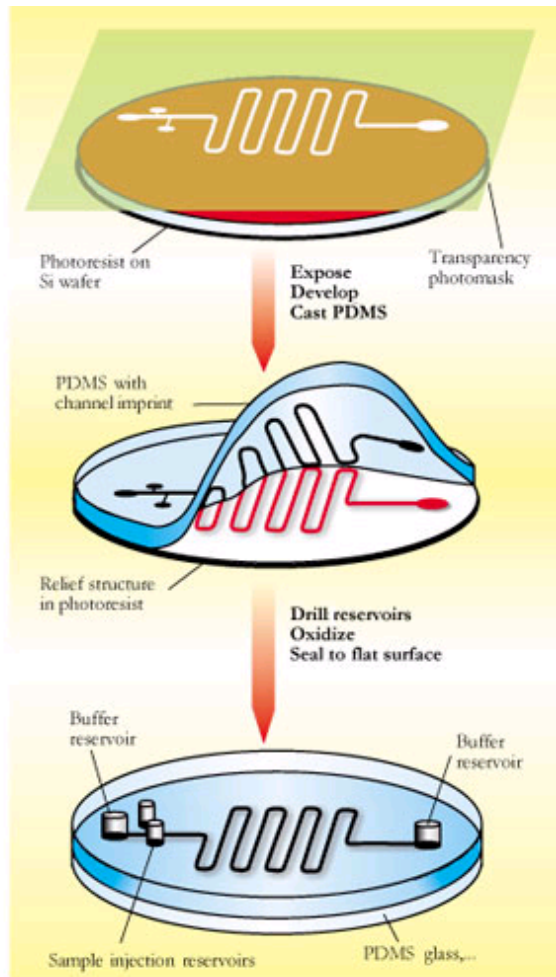


Bertsch et al. (2001) Lab Chip **1**, 56

Stroock et al. (2000) Science **295**, 5555

Copyright 2006 Eric Lauga

Fabrication constraints

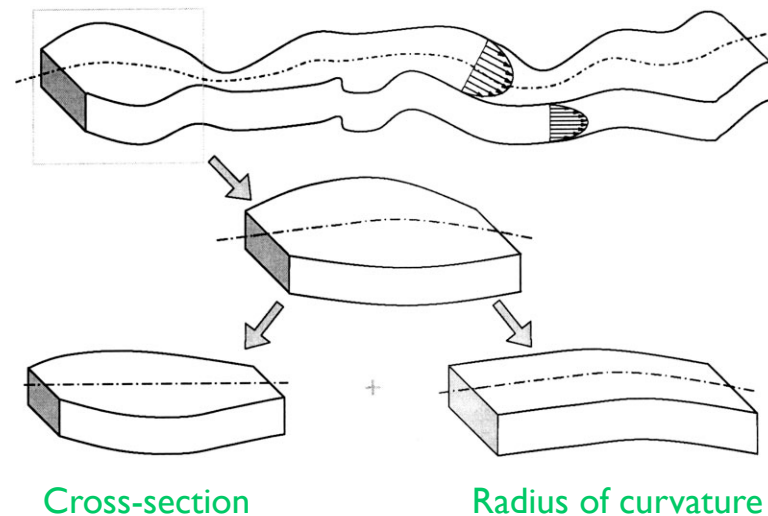


Soft Lithography
Simplest: **one step** of microfabrication



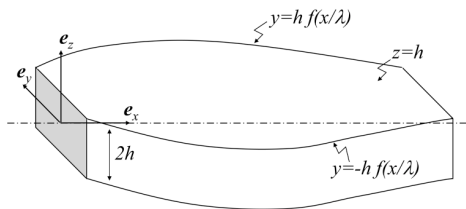
The channels will have a **fixed height**
Will the resulting channel be a good mixer?

Design: two degrees of freedom



Whitesides & Stroock (2001) *Physics Today* **54**, 42

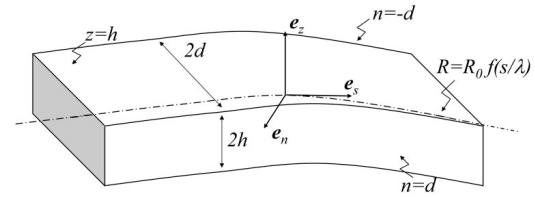
The flow is always 3D



Straight channel of
varying width



Flow is 3D unless constant width



Curved channel of constant width
and varying curvature



Flow is 3D unless constant curvature

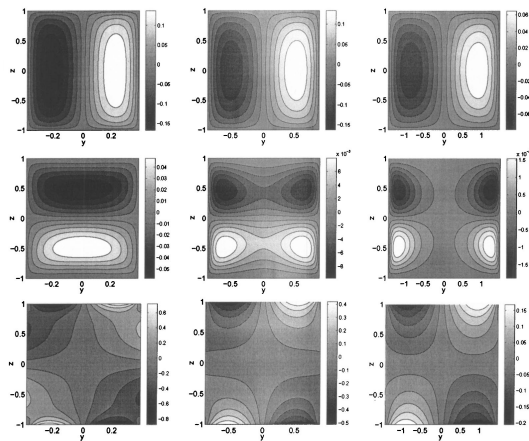
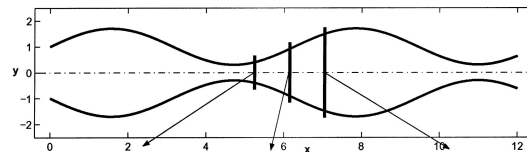


Illustration on periodic channel

Lauga, Stroock & Stone (2004) Phys. Fluids **16**, 3051

Outline

SLIP

The no-slip boundary conditions in hydrodynamics

SWIM

Locomotion of swimming microorganisms near surfaces

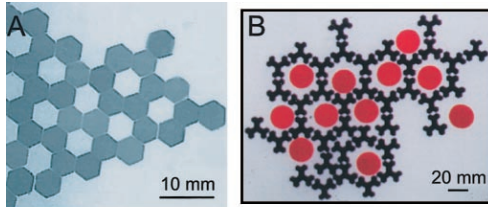
MIX

3D flows in microchannels

PACK

Capillary-driven assembly of microparticles

Self-assembly on small scales



Whitesides et al. (2002) PNAS **99**, 4769

Kralchevsky and Denkov (2001) Curr. Opin. Colloid Int. Sci. **6**, 383

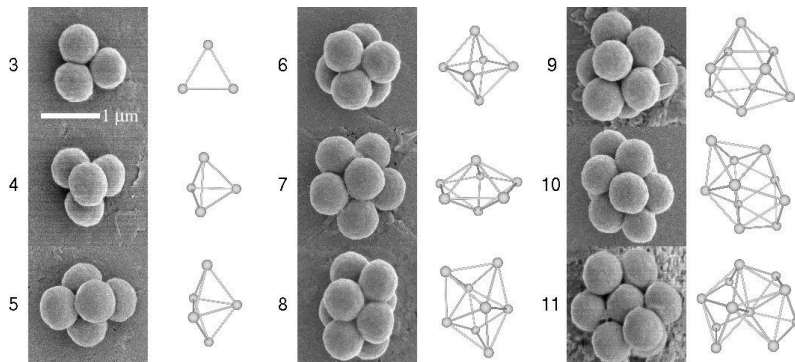
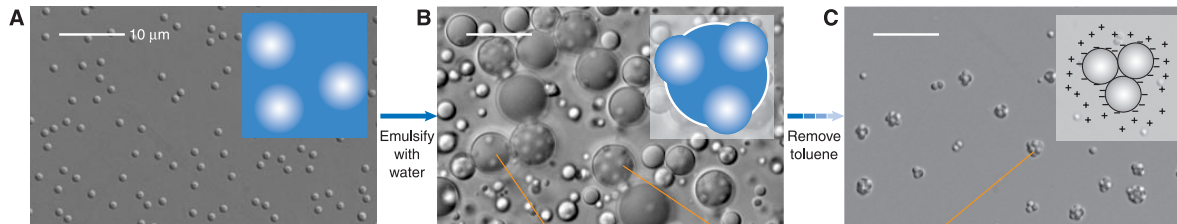
Reducing length scales:
Design an energy landscape

Usually many **local minima**

Ex: Packing of N spheres using vdW forces

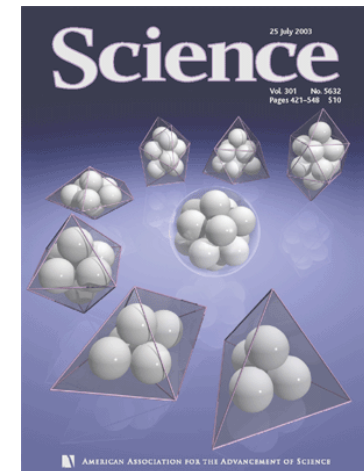
Number of Spheres	6	7	8	9	10	11	12	13
Number of Local minima	2	4	8	18	57	145	366	988

Hoare & McInnes (1976) Faraday Discuss. Chem. Soc. **61**, 12

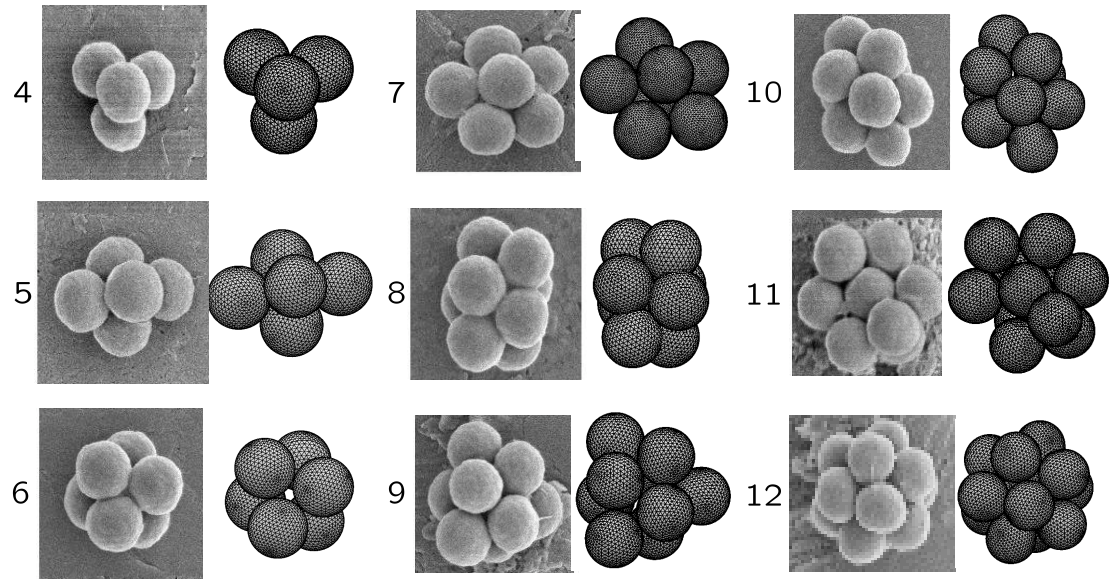
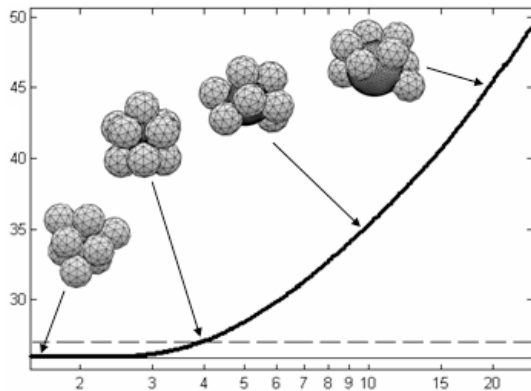
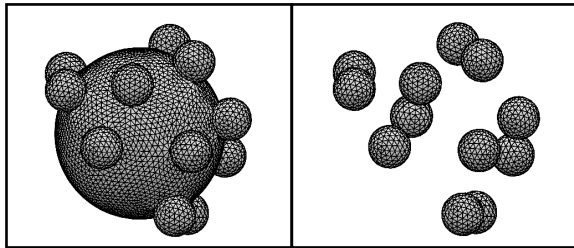


Final packings are unique
(no local minima)

Manoharan et al. (2004)
Science **301**, 483



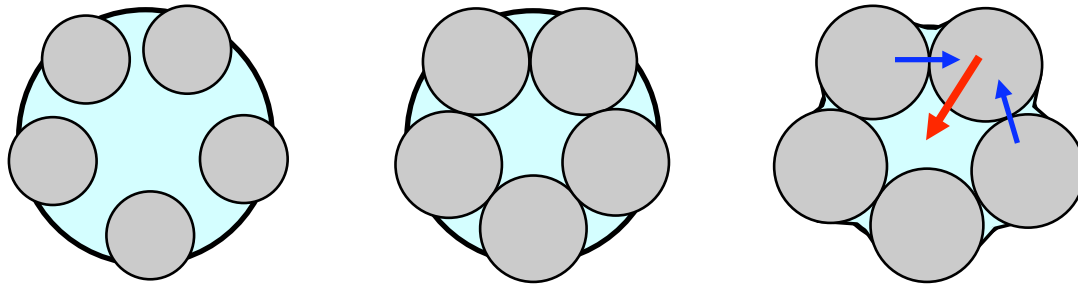
Simulations



Numerical simulations of hard spheres on droplet reproduce the packings obtained by Manoharan et al. [Lauga & Brenner \(2004\) Phys. Rev. Lett. 93, 238301](#)

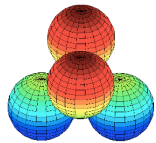
Experimental pictures from [Manoharan et al. \(2004\) Science 301, 483](#)
Copyright 2006 Eric Lauga

Theory

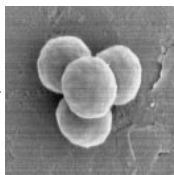


Rearrangement of particles is uniquely determined by geometry.

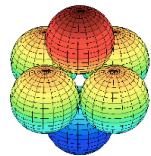
Tetrahedron



$N = 4$



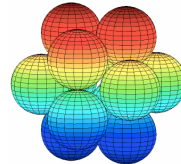
Octahedron



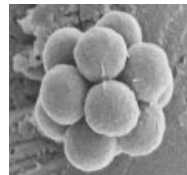
$N = 6$



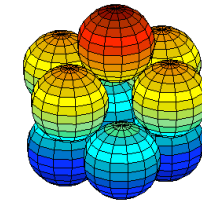
Icosahedron



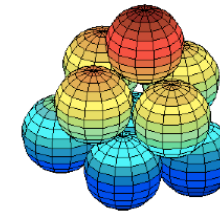
$N = 12$



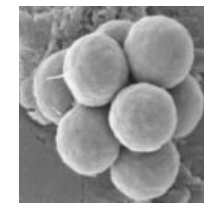
Some final packings correspond to surface jamming



Simple model
 $N = 9$

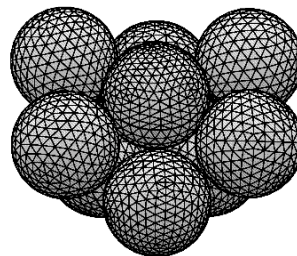
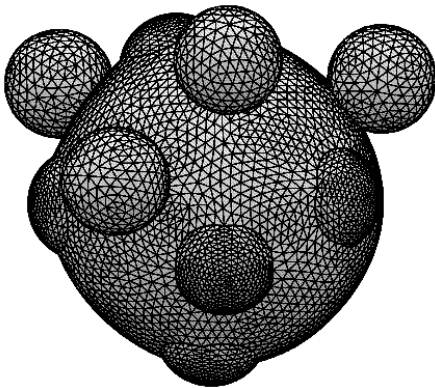
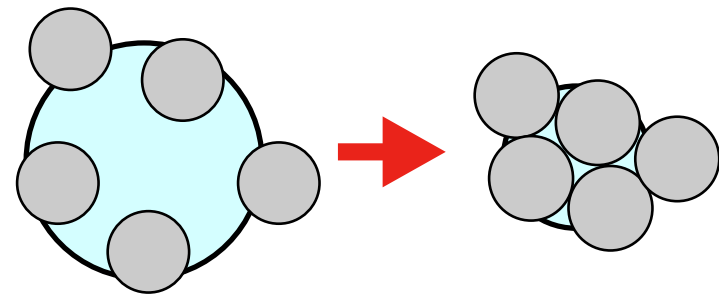
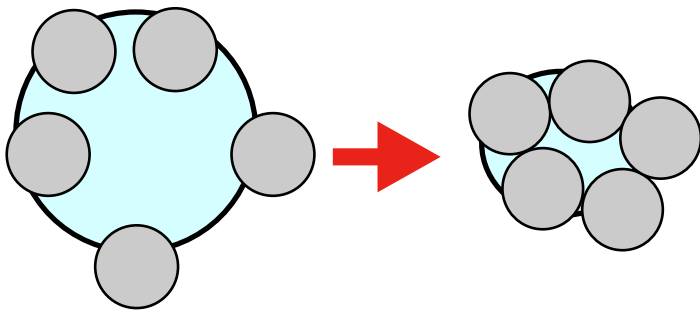


Full simulations



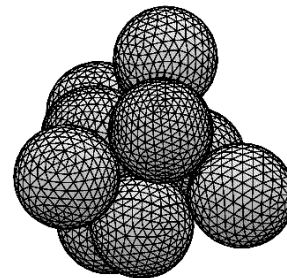
Experiments

How to create different packings



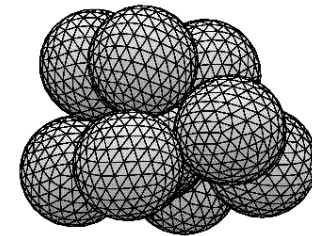
(a)

4 spheres with 160°
5 spheres with 20°



(b)

6 spheres with 160°
3 spheres with 20°

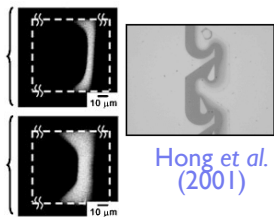


(c)

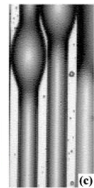
8 spheres with 160°
1 spheres with 20°

Lauga & Brenner (2004) *Phys. Rev. Lett.* **93**, 238301
Schnall-Levin, Lauga & Brenner (2006) *Langmuir* **22**, 4547

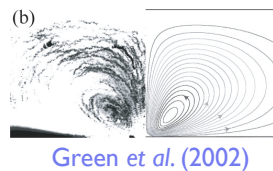
Fluid Mechanics at the micron scale...



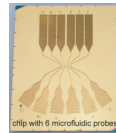
Hong et al. (2001)



Gau et al. (1999)

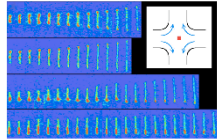


Green et al. (2002)

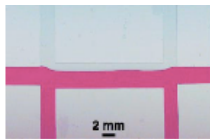


Delamarche et al. (2005)

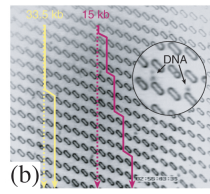
Ismagilov et al. (2000)



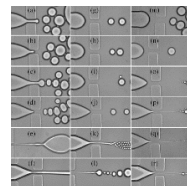
Perkins et al. (1997)



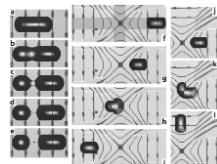
Zhao et al. (2001)



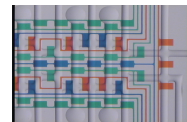
Chou et al. (2000)



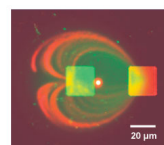
Anna et al. (2003)



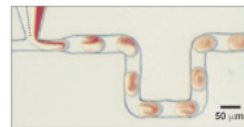
Darhuber et al. (2003)



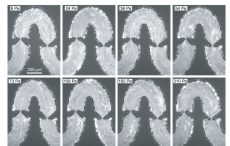
S. Maerkl



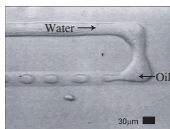
Juncker et al. (2005)



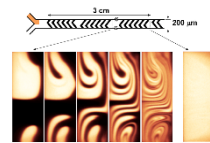
Song et al. (2003)



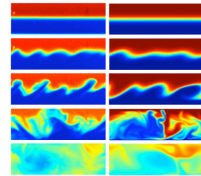
Groisman et al. (2003)



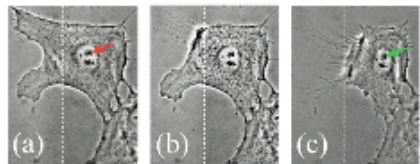
Thorsen et al. (2001)



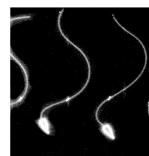
Stroock et al. (2002)



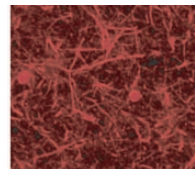
Lin et al. (2004)



Takayama et al. (2003)



C. Brokaw



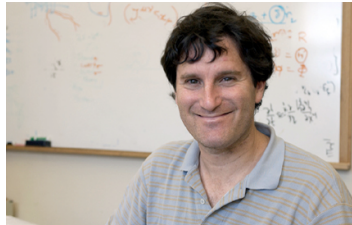
Shin et al. (2004)

... is not just the study of viscous flows:
there is a lot of additional physics

- Surface effects and wetting
- Interface deformations and two-phase flow
- Advective transport vs. diffusion
- Slip and non-continuum effects
- Elastic and non-Newtonian forces
- Electrokinetics and electrical forces
- Intermolecular forces
- Heat transfer
- Porous material
- Acoustic streaming and sound waves
- Suspensions
- Living cells
- Chemical reactions

Beebe et al. (2002) *Annu. Rev. Biomed. Eng.* **4**, 261
 Stone, Stroock & Adjadri (2004) *Ann. Rev. Fluid Mech.* **36**, 381
 Darhuber & Troian (2005) *Ann. Rev. Fluid Mech.* **37**, 425
 Delamarche et al. (2005) *Adv. Mat.* **17**, 2911
 Squires & Quake (2005) *Rev. Mod. Phys.* **77**, 977

Thank you



Michael P. Brenner



Howard A. Stone



Carlo
Cossu



Willow
DiLuzio



Todd
Squires



Abraham
Stroock

Manouk Abkarian, Silas Alben, Shelley Anna, Jacquie Ashmore, Raymond Bergmann, Thomas Bewley, Nathalie Bontoux, Henry Chen, Marc Durand, Joel Frenzer, Jose Gordillo, Patrick Huerre, John Hutchinson, Stephan Koehler, Srinivas Paruchuri, Marcus Roper, Jim Rice, Michael Schnall-Levin, Steven Subotnick, Thomas Ward, George Whitesides, Kate Zirpolo, and many others...

Funding: NSF, Harvard MRSEC, ONR, Ecole des Mines de Paris

American Physical Society