

MORPHODYNAMICS OF RIVERS AND TURBIDITY CURRENTS:

AN ELEGANT CONVERSATION BETWEEN WATER AND SEDIMENT



National Center for Earth-surface Dynamics



Image courtesy A. Alabyan



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University of Illinois



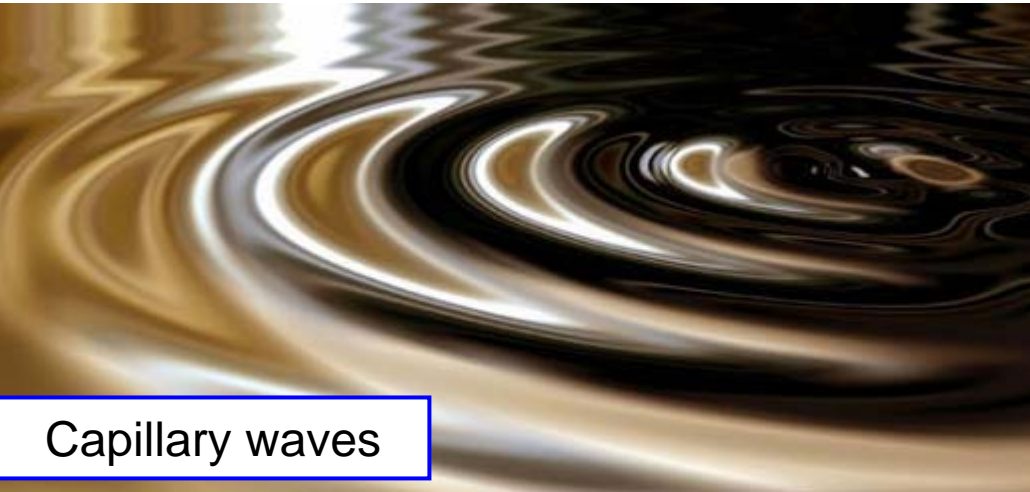
A CIVIL ENGINEER/GEOLOGIST GIVING AN INVITED TALK AT THE AMERICAN PHYSICAL SOCIETY

IS LIKE A COUNTRY PRIEST GRANTED AN AUDIENCE WITH
THE POPE

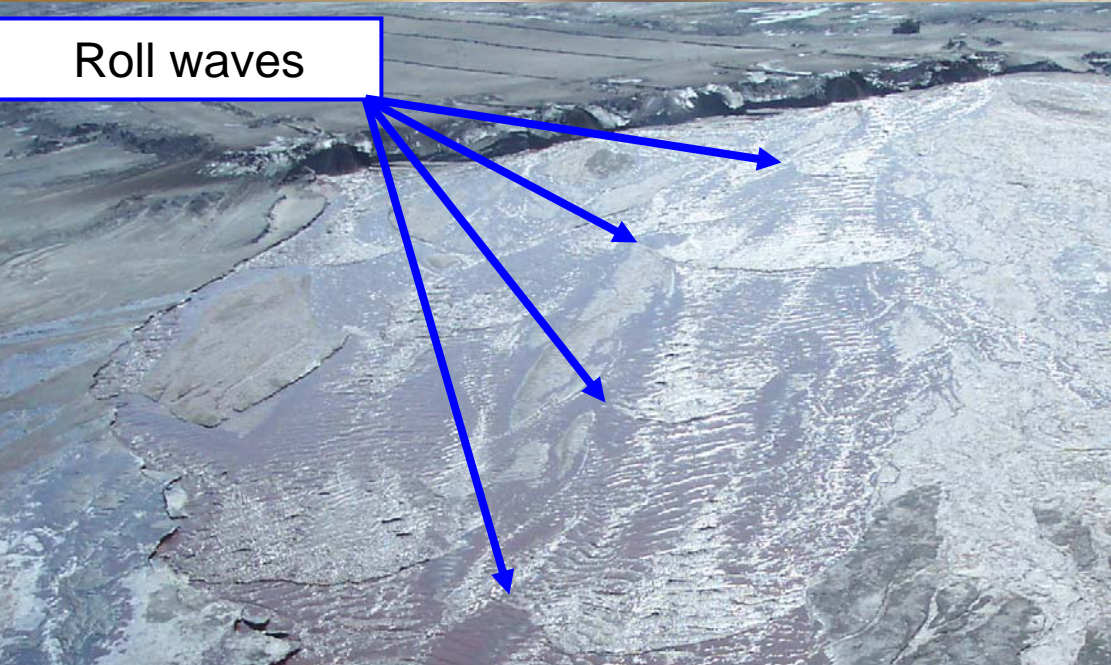


"PURE FLUID MECHANICS"

Haboob dust storm



Capillary waves



Roll waves

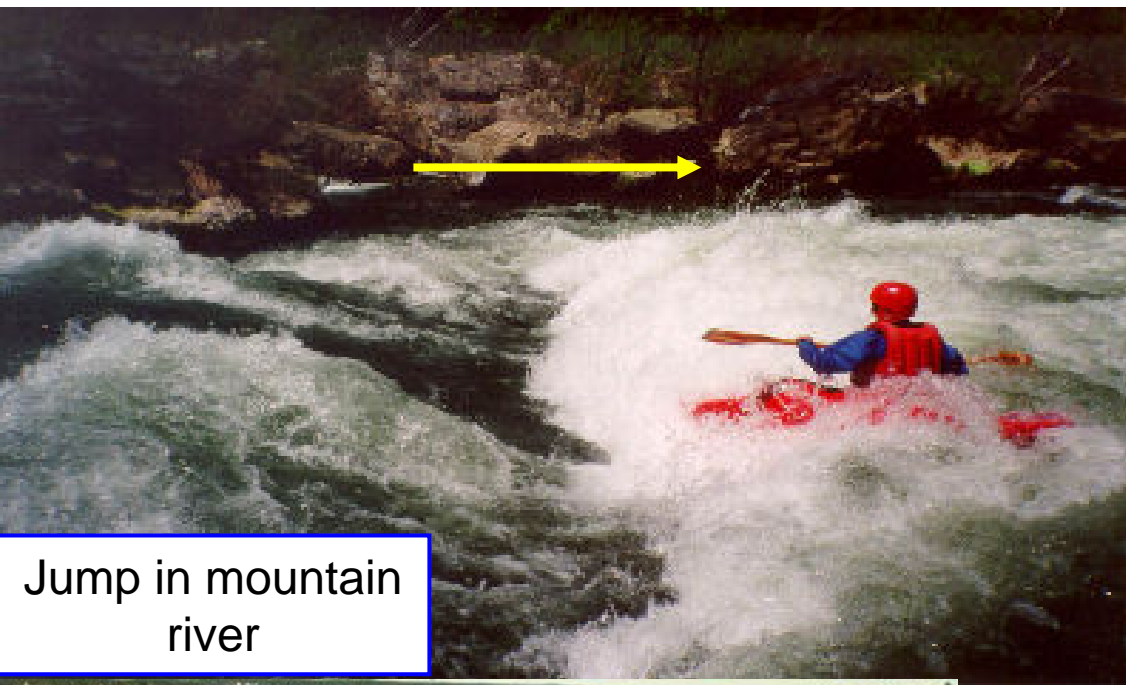


Wind ripples

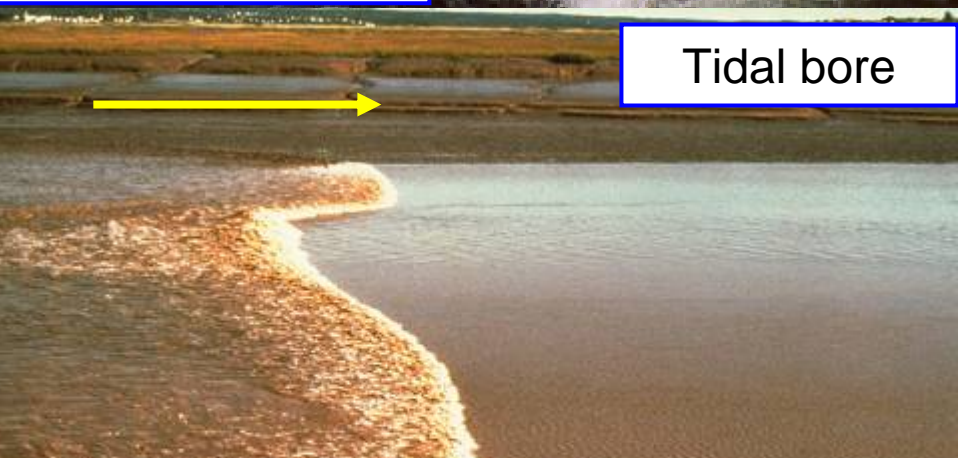
Image courtesy IOCC

<http://www.rikenresearch.riken.jp/research/223/images/2234070426115631.jpg>
http://scribalterror.blogspot.com/scribal_terror/images/2007/05/02/dust_2.jpg
<http://images.jupiterimages.com/common/detail/66/41/23354166.jpg>

HYDRAULIC JUMPS AND BORES



Jump in mountain river

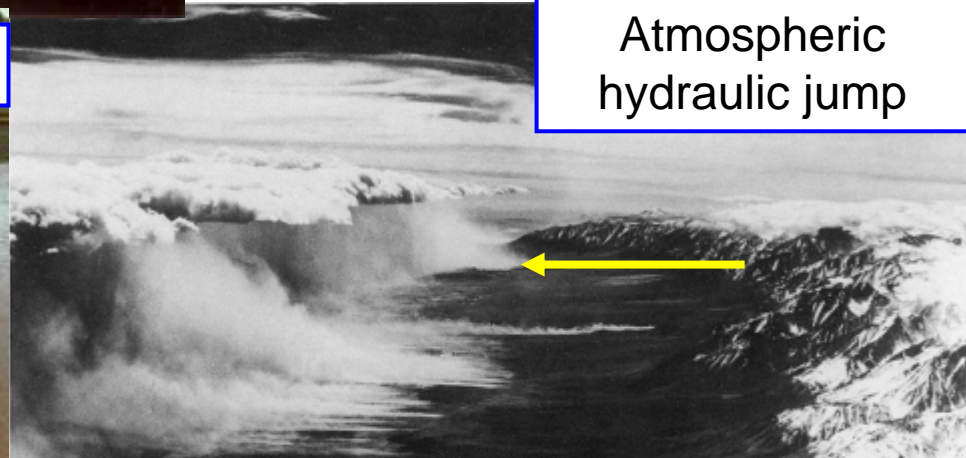


Tidal bore

Circular jump in kitchen sink



Atmospheric hydraulic jump



NON-SEDIMENT FLUID-BOUNDARY INTERACTION: MEANDERING CHANNELS IN ICE

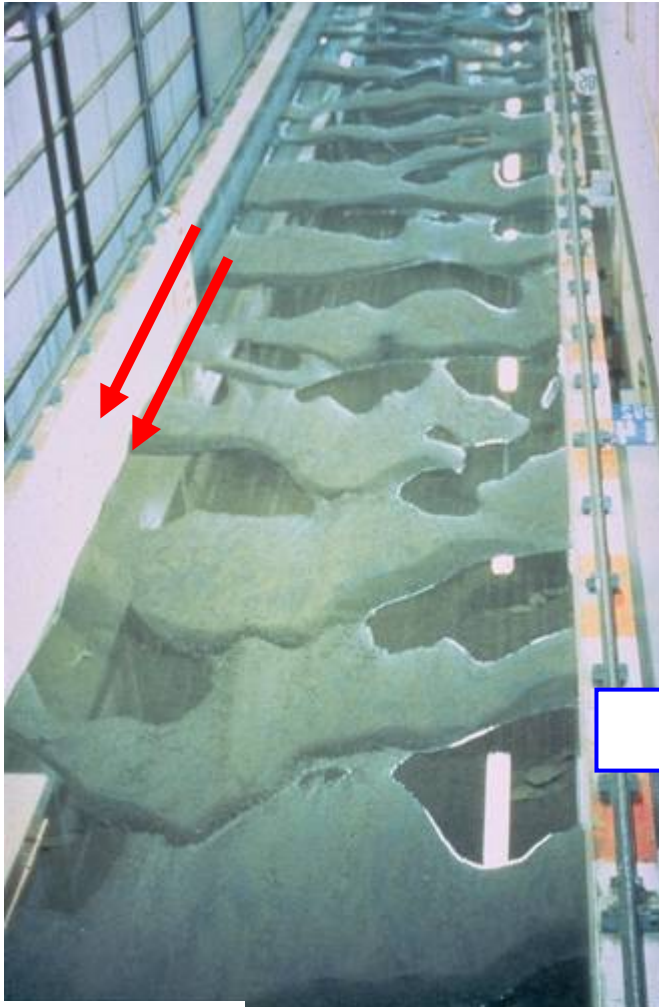
<http://people.whitman.edu/~carsonrj/researchpics/MendenhallAK2.jpg>

[http://crevassezone.org/Photos/Graphics/3054L-\(Meander\).jpg](http://crevassezone.org/Photos/Graphics/3054L-(Meander).jpg)



THE EFFECT OF INCREASING THE WIDTH-DEPTH RATIO B/H

Flume with flow off



Dunes

Tributary of Amazon River



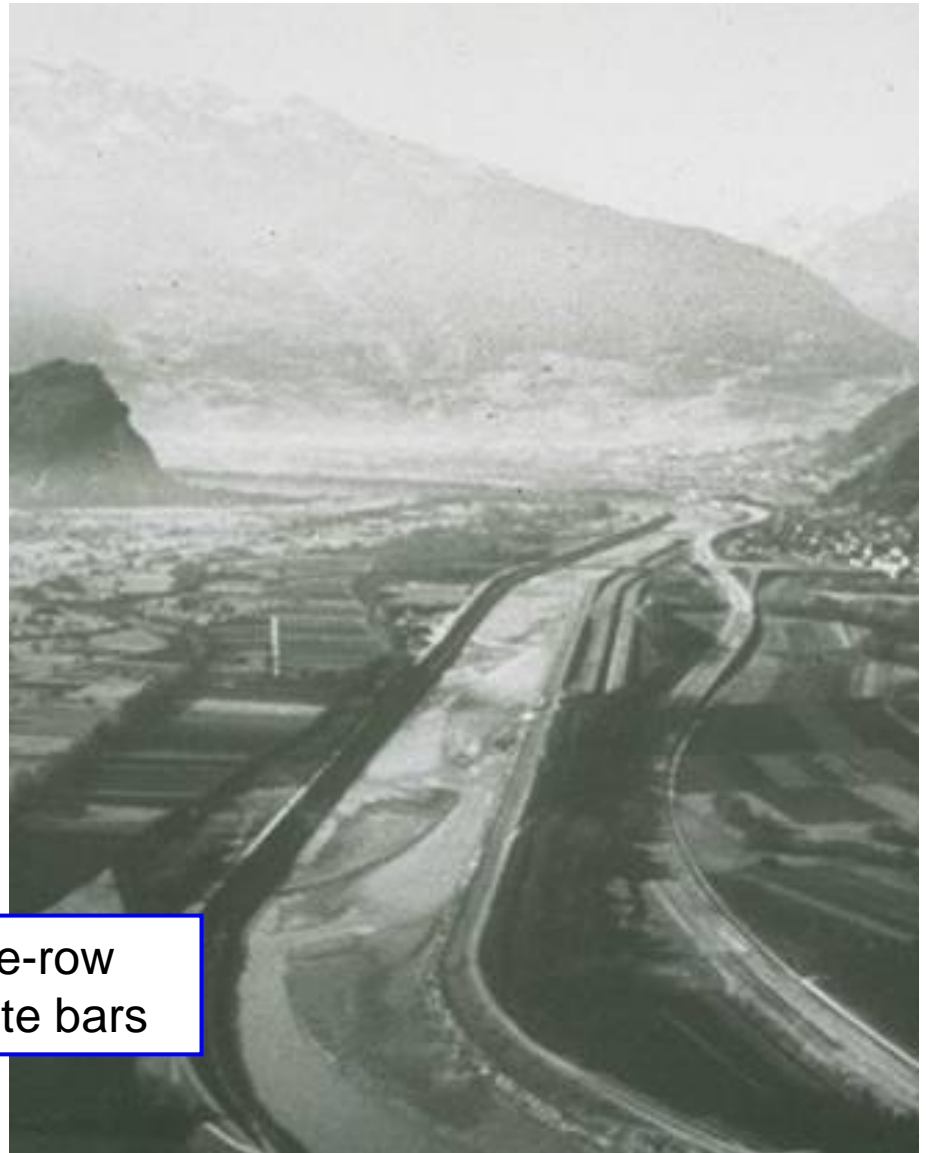
Flume with flow off

B/H → UP

Rhine River, Switzerland



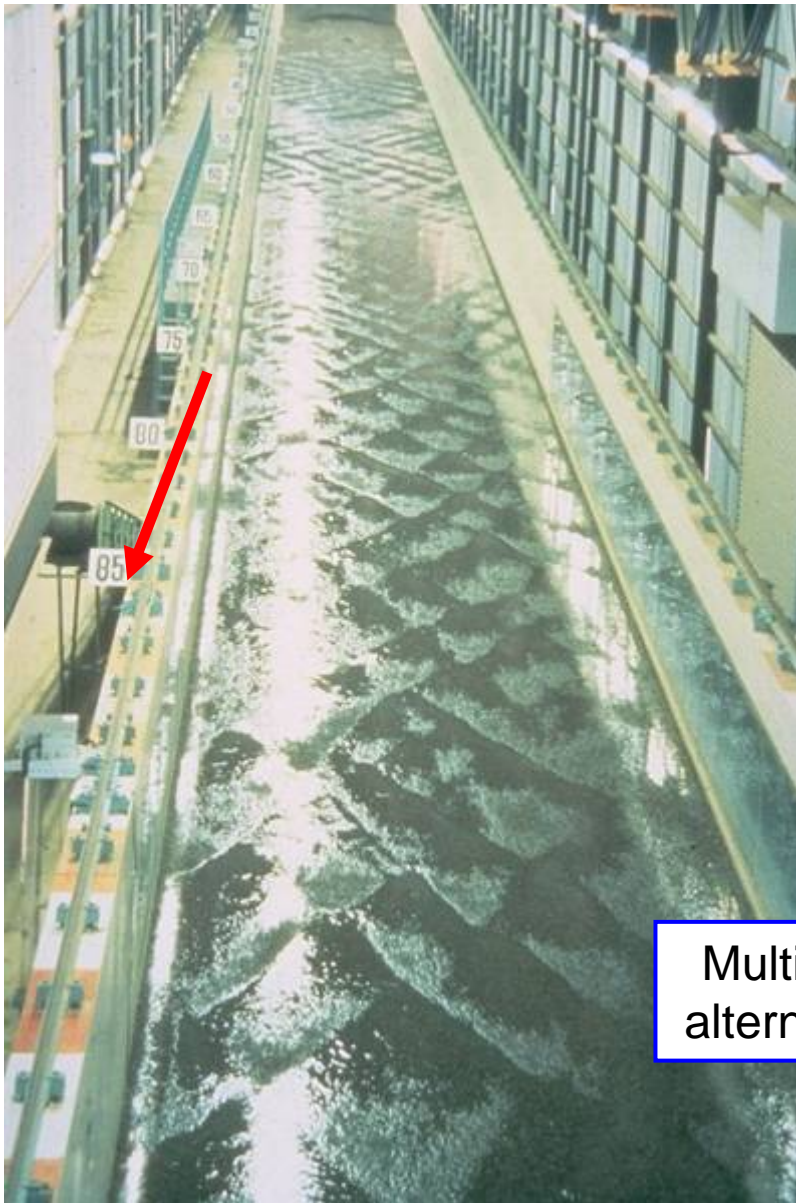
Single-row
alternate bars



Flume with flow off

B/H → UP

Fuefuki River, Japan



Multiple-row
alternate bars



Flume with flow off

B/H → UP

Ohau River, New Zealand



Braiding

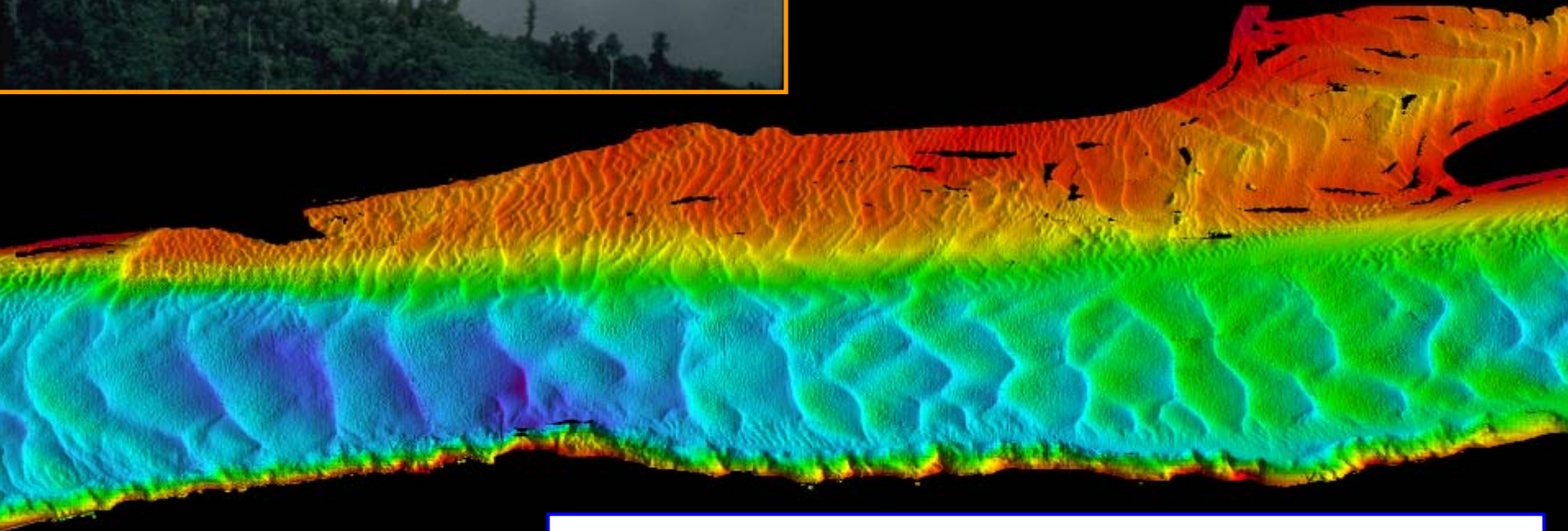


RIVER DUNES



Fly River, Papua New Guinea

Courtesy M. Amsler, J. Best, D. Parsons etc.



Confluence of Parana and Bermejo River, Argentina

DUNES IN THE RHINE DELTA, THE NETHERLANDS

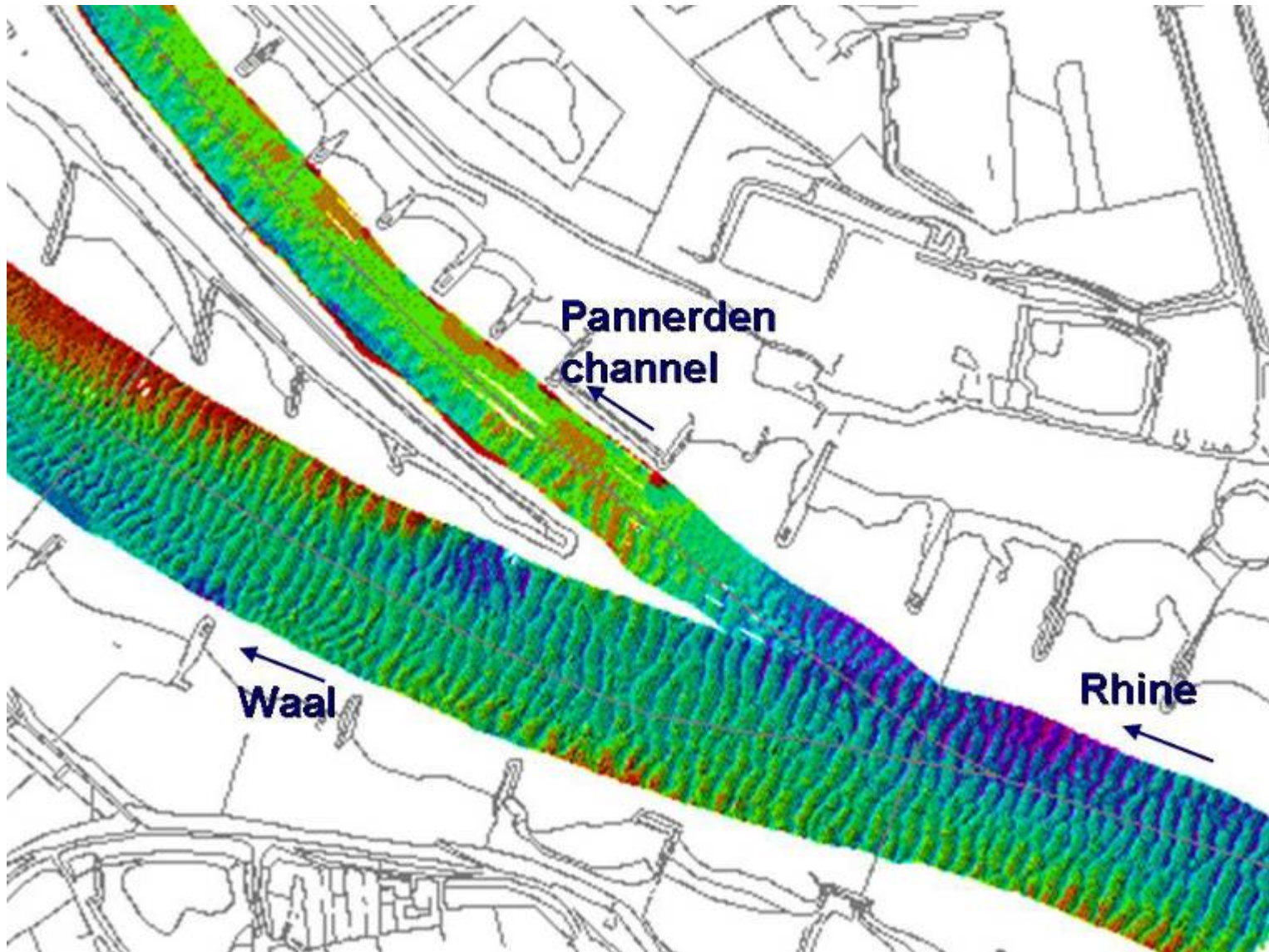
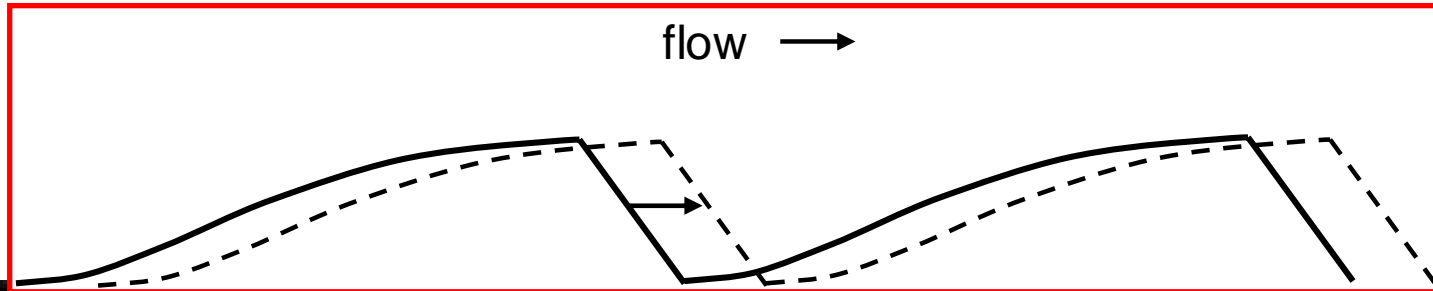
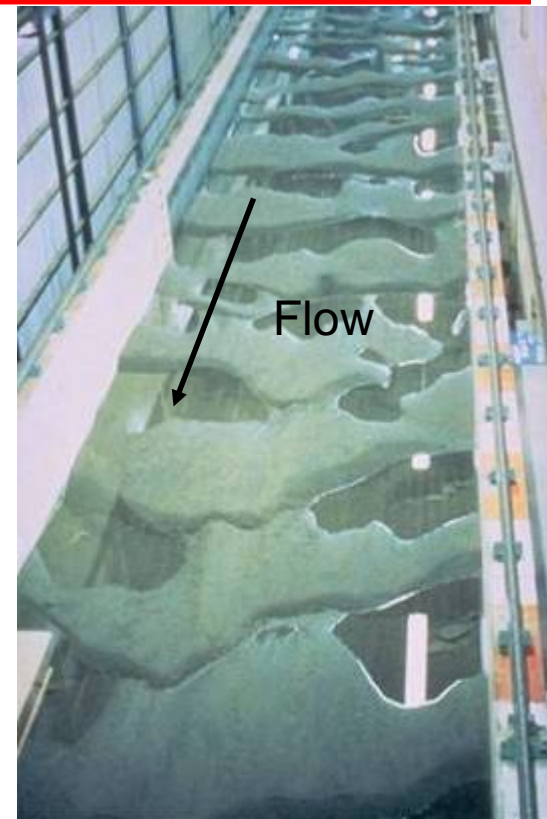


Image courtesy A. Wilbers and A. Blom

DUNE ASYMMETRY



Dunes in a channel at St. Anthony Falls Laboratory, University of Minnesota, USA



Dunes in a channel at Tsukuba University, Japan.
Image courtesy H. Ikeda.

FELIX EXNER: FATHER OF RIVER MORPHODYNAMICS

Exner's Question (1920, 1925): *Why Are Dunes Asymmetric?*

The parameters:

x = streamwise distance [L]

t = time [T]

η = bed elevation [L]

q_t = volume total sediment transport rate per unit stream width [L^2/T]

$\xi = H + \eta$ = water surface elevation [L]

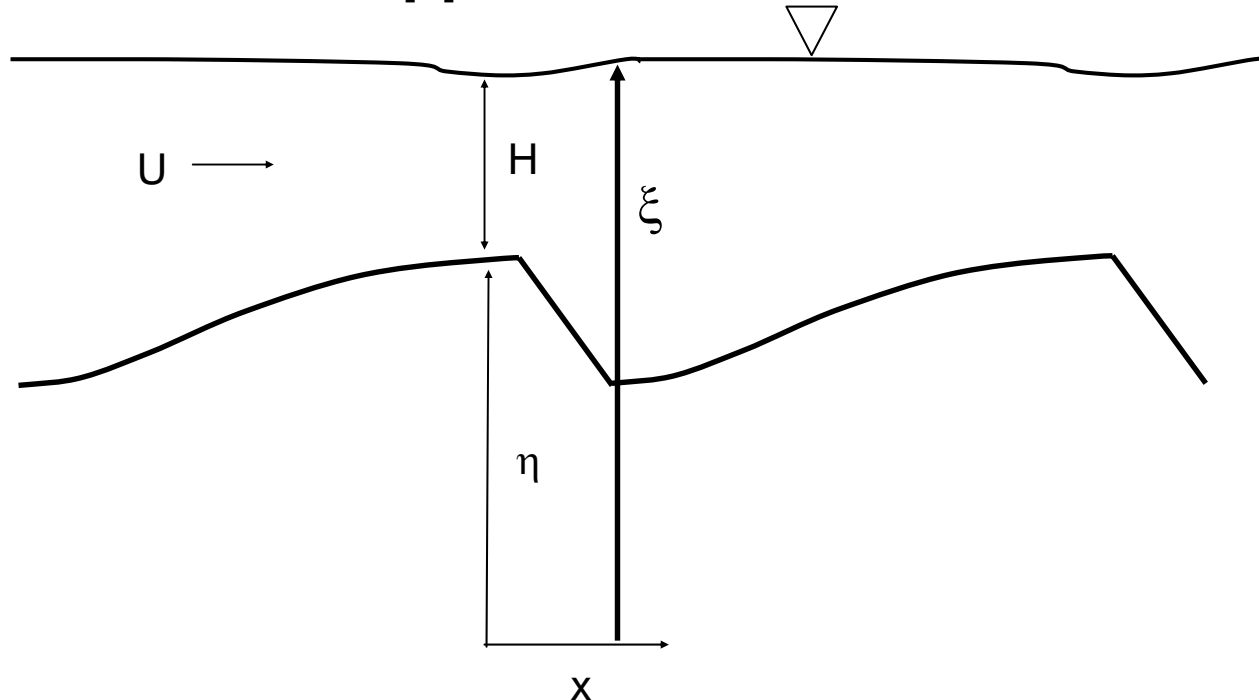
λ_p = bed porosity [1]

g = acceleration of gravity [L/T^2]

H = flow depth [L]

U = depth-averaged flow velocity [L/T]

$q_w = UH$ = water discharge per unit stream width [L^2/T]



OCCAM'S RAZOR: THE MIMINAL FORMULATION TO ANSWER THE QUESTION

Shallow-water inviscid equations of mass and momentum balance

$$\frac{\partial H}{\partial t} + \frac{\partial UH}{\partial x} = 0 \quad \rightarrow \quad UH = q_w = \text{constant}$$

$$\frac{\partial UH}{\partial t} + \frac{\partial U^2H}{\partial x} = -\frac{1}{2}gH\frac{\partial H}{\partial x} - \frac{1}{2}gH\frac{\partial \eta}{\partial x}$$

Quasi-steady assumption: $q_t/q_w \ll 1$

Exner's equation of conservation of bed sediment:

$$(1 - \lambda_p) \frac{\partial \eta}{\partial t} = -\frac{\partial q_t}{\partial x}$$

Relation between sediment transport rate and flow hydraulics:

$$q_t = q_t(U) = \alpha U^n, \quad n > 0$$

Exner's seminal contribution: if more sediment enters a reach than leaves, the bed elevation in the reach increases.

The phenomenon of sediment transport was poorly known in Exner's time. Exner guessed that a higher velocity caused a higher sediment transport rate.

REDUCTION

$$UH = q_w \quad \text{and} \quad H = \xi - \eta \quad \rightarrow \quad U = \frac{q_w}{\xi - \eta}$$

$$\frac{\partial U^2 H}{\partial x} = -\frac{1}{2} g H \frac{\partial H}{\partial x} - \frac{1}{2} g H \frac{\partial \eta}{\partial x} \quad \text{and} \quad H = \xi - \eta \quad \text{and} \quad U = \frac{q_w}{\xi - \eta}$$

→

$$\frac{\partial H}{\partial x} = -\frac{1}{(1 - \mathbf{Fr}^2)} \frac{\partial \eta}{\partial x} \quad \text{and} \quad \frac{\partial \xi}{\partial x} = -\frac{\mathbf{Fr}^2}{(1 - \mathbf{Fr}^2)} \frac{\partial \eta}{\partial x}$$

where

$$\mathbf{Fr} = \frac{U}{\sqrt{gH}} = \text{Froude number}$$

Range for dunes: low Froude number: $\mathbf{Fr}^2 \ll 1$

$$\frac{\partial \xi}{\partial x} = -\frac{\mathbf{Fr}^2}{(1 - \mathbf{Fr}^2)} \frac{\partial \eta}{\partial x} \cong 0 \quad \therefore$$

Constant water surface elevation

MORE REDUCTION

$$U = \frac{q_w}{\xi - \eta} \quad \text{and} \quad q_t = \alpha U^n \quad \text{and} \quad \frac{\partial H}{\partial x} = -\frac{1}{1 - \mathbf{Fr}^2} \frac{\partial \eta}{\partial x}$$

and $\mathbf{Fr}^2 \ll 1$ and $\xi = \text{constant}$

substituted into

$$(1 - \lambda_p) \frac{\partial \eta}{\partial t} = -\frac{\partial q_t}{\partial x}$$

yields

$$\frac{\partial \eta}{\partial t} + c(\eta) \frac{\partial \eta}{\partial x} = 0$$

$$c(\eta) = \alpha q_w^n (\xi - \eta)^{-(n+1)}$$

Since $\xi = \text{constant}$, $q_w = \text{constant}$ and $n > 0$,

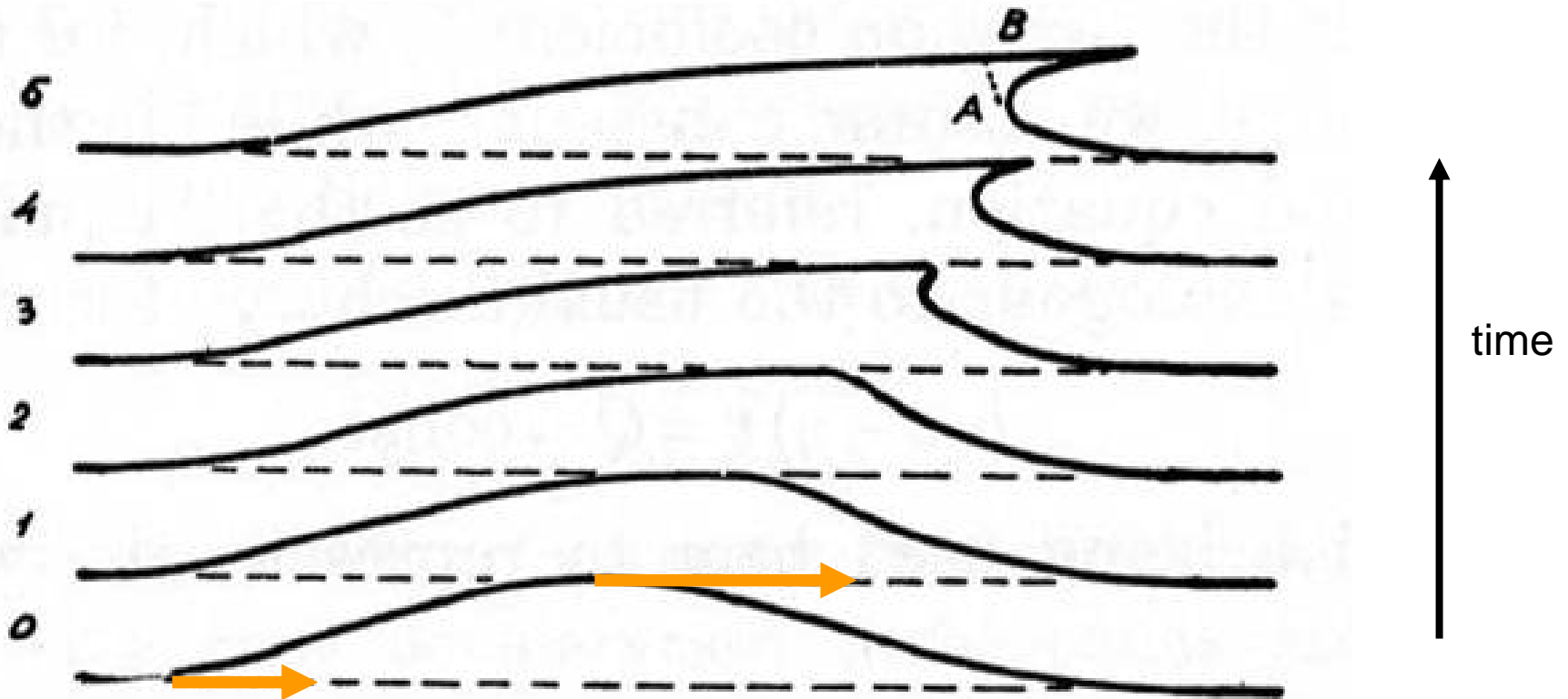
$c > 0$ is an increasing function of η !

THE RESULT

Dunes migrate downstream, and migration speed increases with bed elevation

$$\frac{\partial \eta}{\partial t} + c(\eta) \frac{\partial \eta}{\partial x} = 0$$

And thus the asymmetry!



Exner's original sketch

THE FLOW AND THE BED TALK TO EACH OTHER

The field of **sediment morphodynamics** consists of the class of problems for which the flow over a bed interacts strongly with the shape of the bed, both of which evolve in time.

$$\frac{\partial H}{\partial t} + \frac{\partial UH}{\partial x} = 0$$

$$\frac{\partial UH}{\partial t} + \frac{\partial U^2H}{\partial x} = -\frac{1}{2}gH\frac{\partial H}{\partial x} - gH\frac{\partial \eta}{\partial x}$$

$$(1 - \lambda_p) \frac{\partial \eta}{\partial t} = -\frac{\partial q_t}{\partial x}$$

$$q_t = q_t(U)$$

Quasi-steady assumption:

The flow naturally talks fast, but can also talk slow.

The bed naturally talks slow.

The only part of the flow's talk that the bed hears is the slow part.

(Quasi-steady assumption: $q_t/q_w \ll 1$)

THE CONVERSATION

BLAH*blah*BLAH
*blah*BLAH...

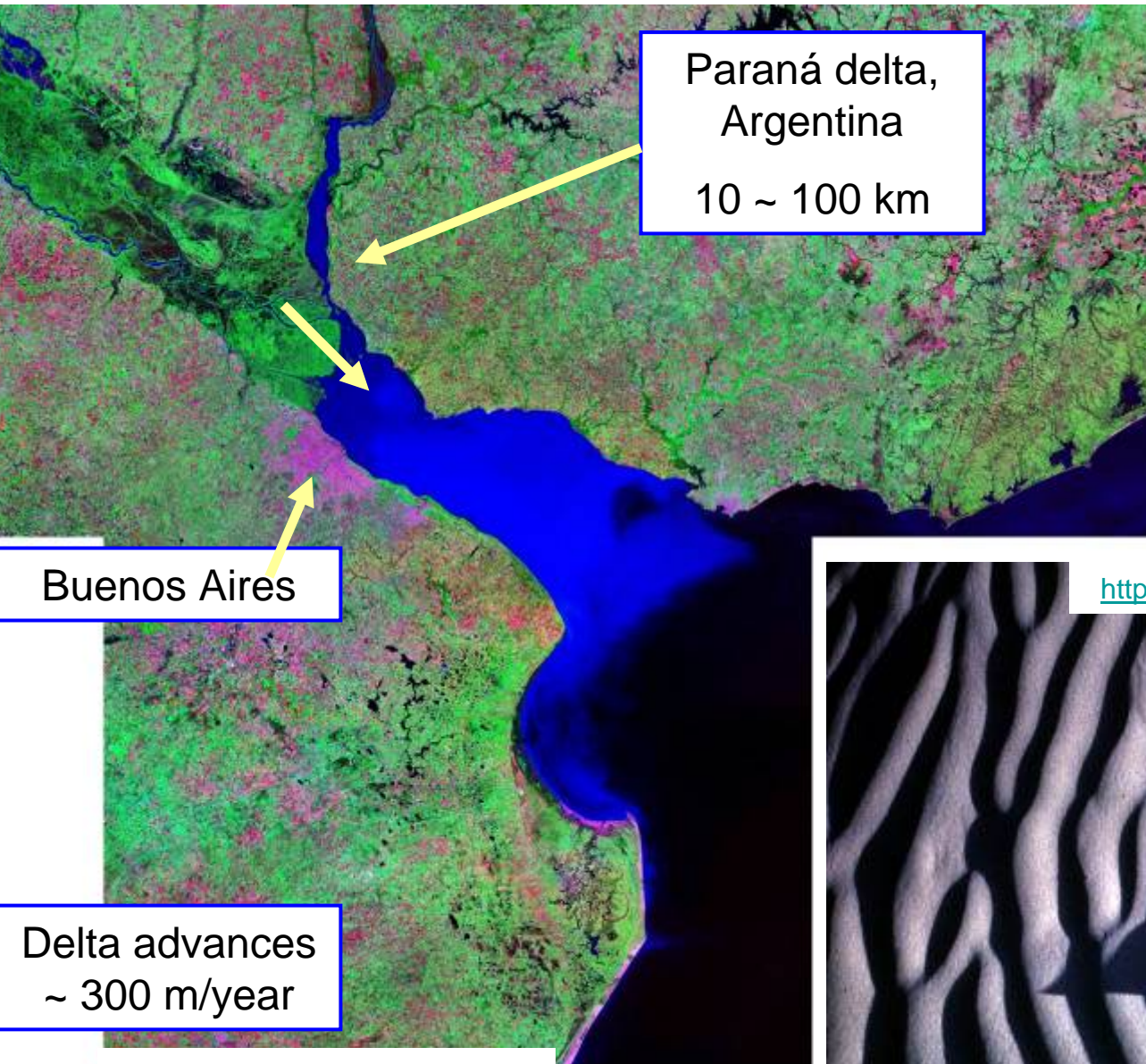
Could you slow down
a bit, I'm having a bit
of trouble trying to
get the gist of it all.

Courtesy IOCC



<http://people.debian.org/~neal/FOSDEM-2005/03-Marcus-IPC-0/03-Marcus-IPC-0.src/conversation.jpg>

SCALES



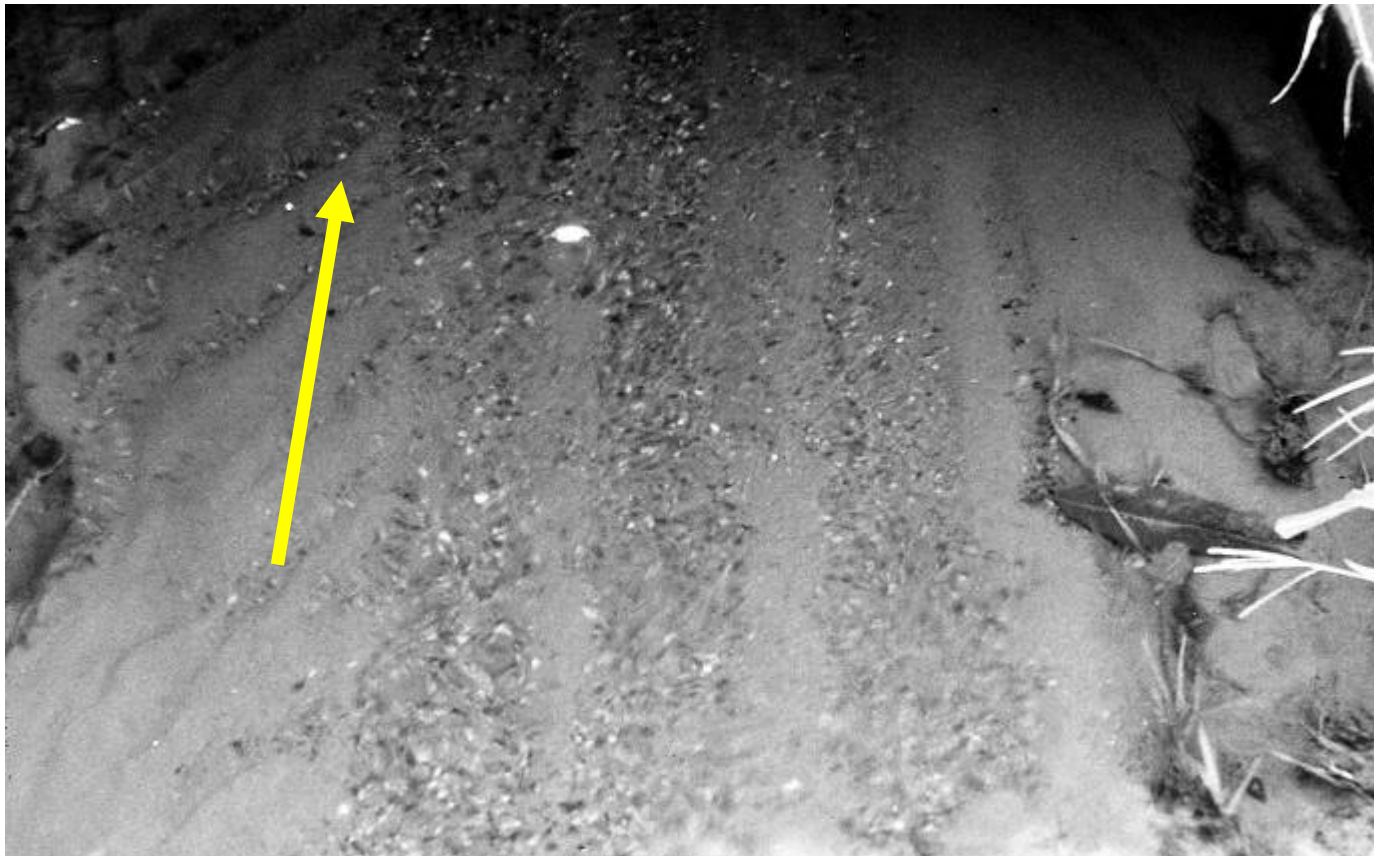
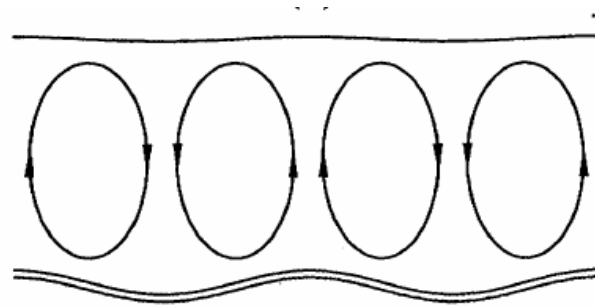
Current ripples:
~ 20 cm
wavelength

<http://www.ux1.eiu.edu/~cfjps/1300/ripples.jpg>

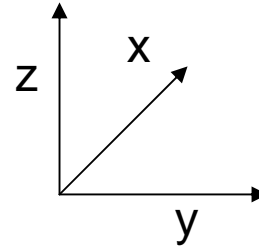
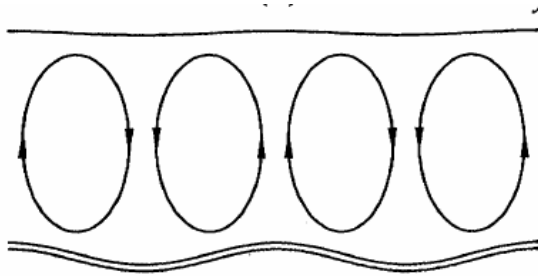


<https://zulu.ssc.nasa.gov/mrsid/>

LONGITUDINAL STREAKS



LONGITUDINAL STREAKS: LINEAR STABILITY ANALYSIS



$$V \frac{\partial U}{\partial y} + W \frac{\partial U}{\partial z} = 1 + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z},$$

$$(1 - \lambda_p) \frac{\partial \eta}{\partial t} = - \frac{\partial q_{tx}}{\partial x} - \frac{\partial q_{tz}}{\partial z}$$

$$V \frac{\partial V}{\partial y} + W \frac{\partial V}{\partial z} = - \frac{\partial P}{\partial y} + \frac{\partial \tau_{yy}}{\partial y} + \frac{\partial \tau_{yz}}{\partial z},$$

$$V \frac{\partial W}{\partial y} + W \frac{\partial W}{\partial z} = - \frac{\partial P}{\partial z} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \tau_{zz}}{\partial z},$$

$$\frac{\partial V}{\partial y} + \frac{\partial W}{\partial z} = 0,$$

3D non-isotropic closure for Reynolds-averaged Navier-Stokes equations: Speziale.

Closure for sediment transport rates q_x and q_z in terms of bed shear stress and slope vectors.

$$\tau_{ij} = -\frac{2}{3}k\delta_{ij} + 2\nu_t D_{ij} + C_D l^2 (D_{im} D_{mj} - \frac{1}{3} D_{mn} D_{mn} \delta_{ij}) + C_E l^2 (\hat{D}_{ij} - \frac{1}{3} \hat{D}_{mm} \delta_{ij}).$$

Here D_{ij} is the mean rate of strain tensor, $\nu_t = \frac{1}{2} k^{1/2} l$ is the eddy viscosity and

$$\hat{D}_{ij} = \frac{D D_{ij}}{D t} - \frac{\partial U_i}{\partial x_k} D_{kj} - \frac{\partial U_j}{\partial x_k} D_{ki}$$

DUNES, ANTIDUNES

Dunes: Middle
Loup River,
Nebraska



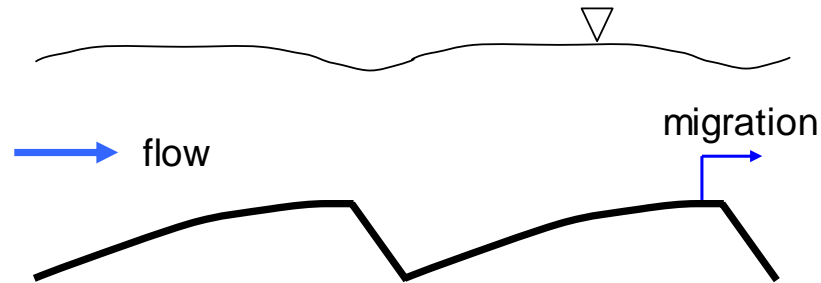
Image courtesy D. Mohrig

Antidunes,
Brittany, France

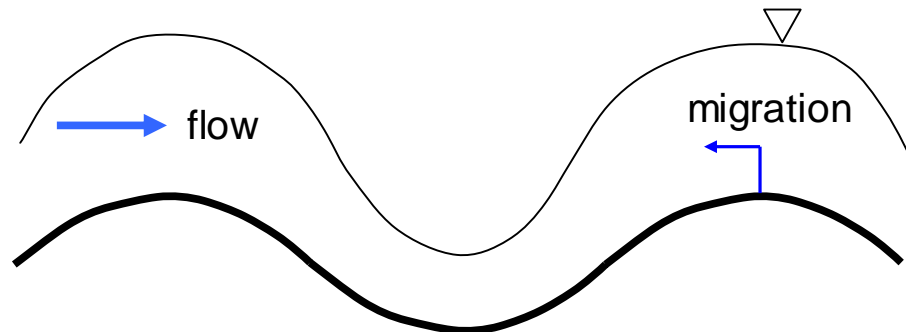


DEFINITION OF DUNES AND ANTIDUNES

Dunes are 1D (or quasi-1D) bedforms for which the water surface fluctuations are approximately *out of phase* with the bed fluctuations. That is, the water surface is high where the bed is low and vice versa. As is shown below dunes migrate downstream.



Antidunes are 1D (or quasi-1D) bedforms for which the water surface fluctuations are approximately *in phase* with the bed fluctuations. That is, the water surface is high where the bed is high and vice versa. As shown below, most antidunes migrate upstream, but there is a regime within which they can migrate downstream.



REGIME DIAGRAM: POTENTIAL FLOW OVER A WAVY BED

x = streamwise direction

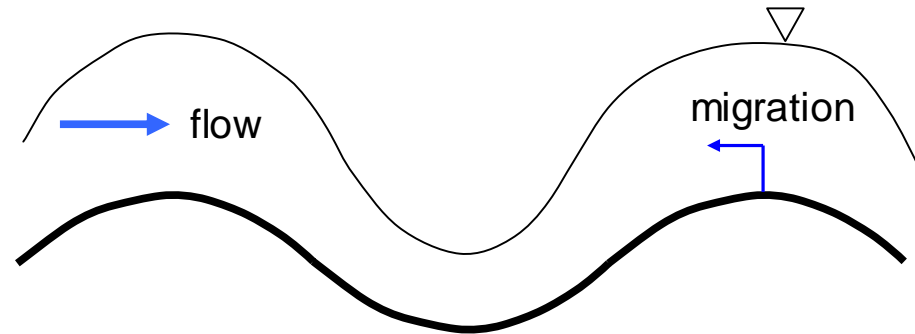
y = vertical direction

u = streamwise velocity

v = vertical velocity

p = pressure

g = gravitational acceleration



η_o = amplitude of bed perturbation

H_o = unperturbed depth

Linearized potential flow analysis is sufficient to explain existence regimes, but not formation (gives neutral stability)

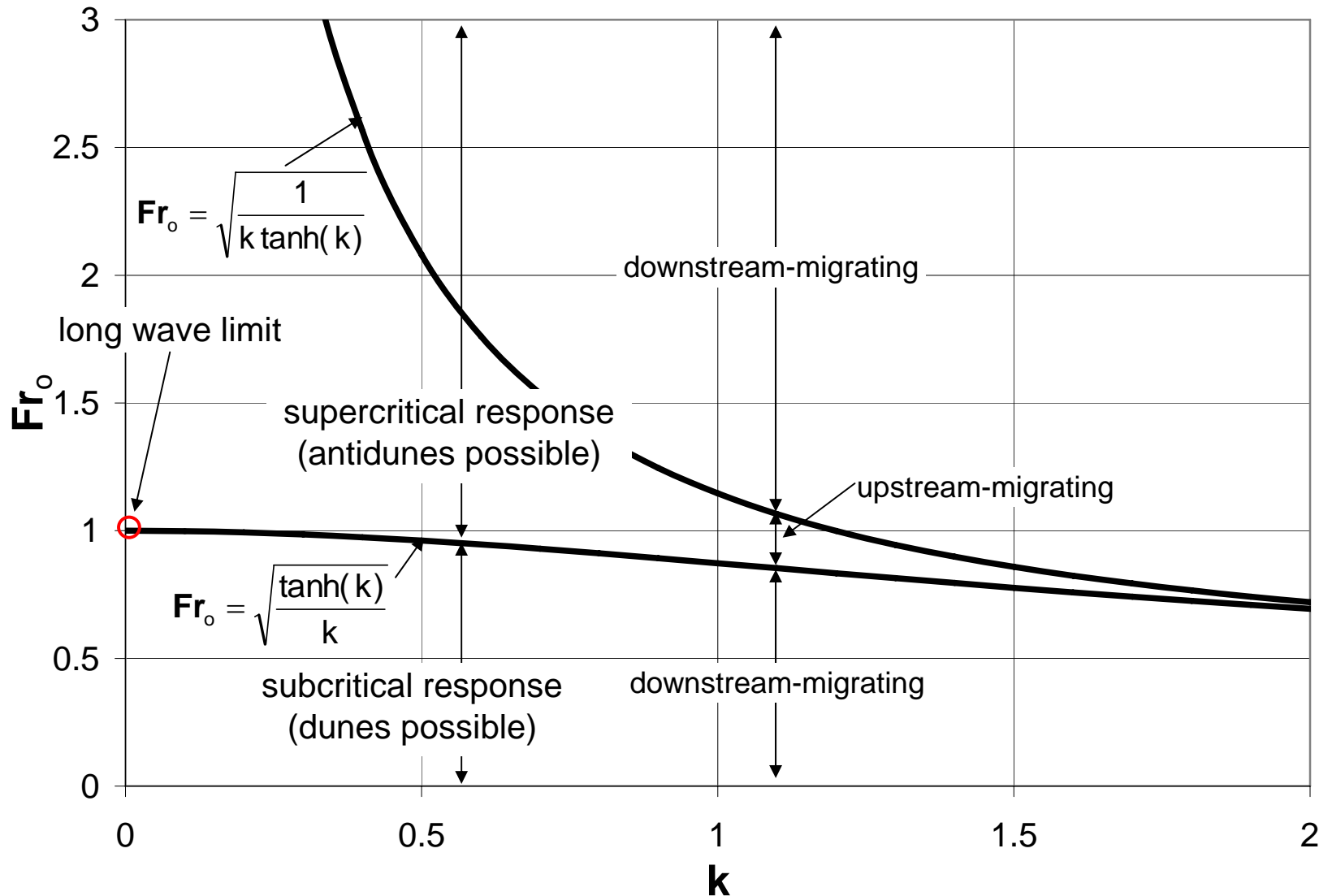
$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial x} - g \frac{\partial \eta}{\partial x}$$

$$u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial y} - g$$

$$\eta = \eta_o \sin(kx) \quad , \quad k = \frac{2\pi H_o}{\lambda}$$

PHASE DIAGRAM FOR DUNES AND ANTIDUNES BASED ON LINEAR POTENTIAL THEORY OVER A WAVY BED



FLOW IN THE DUNE REGIME

$$Fr_o < [\tanh(k)/k]^{1/2}$$

$$k = 2\pi H/\lambda$$

H = depth, λ = wavelength

Water surface is out of phase with the bed.

Depth variation is out of phase with the bed

Flow accelerates from trough to crest.

Sediment transport increases from trough to crest.

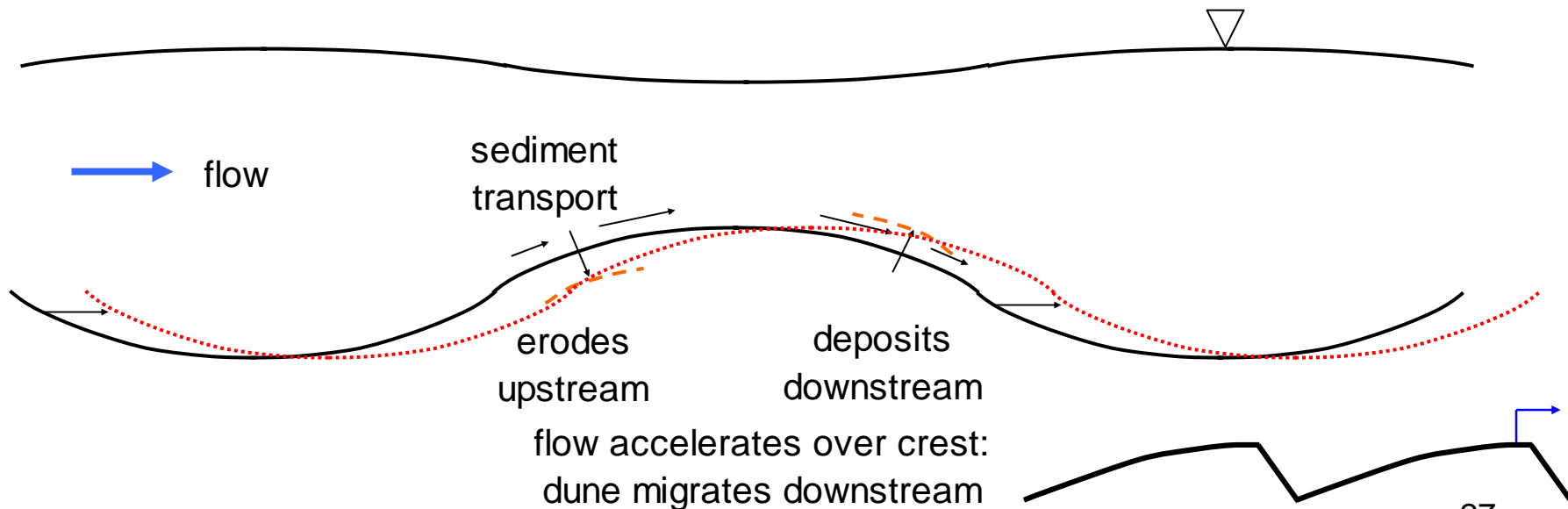
Bedform migrates downstream.

Bedform becomes asymmetric.

$k \rightarrow 0$ (shallow water)

$$\frac{\partial H}{\partial x} = -\frac{1}{(1 - Fr^2)} \frac{\partial \eta}{\partial x}$$

$$\frac{\partial \xi}{\partial x} = -\frac{Fr^2}{(1 - Fr^2)} \frac{\partial \eta}{\partial x}$$



FLOW IN THE UPSTREAM-MIGRATING ANTIDUNE REGIME

$$[\tanh(k)/k]^{1/2} < \mathbf{Fr}_o < [k \tanh(k)]^{-1/2}$$

Water surface is in phase with the bed.

Depth variation is in phase with the bed

Flow decelerates from trough to crest.

Sediment transport decreases from trough to crest.

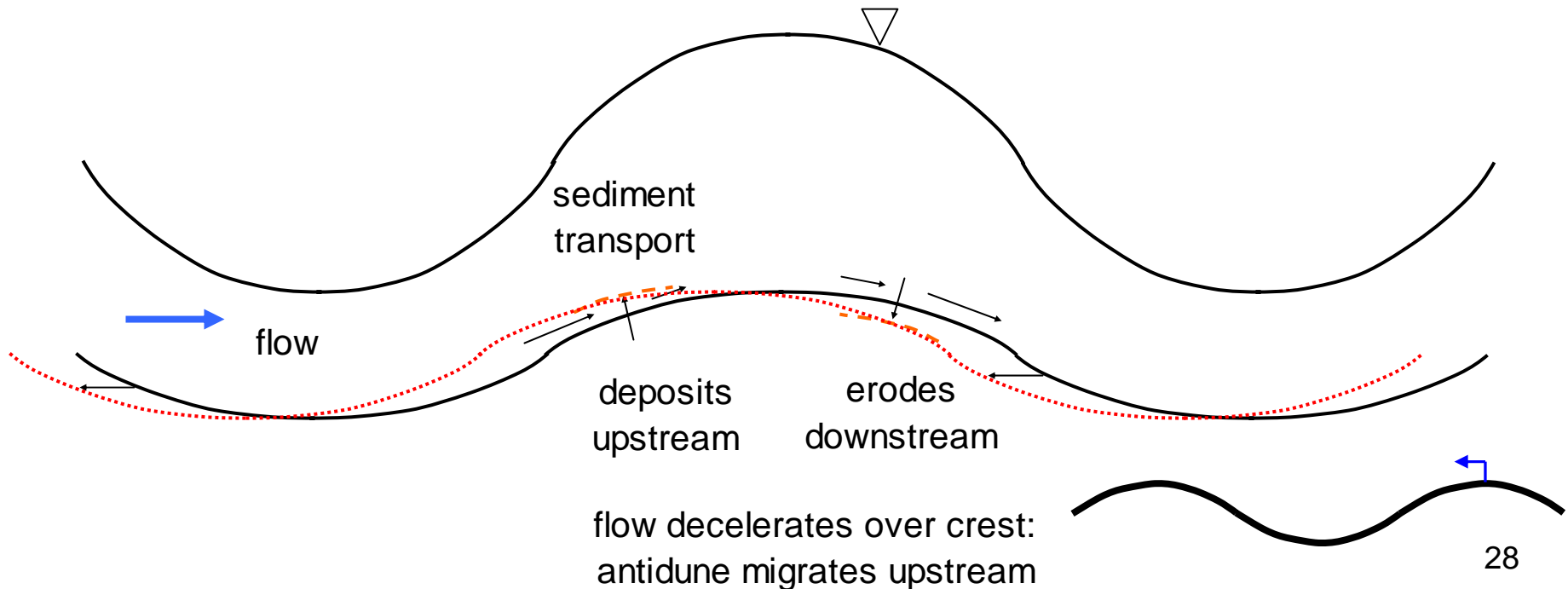
Bedform migrates upstream (or hardly at all).

Bedform stays symmetric.

$k \rightarrow 0$ (shallow water)

$$\frac{\partial H}{\partial x} = -\frac{1}{(1 - \mathbf{Fr}^2)} \frac{\partial \eta}{\partial x}$$

$$\frac{\partial \xi}{\partial x} = -\frac{\mathbf{Fr}^2}{(1 - \mathbf{Fr}^2)} \frac{\partial \eta}{\partial x}$$



FLOW IN THE DOWNSTREAM-MIGRATING ANTIDUNE REGIME

$$[k \tanh(k)]^{-1/2} < Fr_0$$

Water surface is in phase with the bed.

Depth variation is out of phase with the bed.

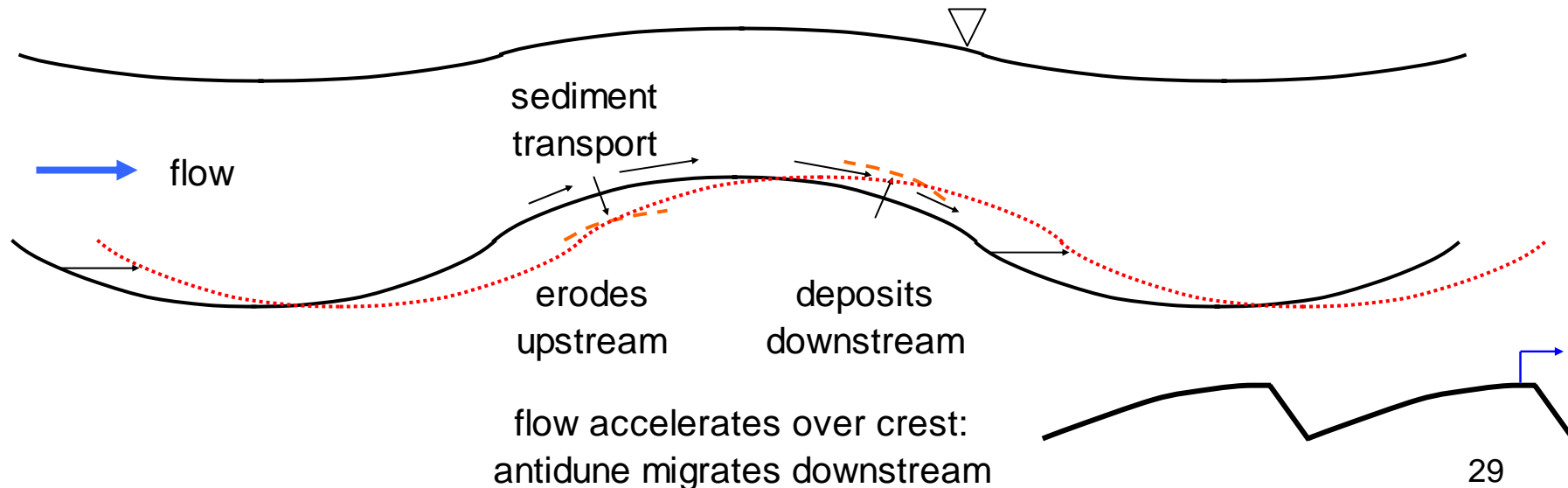
Flow accelerates from trough to crest.

Sediment transport increases from trough to crest.

Bedform migrates downstream.

Bedform becomes asymmetric.

These are antidunes that look like dunes: not too common, but they are observed.



STABILITY ANALYSIS FOR DUNES AND ANTIDUNE FORMATION: OCCAM'S RAZOR

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial x} - g \frac{\partial \eta}{\partial x} + \frac{\partial}{\partial x} \left(v_t \frac{\partial u}{\partial x} \right) + \frac{\partial}{\partial x} \left(v_t \frac{\partial u}{\partial y} \right) - g \frac{\partial \eta}{\partial x}$$

$$u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial y} - g + \frac{\partial}{\partial x} \left(v_t \frac{\partial v}{\partial x} \right) + \frac{\partial}{\partial x} \left(v_t \frac{\partial v}{\partial y} \right) - g \frac{\partial \eta}{\partial x}$$

$$\eta = \eta_o e^{\alpha t} \sin(kx) \quad , \quad k = \frac{2\pi H_o}{\lambda}$$

Closure for v_t : a constant value that gives a result close to the logarithmic law

$$(1 - \lambda_p) \frac{\partial \eta}{\partial t} = -\frac{\partial q_t}{\partial x}$$

$$q_t = f\left(\tau_b, \frac{\partial \eta}{\partial x}\right)$$

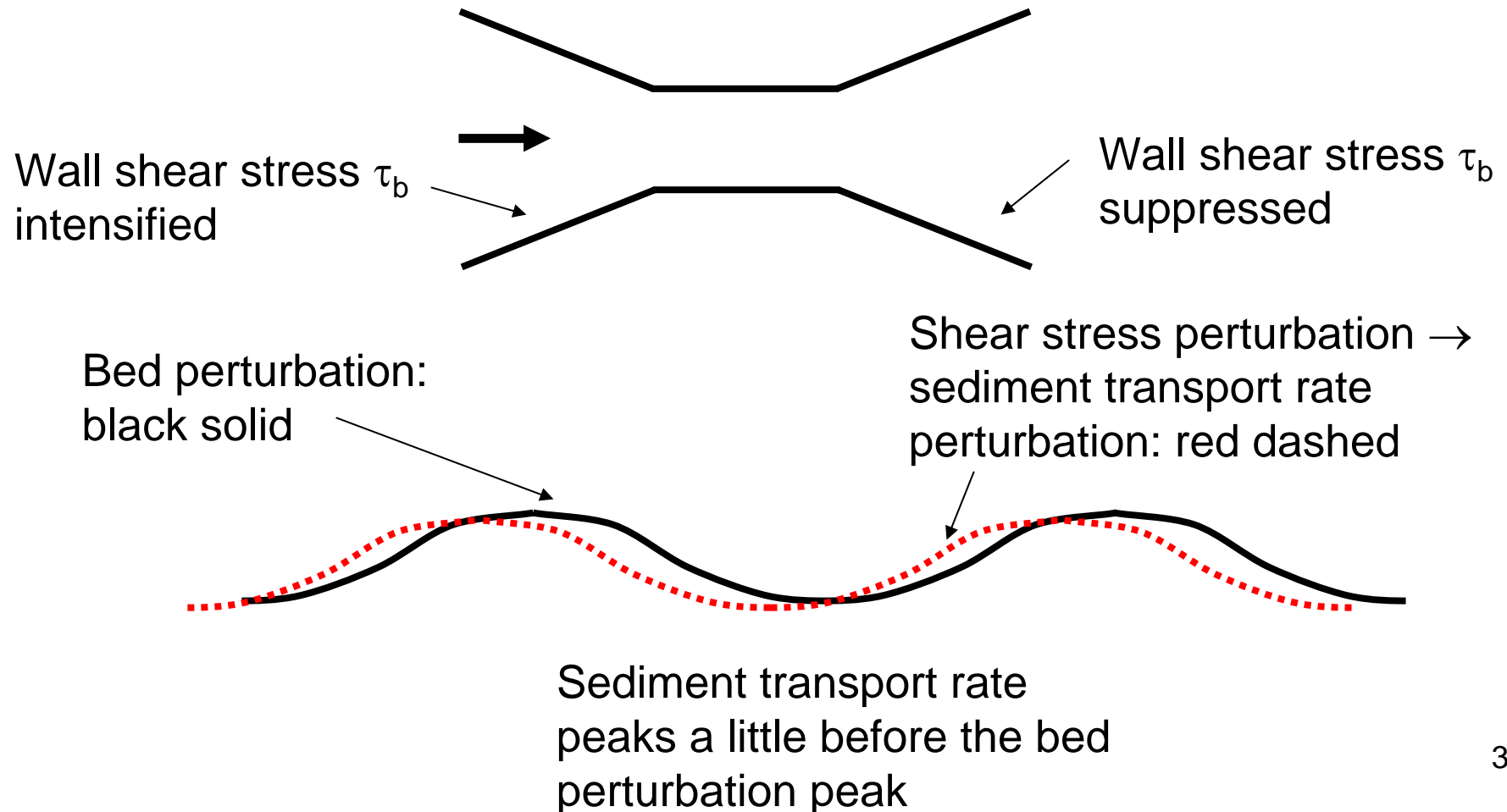
In sediment transport law, τ_b = bed shear stress

INSTABILITY MECHANISM FOR DUNES

Consider flow into a Venturi contraction.

The favorable pressure gradient on the upstream side intensifies the bed shear stress.

The adverse pressure gradient on the downstream side suppresses shear stress.

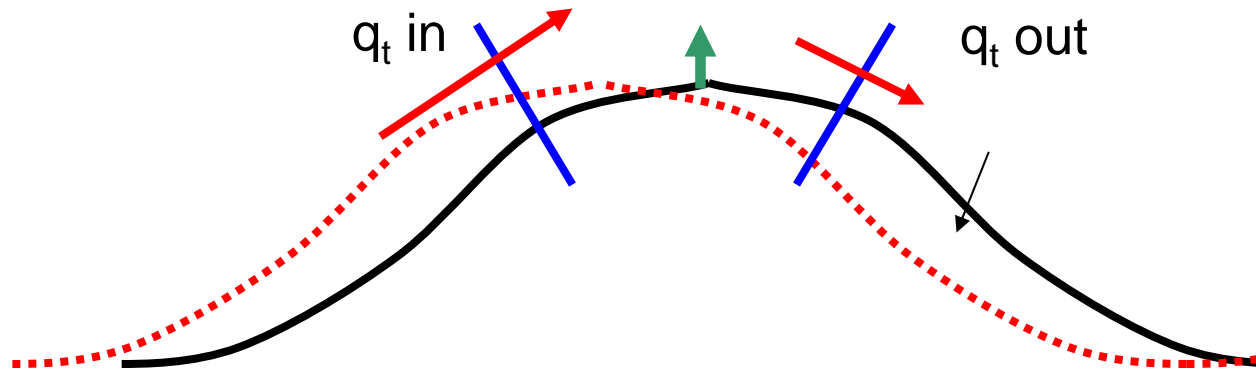


NET DEPOSITION AT APEX: LINEAR MODEL

The bed shear stress perturbation, and thus the sediment transport rate perturbation, lead the bed elevation perturbation.

There is thus net deposition at the apex (and net erosion at the trough), and so amplitude increases in time.

A nonlinear analysis including flow separation on the lee side of dunes is necessary to explain nonlinear equilibrium: numerical, e.g. $k-\varepsilon$



Bed perturbation:
black solid

sediment transport rate
perturbation: red dashed

NONLINEAR PHENOMENON OF ANTIDUNES



SINGLE-ROW AND MULTIPLE-ROW ALTERNATE BARS

Occam's razor minimal analysis:
2D shallow water equations +
2D sediment transport formulation

Controlling parameter: width-depth ratio B/H
No bars \rightarrow single-row bars \rightarrow multiple-row bars

Naka River, Japan



Image courtesy S. Ikeda

Hii River, Japan



Image courtesy H.
Takebayashi

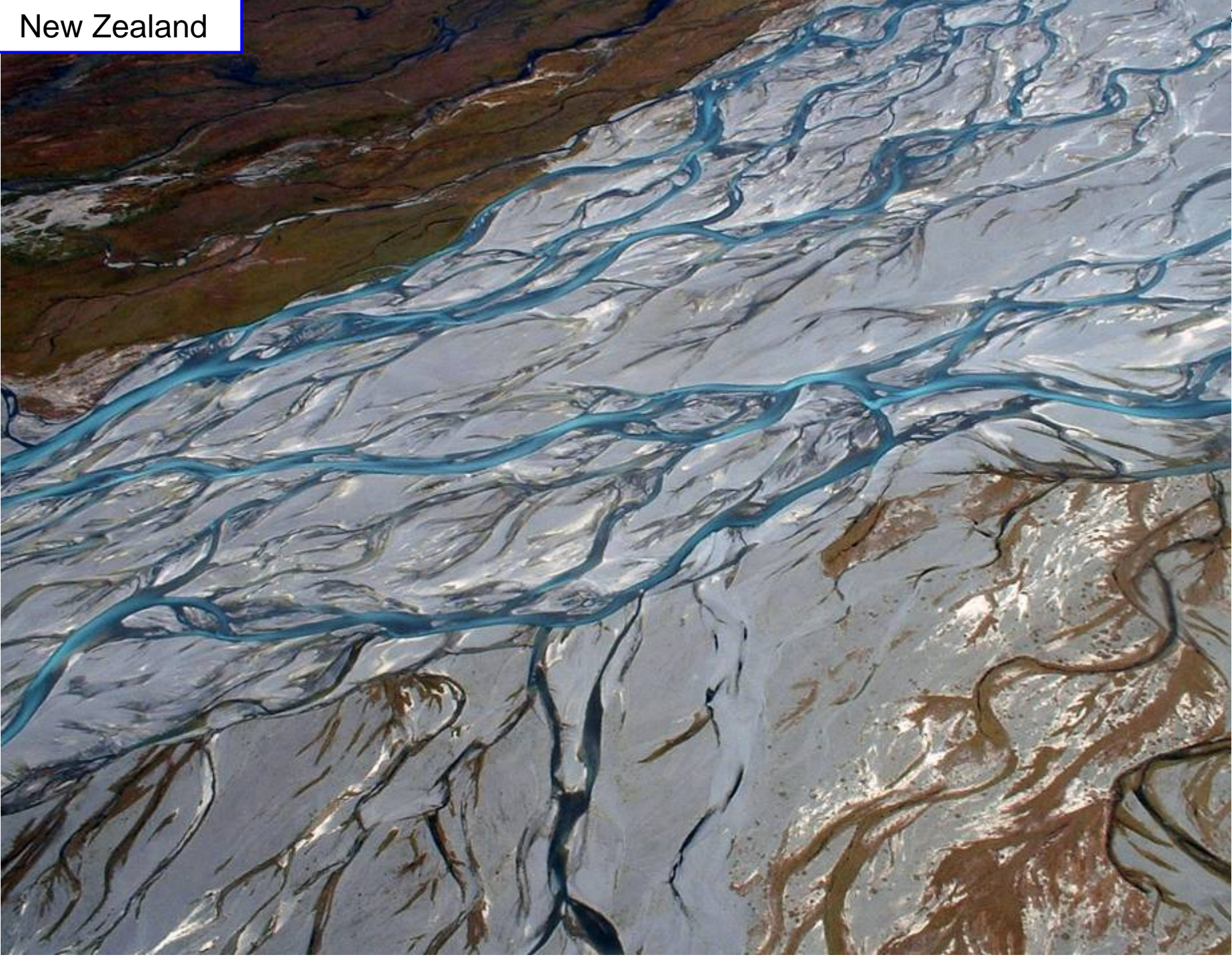
New Zealand

BRAIDING



New Zealand







BRAIDING MECHANISM: CONFLUENCES



Skeithara Sandur, Iceland

Image courtesy H. Johannesson

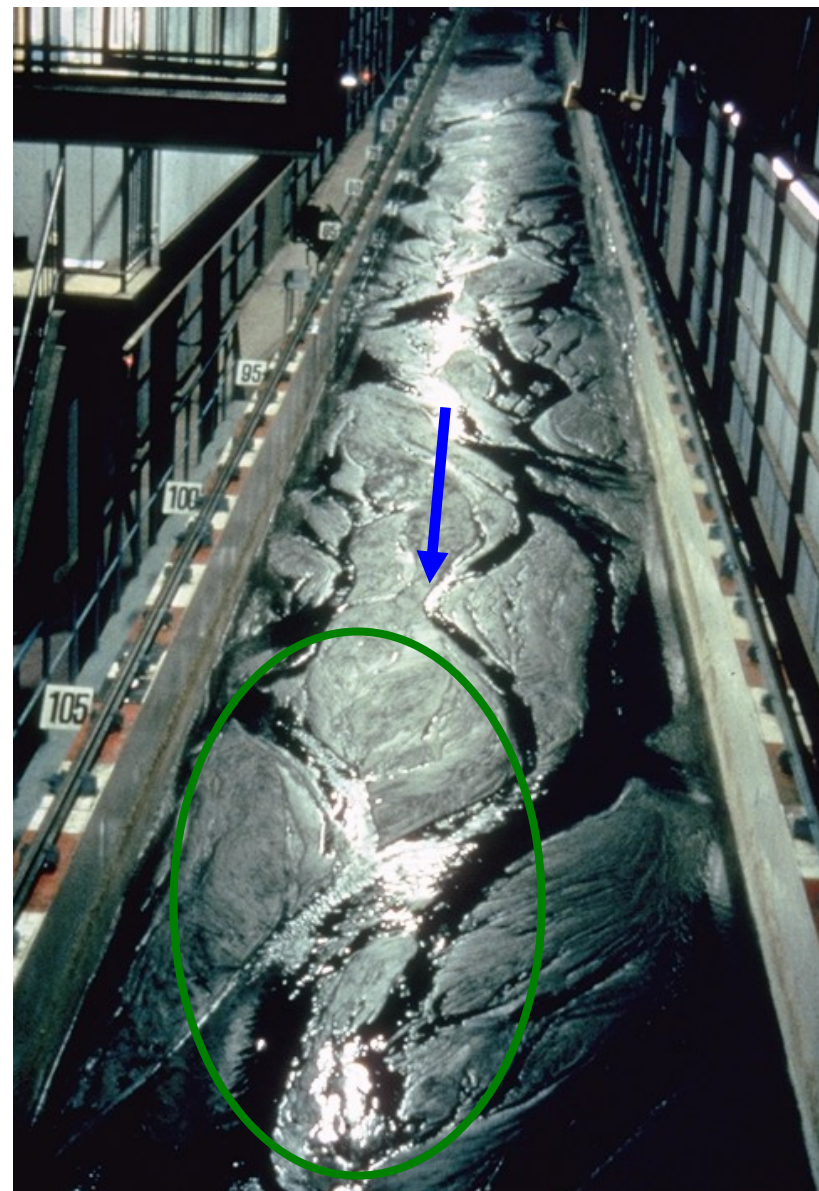
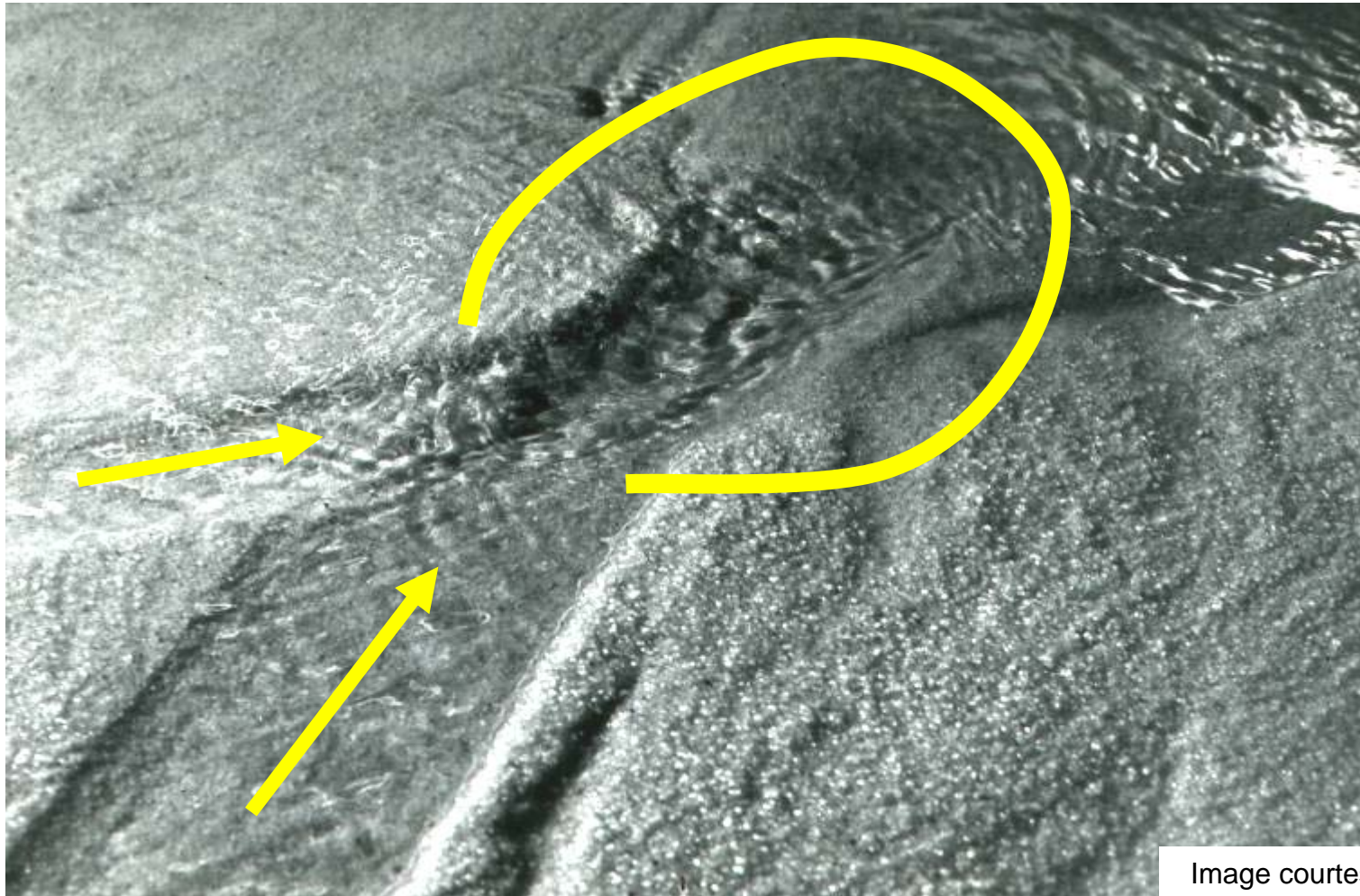


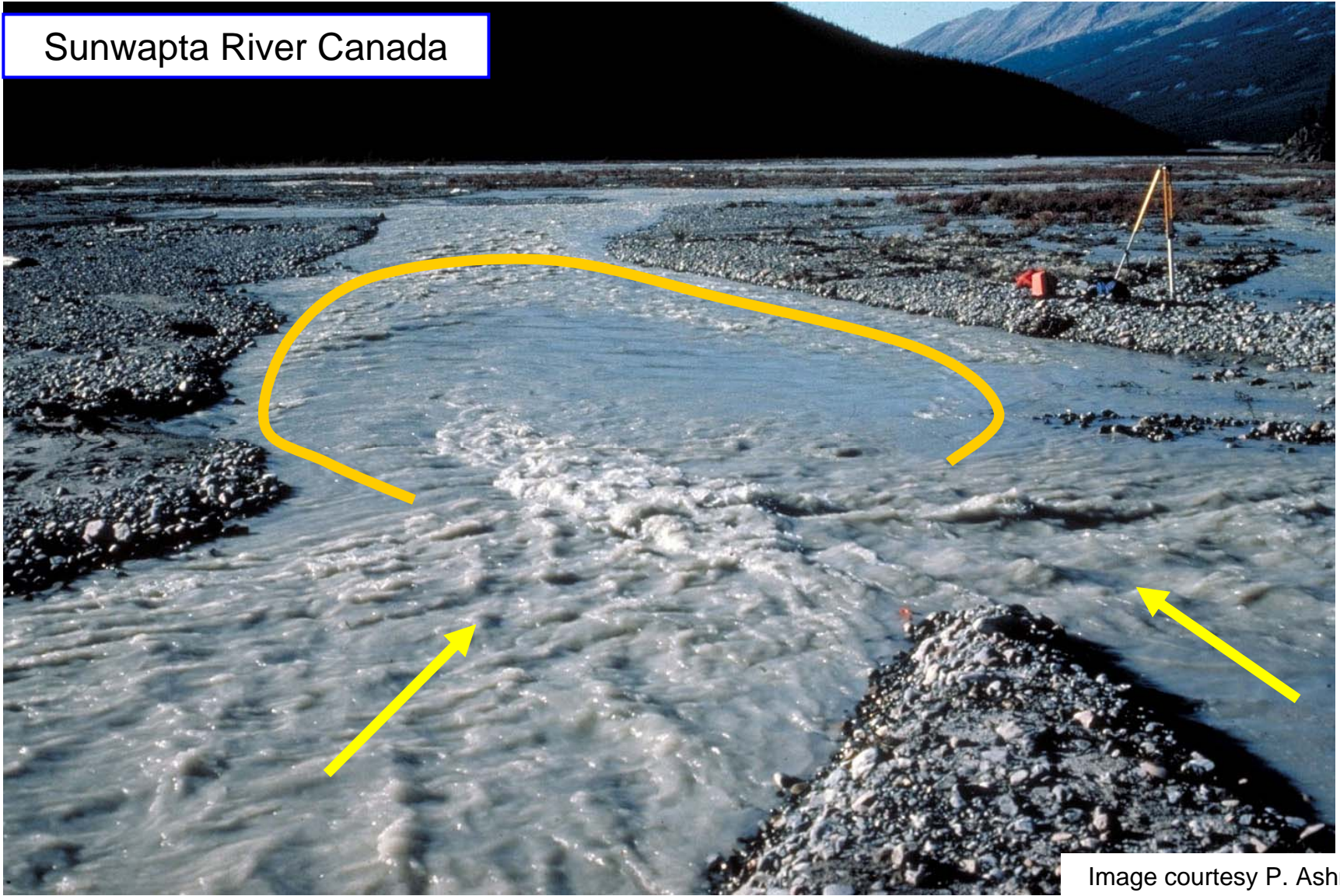
Image courtesy H. Ikeda

NONLINEAR INTENSIFICATION OF SEDIMENT TRANSPORT RATE AT CONFLUENCES: SCOUR



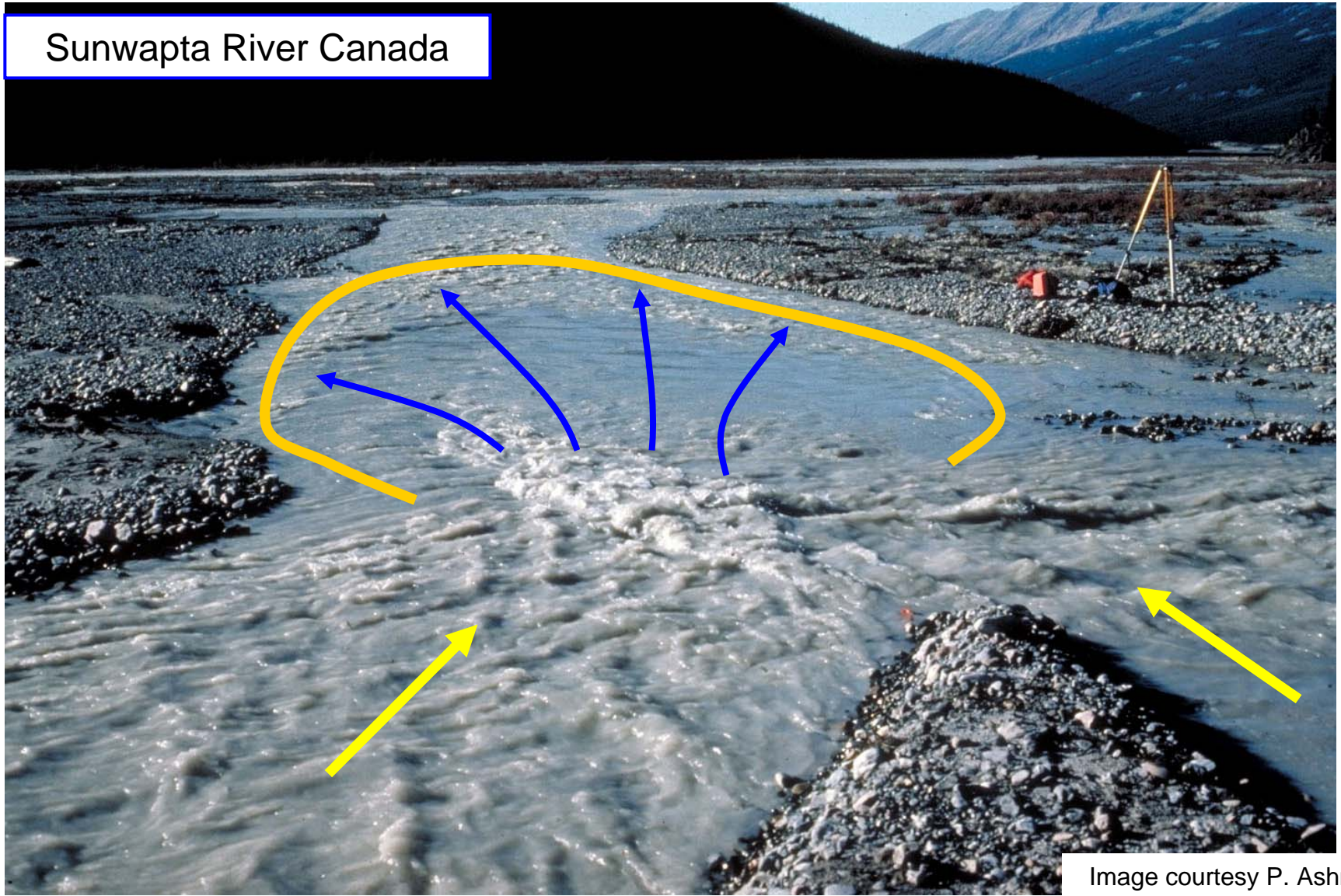
DOWNSTREAM, THE FLOW EXPANDS AND DEPOSITS A MINI-FAN

Sunwapta River Canada



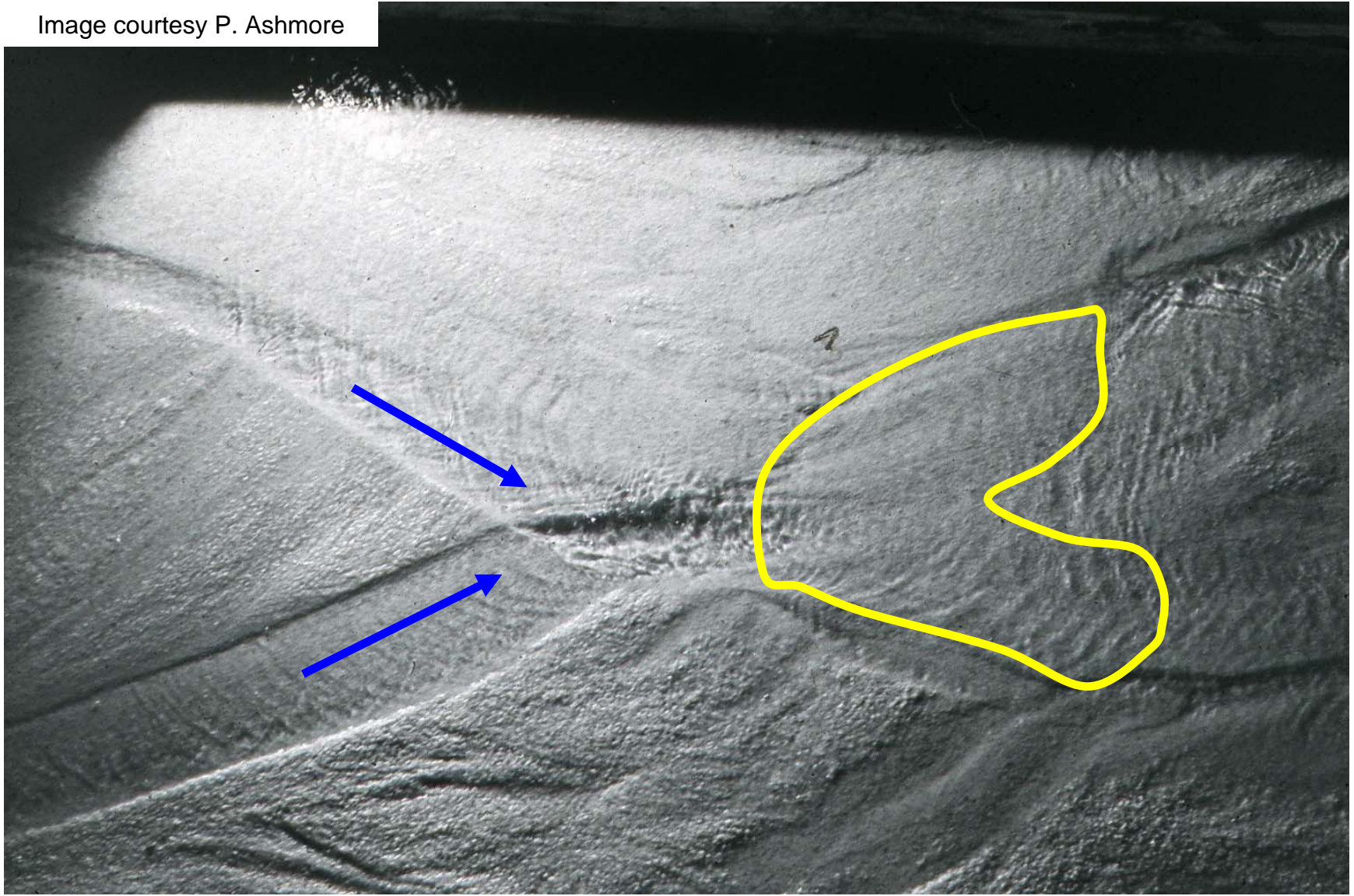
DOWNSTREAM, THE FLOW EXPANDS AND DEPOSITS A MINI-FAN

Sunwapta River Canada



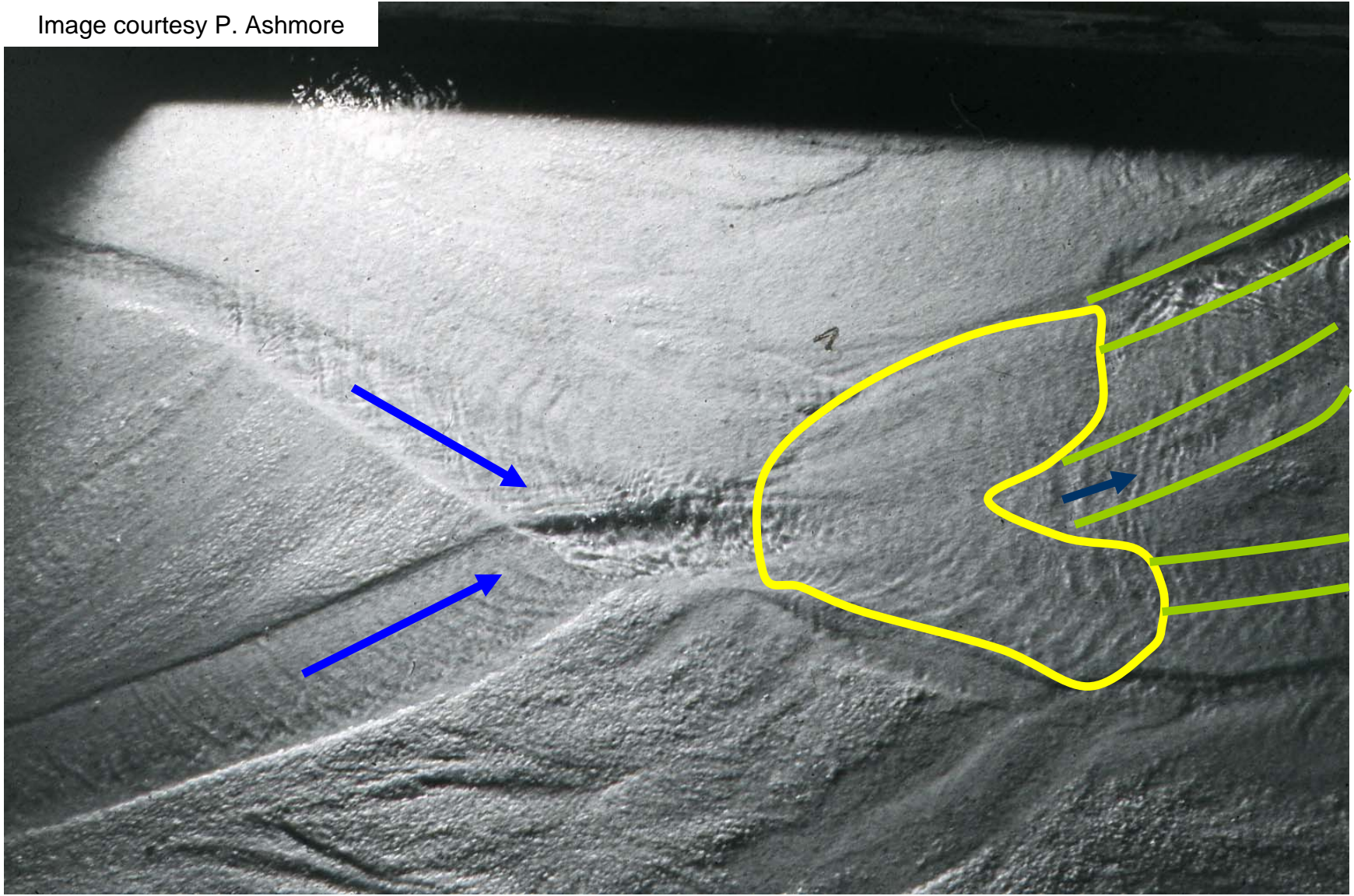
AS THE FLOW GETS WIDER AND SHALLOWER, IT BECOMES UNSTABLE AND BIFURCATES INTO ONE OR MORE CHANNELS

Image courtesy P. Ashmore



AS THE FLOW GETS WIDER AND SHALLOWER, IT BECOMES UNSTABLE AND BIFURCATES INTO ONE OR MORE CHANNELS

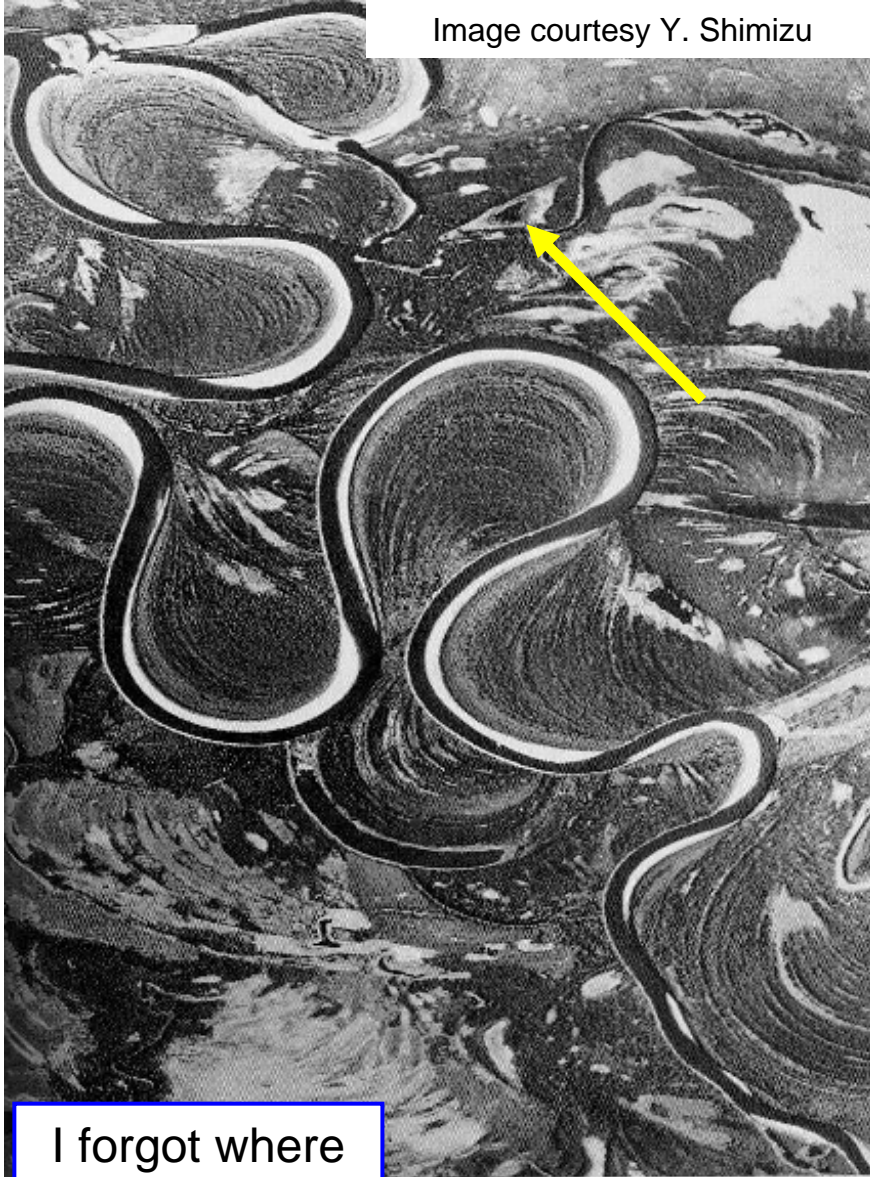
Image courtesy P. Ashmore



MEANDERING

Mississippi River, USA

Image courtesy Y. Shimizu



From maps of H. Fisk



I forgot where

MEANDERING MECHANISM

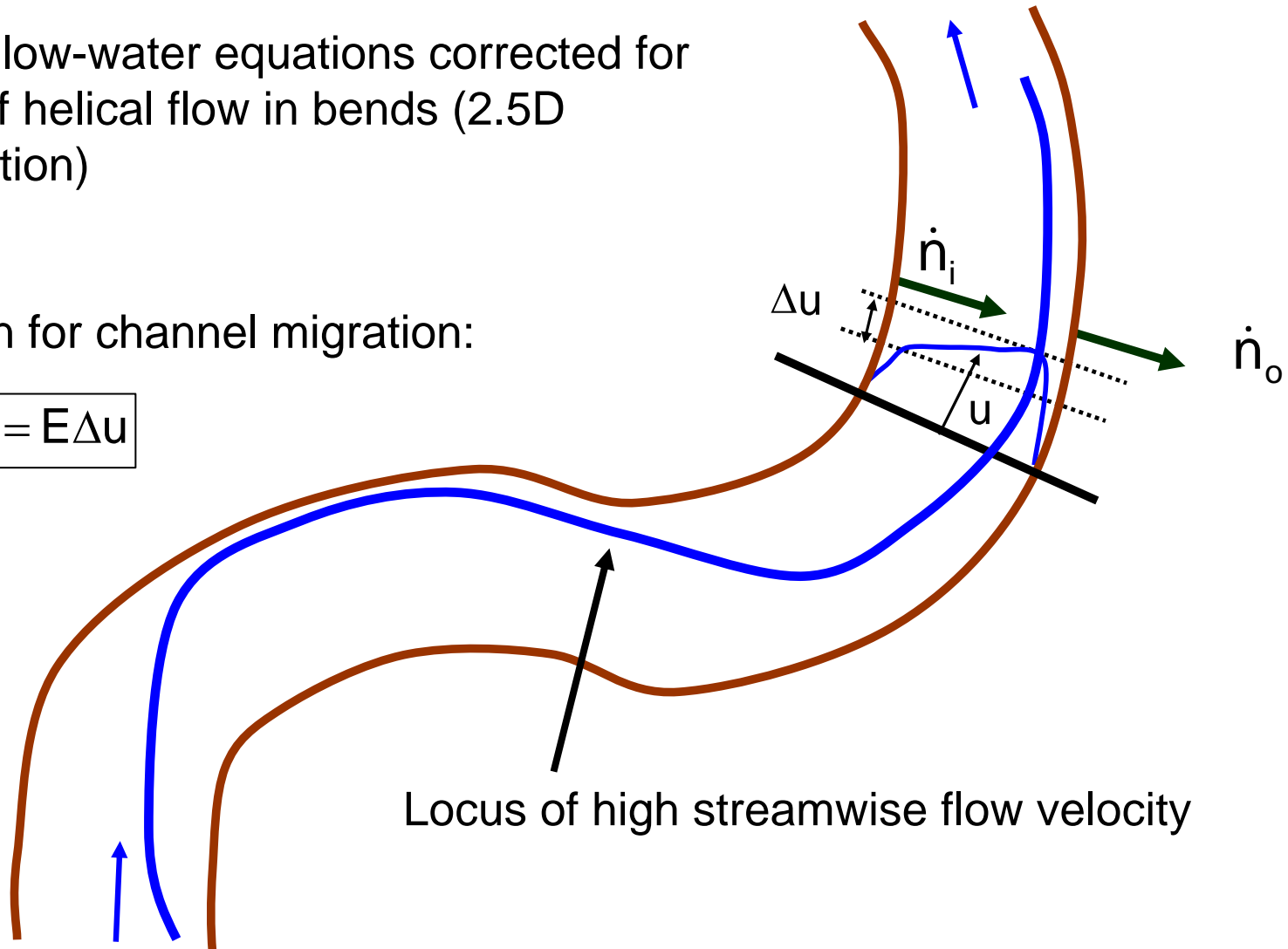
Occam's razor first analysis:

2D shallow-water equations corrected for
effect of helical flow in bends (2.5D
formulation)

+

Relation for channel migration:

$$\dot{n}_o = E\Delta u$$



ONLY BRAIDING IS POSSIBLE IN THE ABSENCE OF BANK STABILIZATION

Image courtesy B. Murray



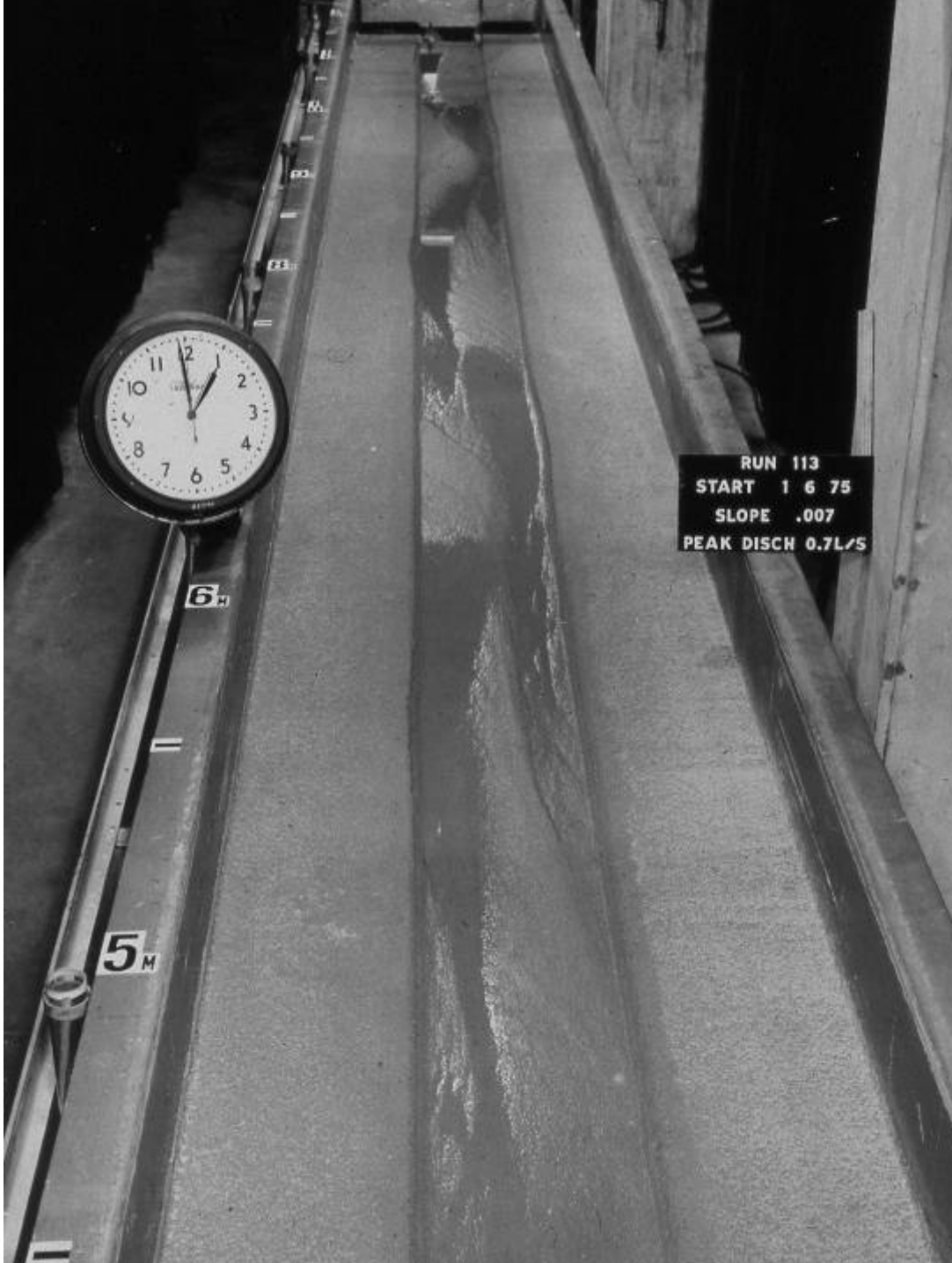
Braided stream on the North Slope, Brooks Range, Alaska

SWEET LITTLE LIES

Yes, honey, I'm with you.



<http://images.google.com/imgres?imgurl=http://jeroenarendsen.nl/pics/Conversation-Icelanders-pianist.gif&imgrefurl=http://jeroenarendsen.nl/category/countries/&h=327&w=387&sz=94&hl=en&start=14&um=1&tbnid=4OxIA6BGtduRcM:&tbnh=104&tbnw=123&prev=/images%3Fq%3Dconversation%26ndsp%3D21%26svnum%3D10%26um%3D1%26hl%3Den%26client%3Dfirefox-a%26rls%3Dorg.mozilla:en-US:official%26sa%3DN>



Inside bank to outside
bank:

“Yes, I’m following you.”



Inside bank to outside
bank:

“Yes, I’m following you.”

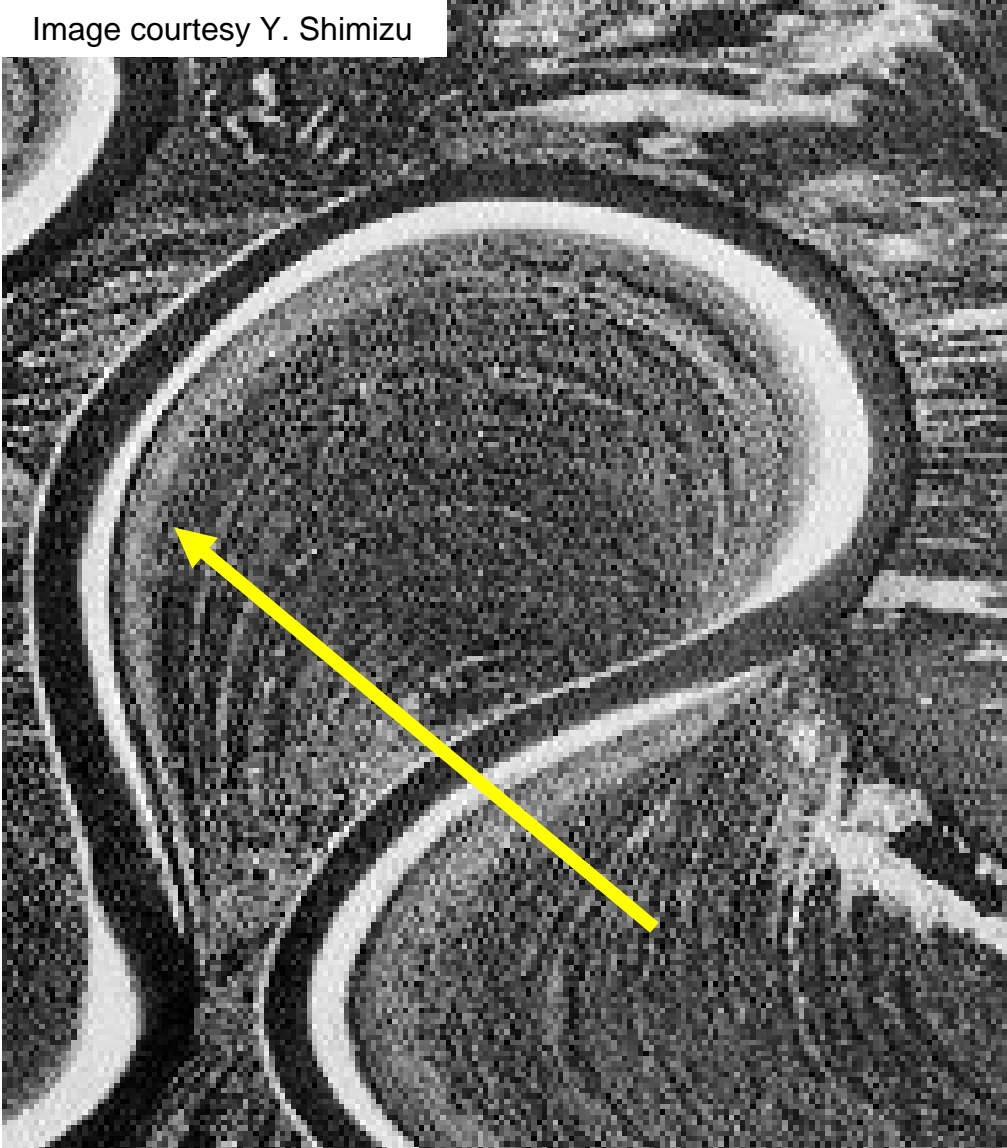
Yeah, while you were
pretending to listen,
***look at the mess we
got ourselves into!***



Image courtesy H. Johannesson

BEND SKEWING: SUBCRITICAL BIFURCATION OF NONLINEAR STABILITY ANALYSIS

Image courtesy Y. Shimizu



Bends grow until cutoff:
There is no nonlinear
stable state

Old oxbow lake due
to cutoff

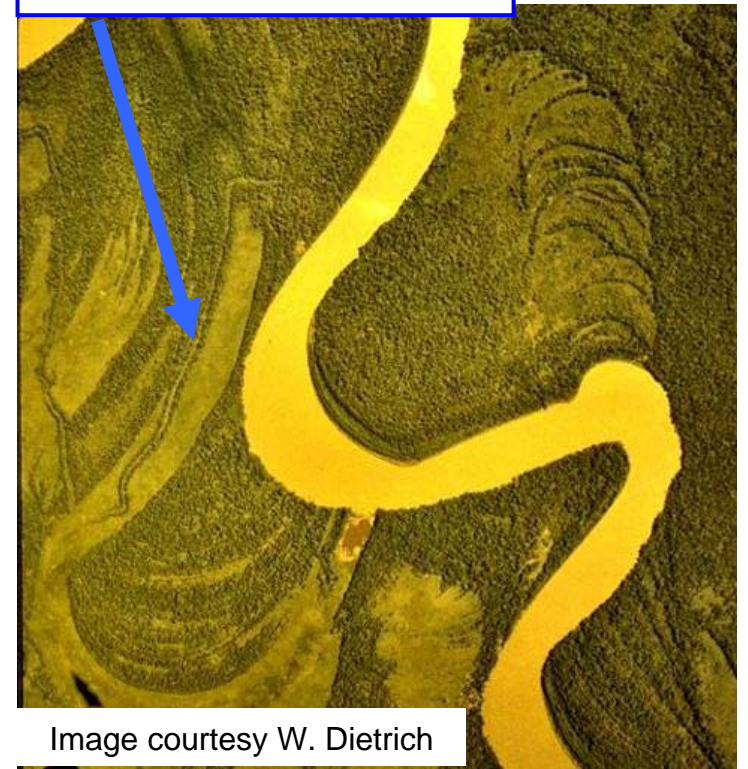


Image courtesy W. Dietrich

SLUMPING ON OUTSIDE SLOWS DOWN EROSION SO TRAPPING OF SEDIMENT BY VEGETATION ON INSIDE CAN KEEP UP



Vermilion River, USA



I slump

Do we talk to
each other?

I trap

Vermilion River, USA

NOW FOR A TOUR OF MORPHODYNAMIC PHENOMENA WITHOUT DETAIL AS TO HOW THEY ARE SOLVED

Yes, they are tractable to various degrees



SCROLL BARS

Image courtesy Y. Shimizu



Image courtesy W. Dietrich

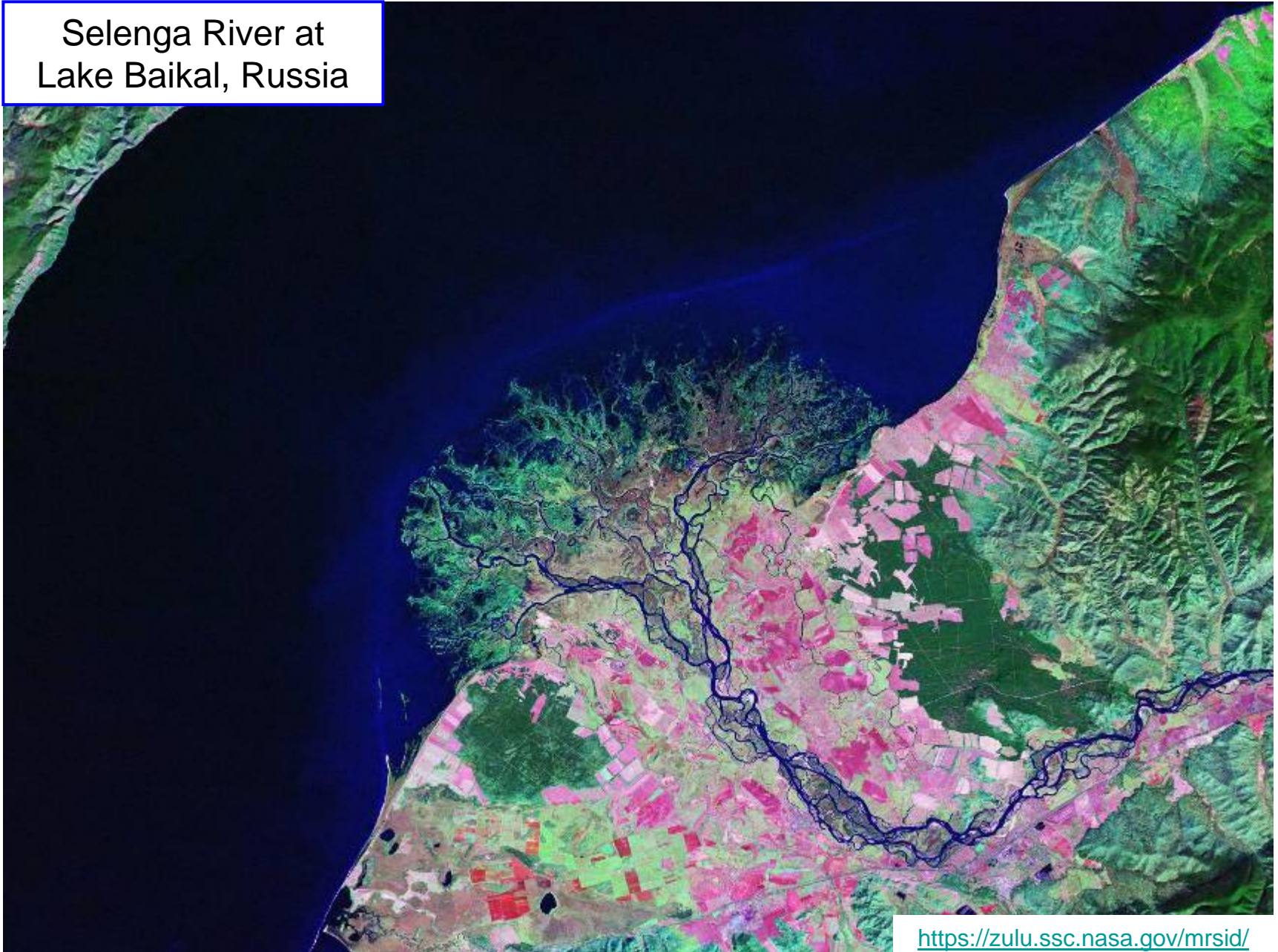


Strickland River, Papua
New Guinea

I still can't remember

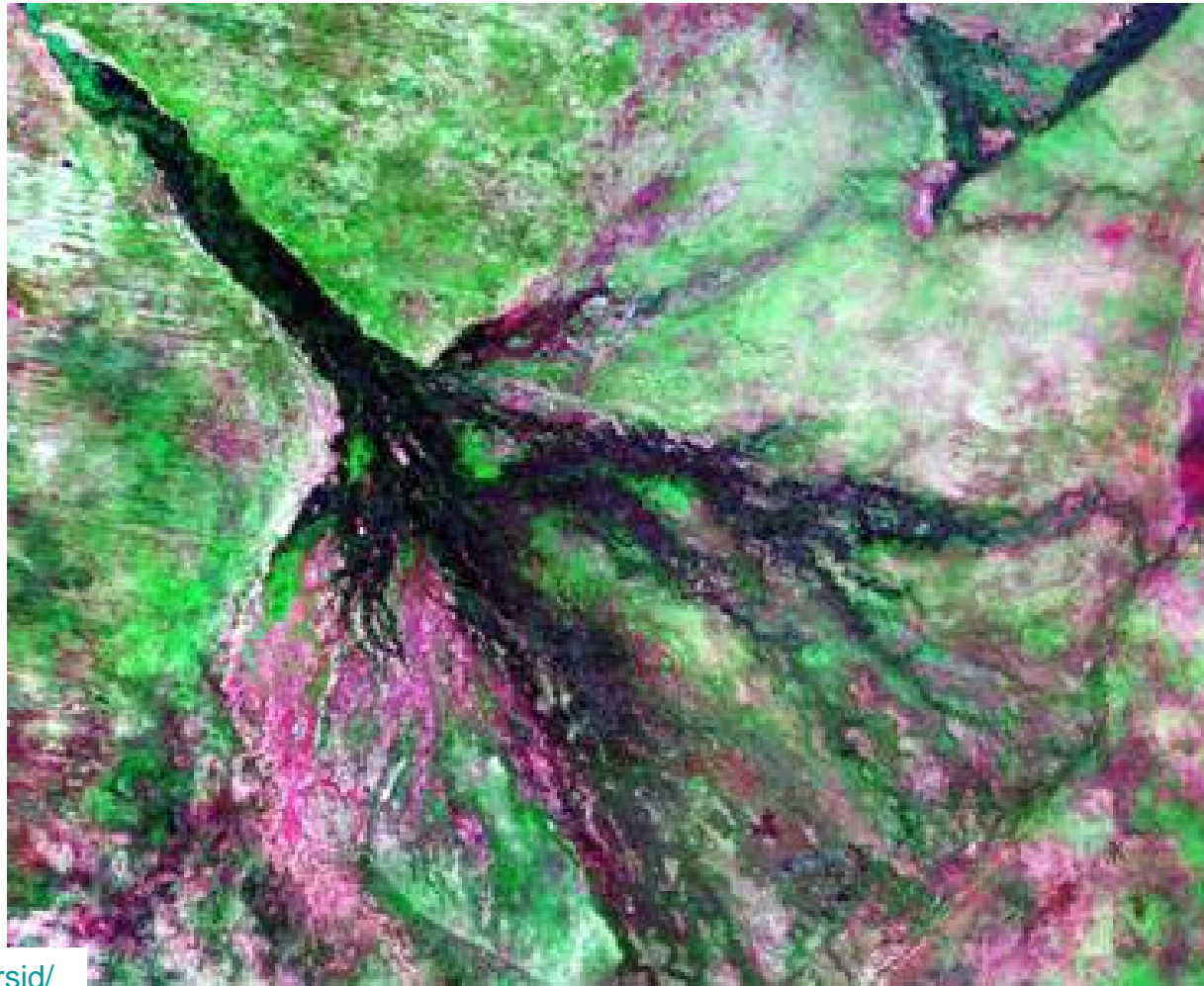
ALLUVIAL FANS AND FAN-DELTAS

Selenga River at
Lake Baikal, Russia

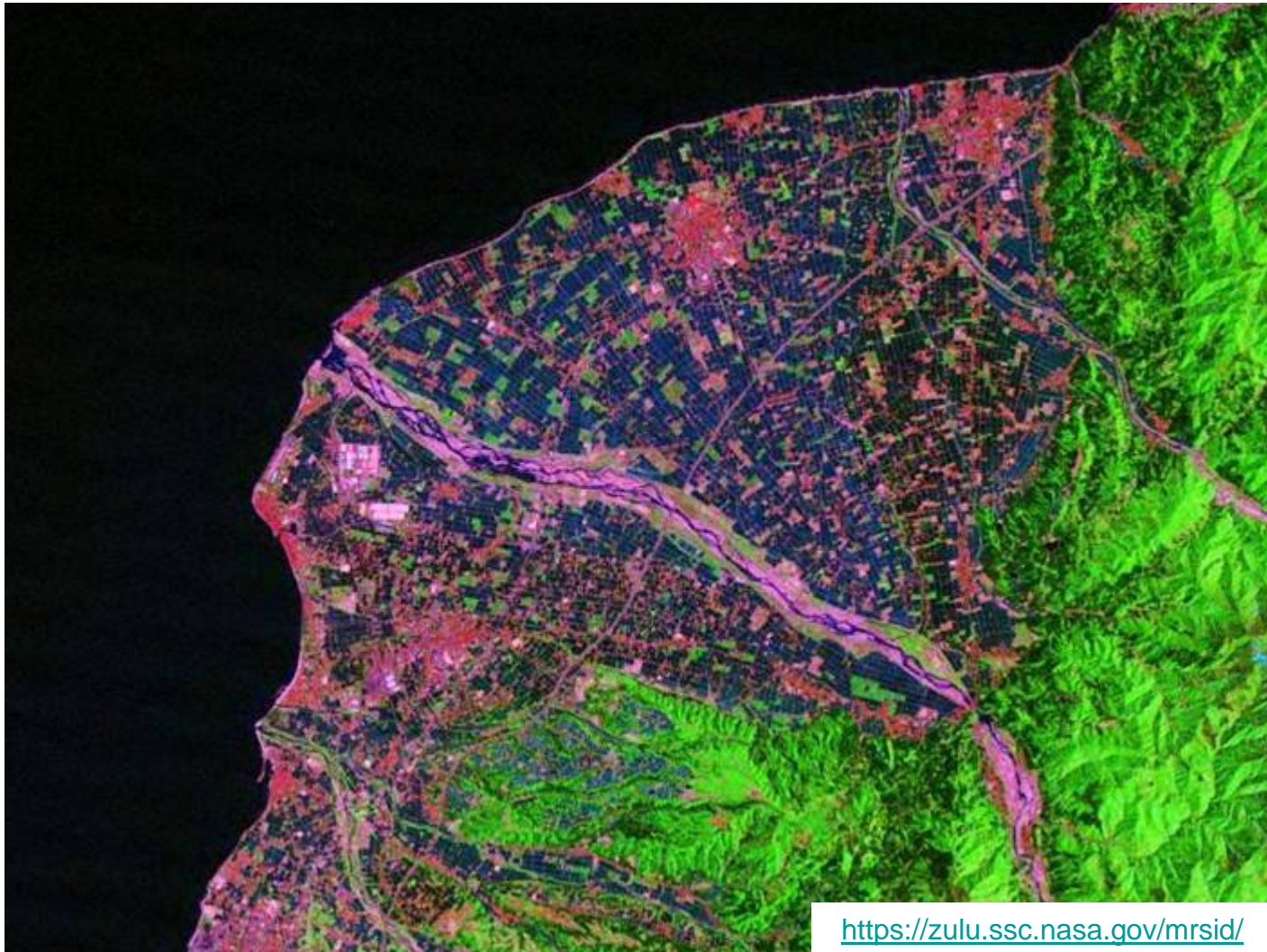


THE OKAVANGO INLAND FAN, BOTSWANA, AFRICA

Graben:
subsidence



THE FAN-DELTA OF THE KUROBE RIVER, JAPAN



<https://zulu.ssc.nasa.gov/mrsid/>

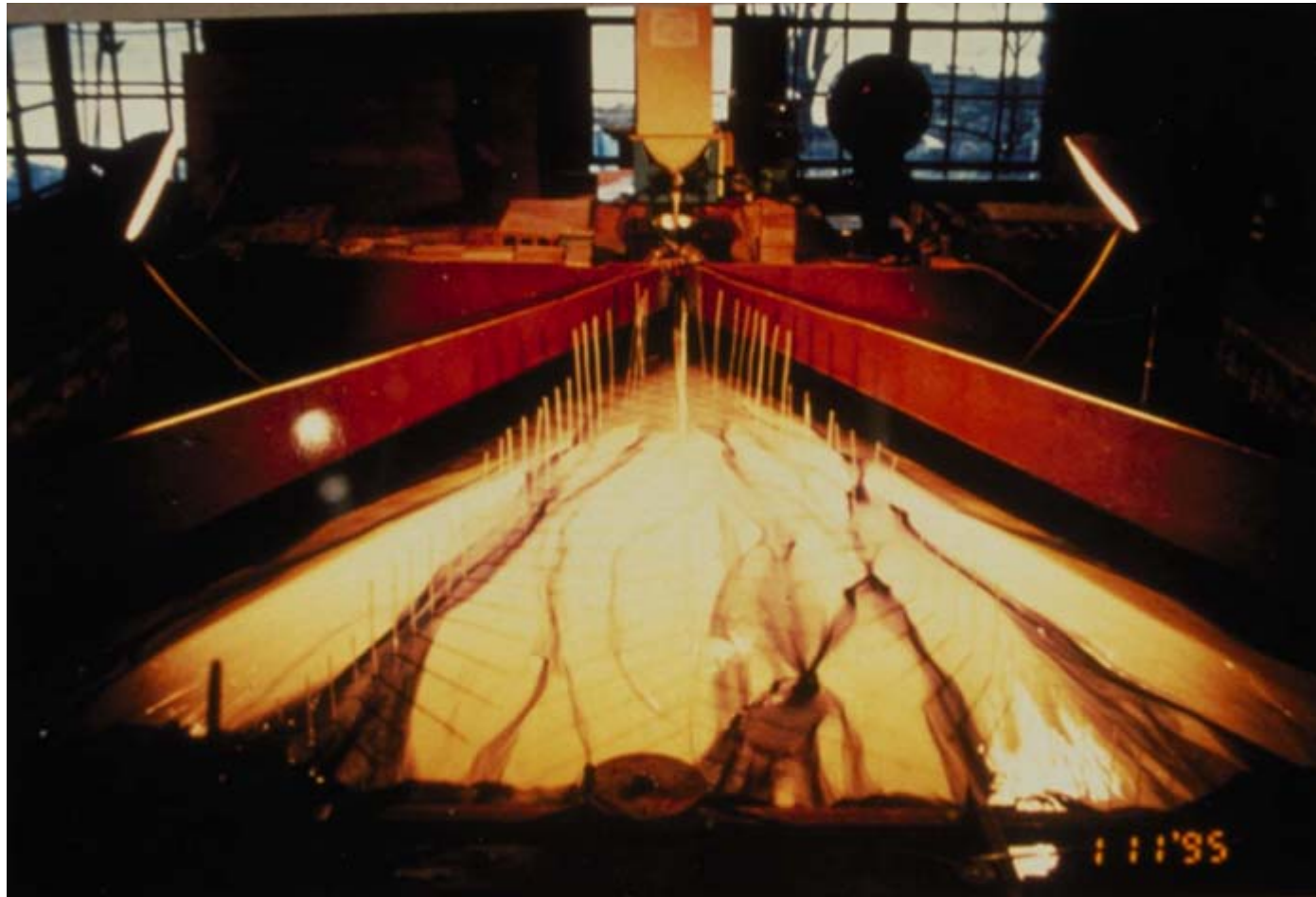
THE FAN-DELTA OF THE IOCC IRON MINE, LABRADOR, CANADA



THE FAN IN THE DELTA



FANS AND FAN-DELTAS AT VARIOUS SCALES



Laboratory fan-delta, ~ 3 m.

Image taken at St. Anthony Falls Laboratory, University of Minnesota USA.

FANS AND FAN-DELTA AT VARIOUS SCALES contd.



Fan created by runoff from cultivated field; ~ 6 m.
Image taken by author near Pigeon Point, California.

FANS AND FAN-DELTA AT VARIOUS SCALES contd.



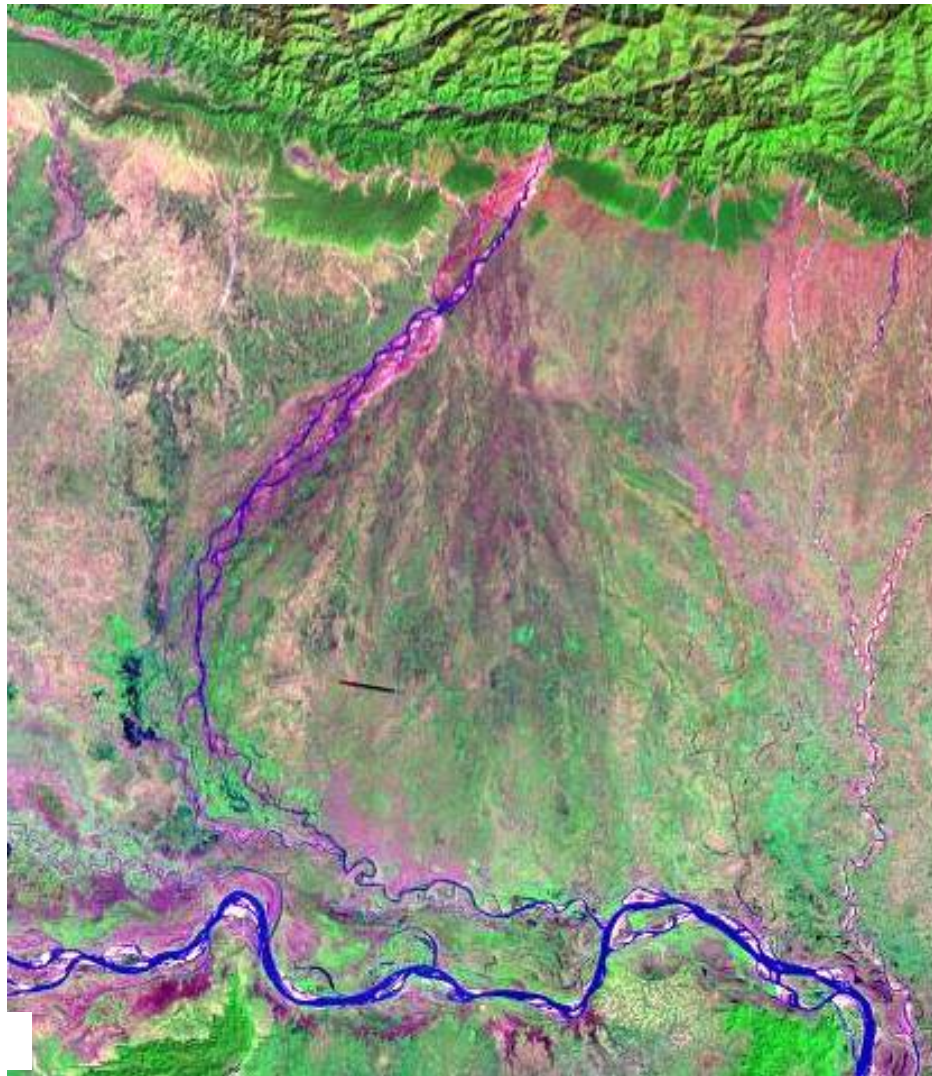
Fan in Idaho, USA created by runoff from burned hillside, ~ 50 m.

FANS AND FAN-DELTA AT VARIOUS SCALES contd.



Copper Creek Fan, Death Valley, USA; ~ 10 km.
Image courtesy Roger Hooke.

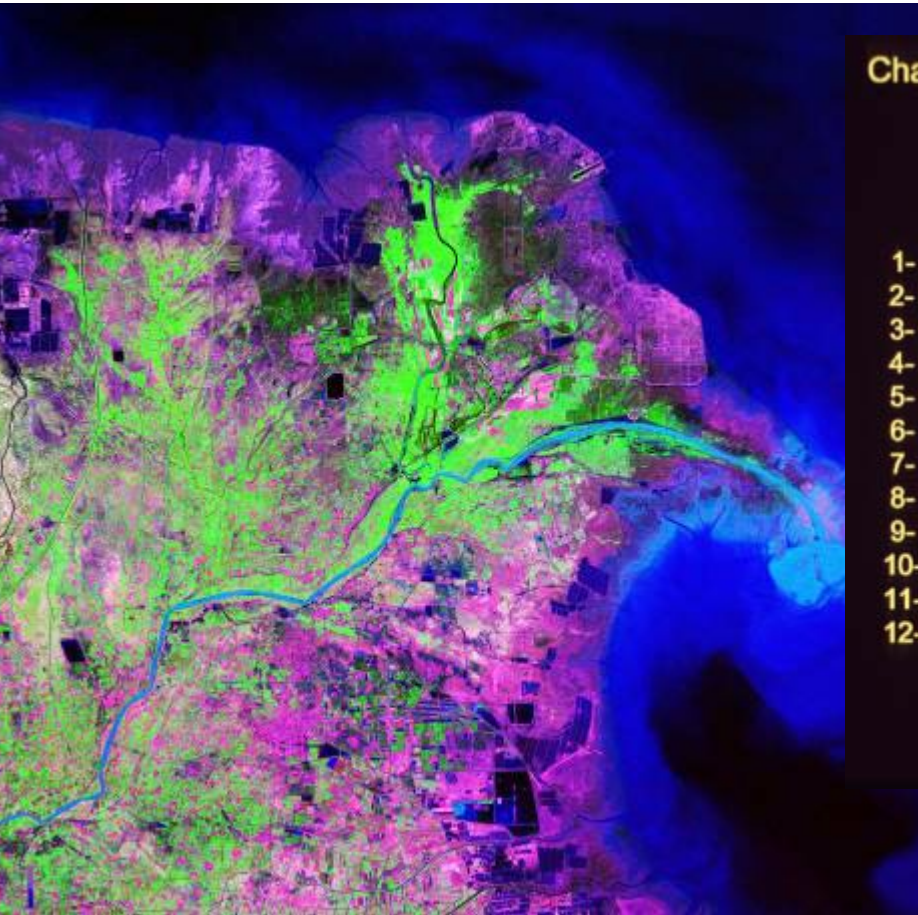
FANS AND FAN-DELTA AT VARIOUS SCALES contd.



<https://zulu.ssc.nasa.gov/mrsid/>

Kosi River Fan, India; ~ 125 km.

RIVER MIGRATION AND AVULSION MAKES FANS



Changes of River Course of Yellow River Fan-Delta in Recent History

- 1- 1855.6-1889.3
- 2- 1889.3-1897.5
- 3- 1897.5-1904.6
- 4- 1904.6-1917.7
- 5- 1917.7-1926.6
- 6- 1926.6-1929.8
- 7- 1926.6-1934.8
- 8- 1934.8-1953.7
- 9- 1953.7-1960.8
- 10- 1960.8-1964.1
- 11- 1964.1-1976.5
- 12-1976.6

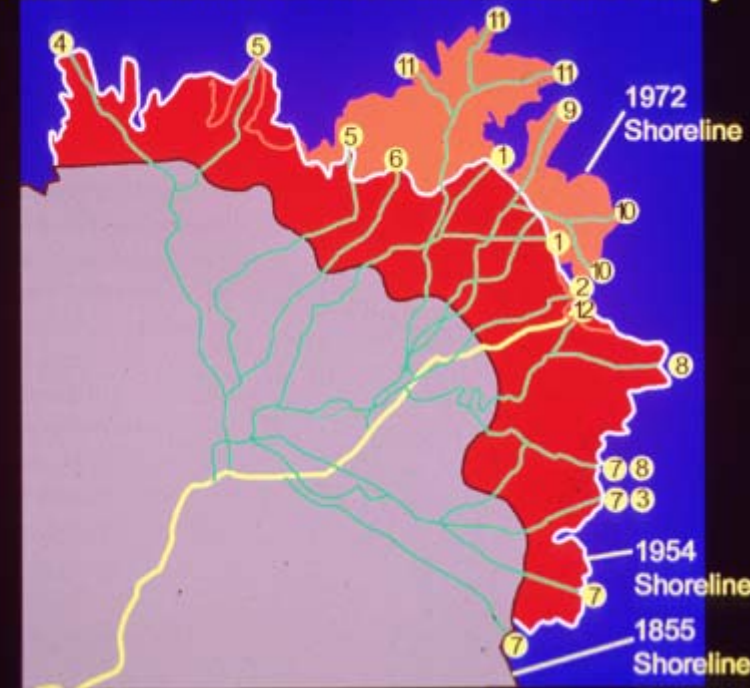
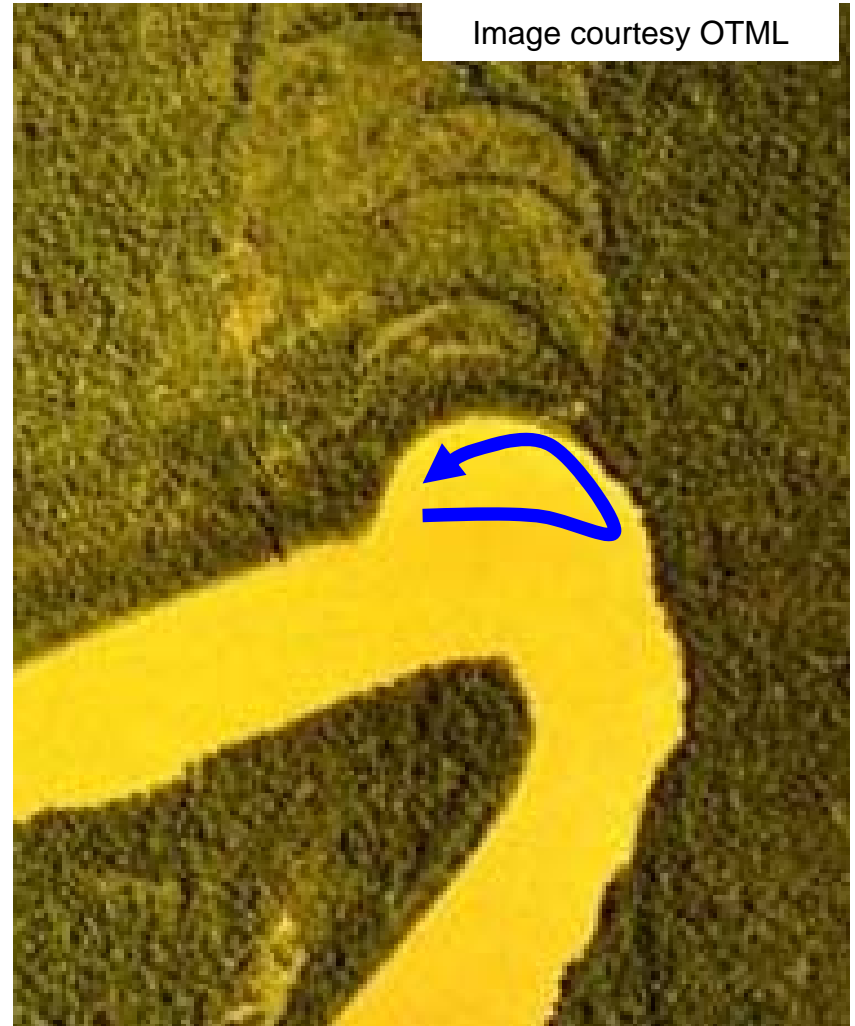


Image courtesy OTML



CONCAVE BANK BENCHES

Image courtesy OTML



SELF-CHANNELIZATION: NATURAL LEVEES

Image courtesy National Geographic



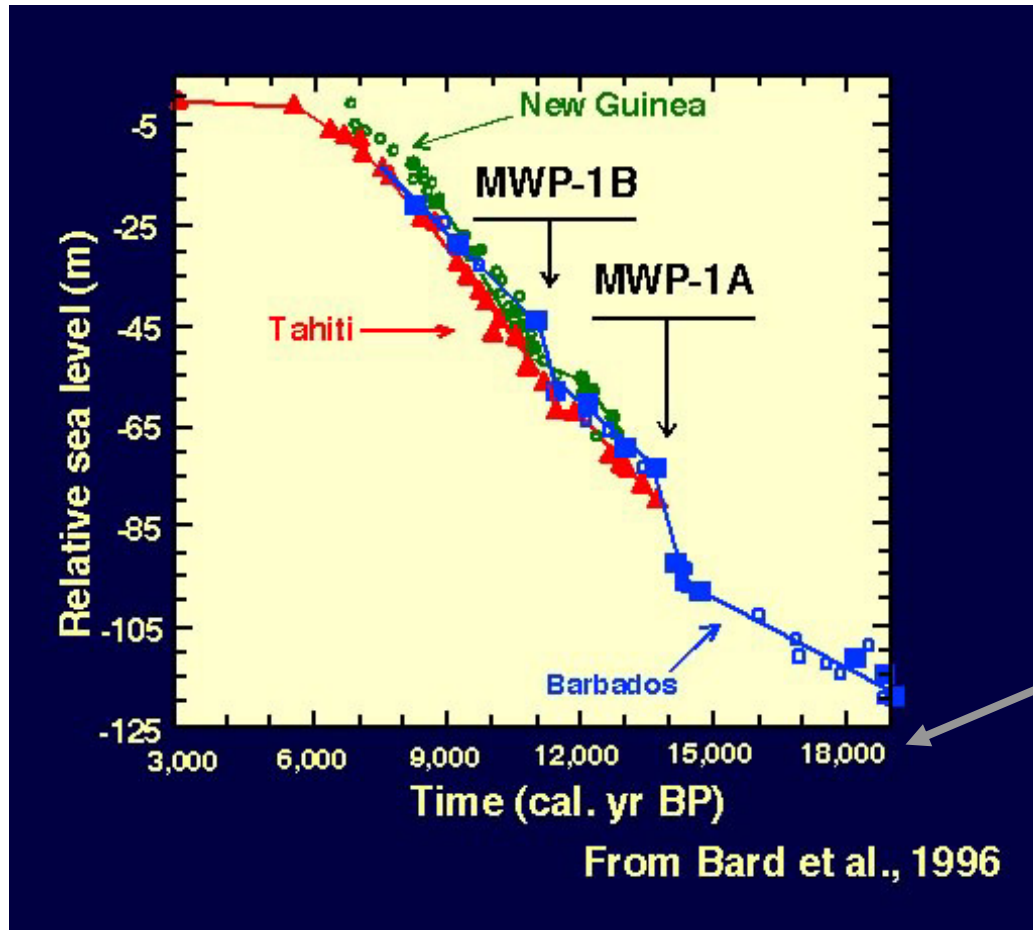
Gilgal Abey River,
Ethiopia



Mississippi River, USA



SEA LEVEL ROSE SOME 120 M SINCE THE END OF THE LAST GLACIATION



Years before present

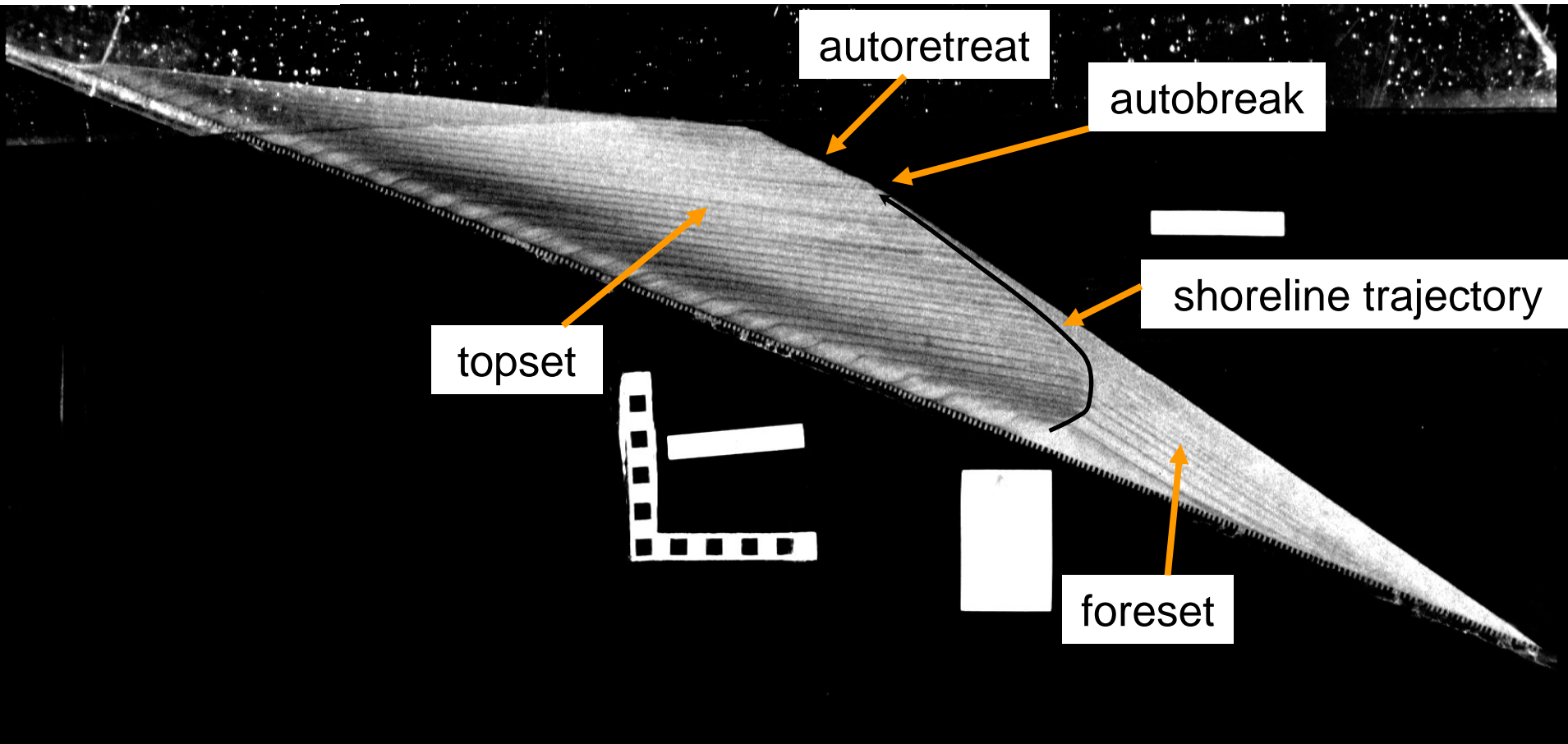
How does a river mouth respond to sea level rise?

- Does a delta continue to prograde into the ocean?
- Or does the sea drown the delta and invade the river valley (transgression)?

DELTA AND SEA LEVEL RISE

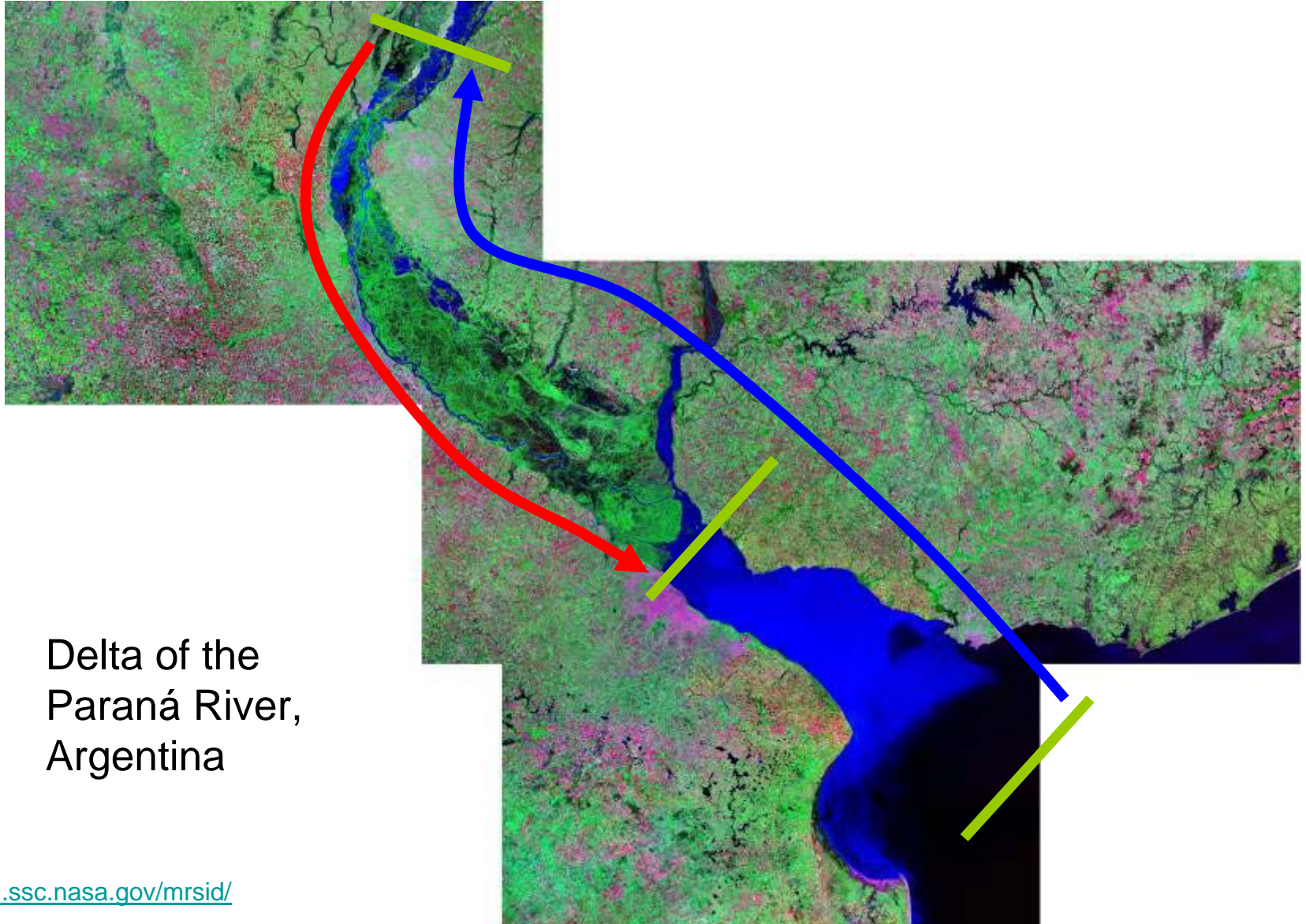
Experiment on effect of base level rise on delta

Image courtesy T. Muto



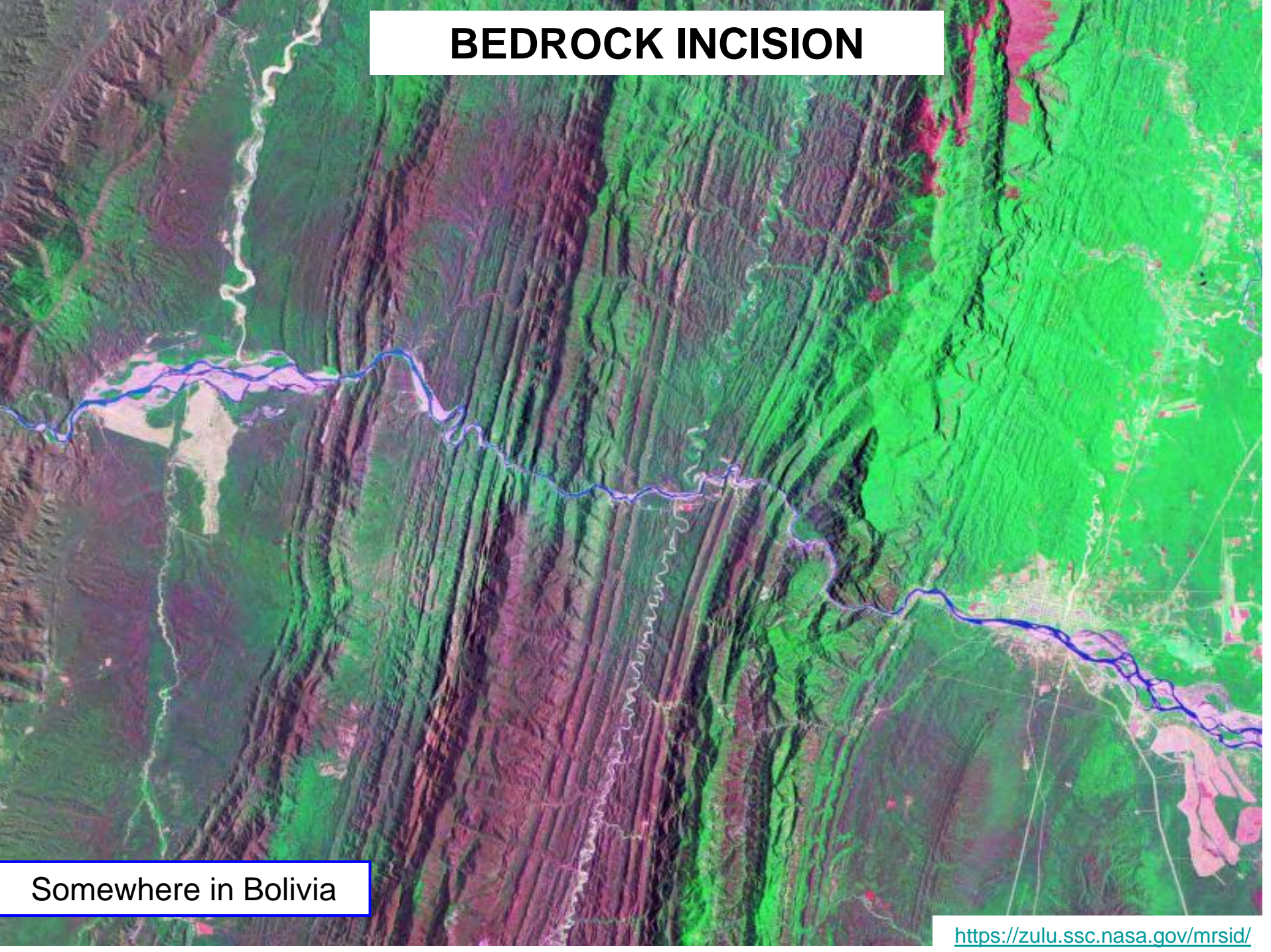
AUTORETREAT

HOW DID THE DELTAS OF MAJOR RIVERS RESPOND?



Delta of the
Paraná River,
Argentina

BEDROCK INCISION

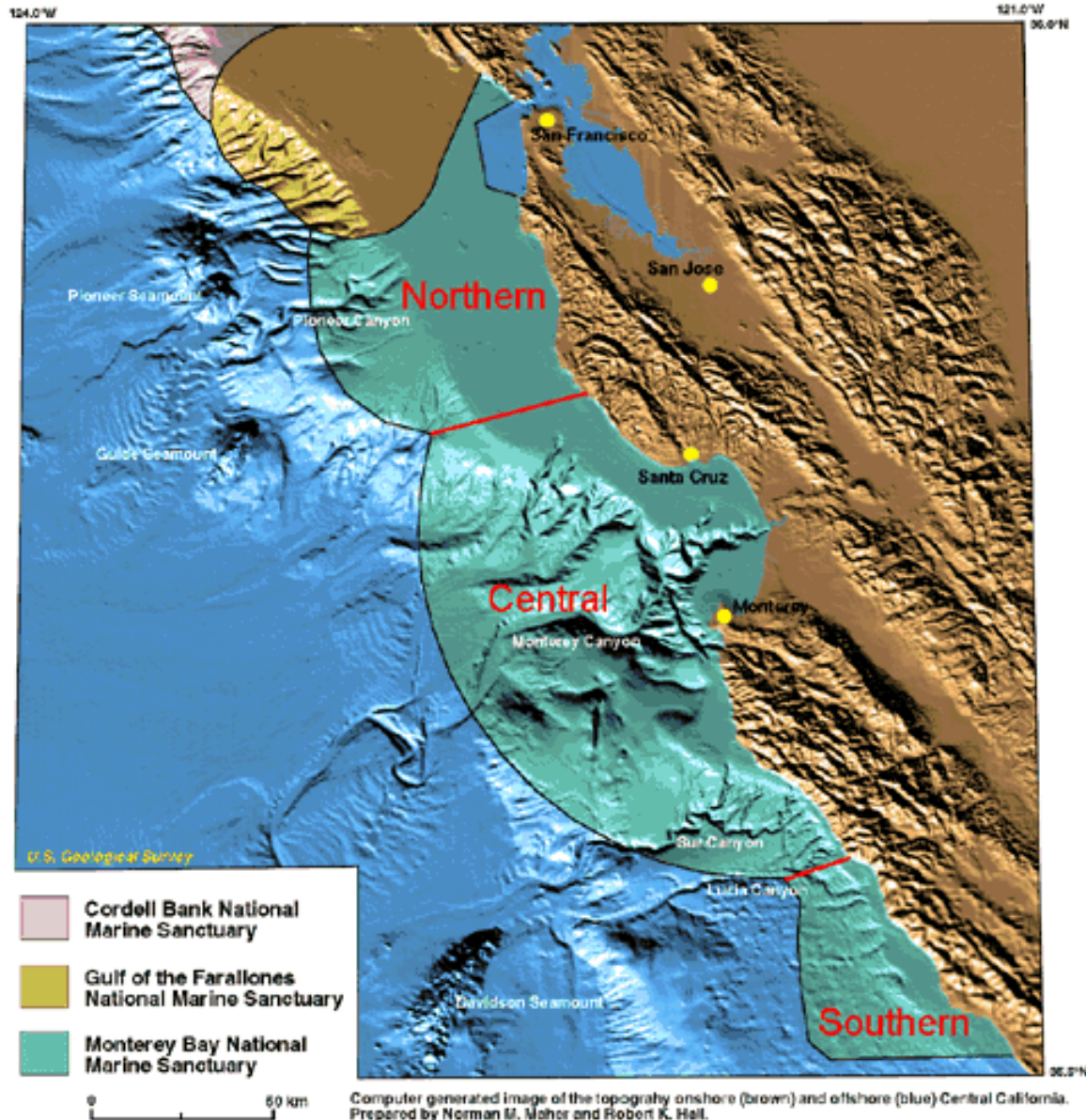


Somewhere in Bolivia



Hirose River, Japan

SUBMARINE MORPHODYNAMICS DUE TO TURBIDITY CURRENTS



California Margin

Image courtesy MBARI



CANYON EXCAVATION

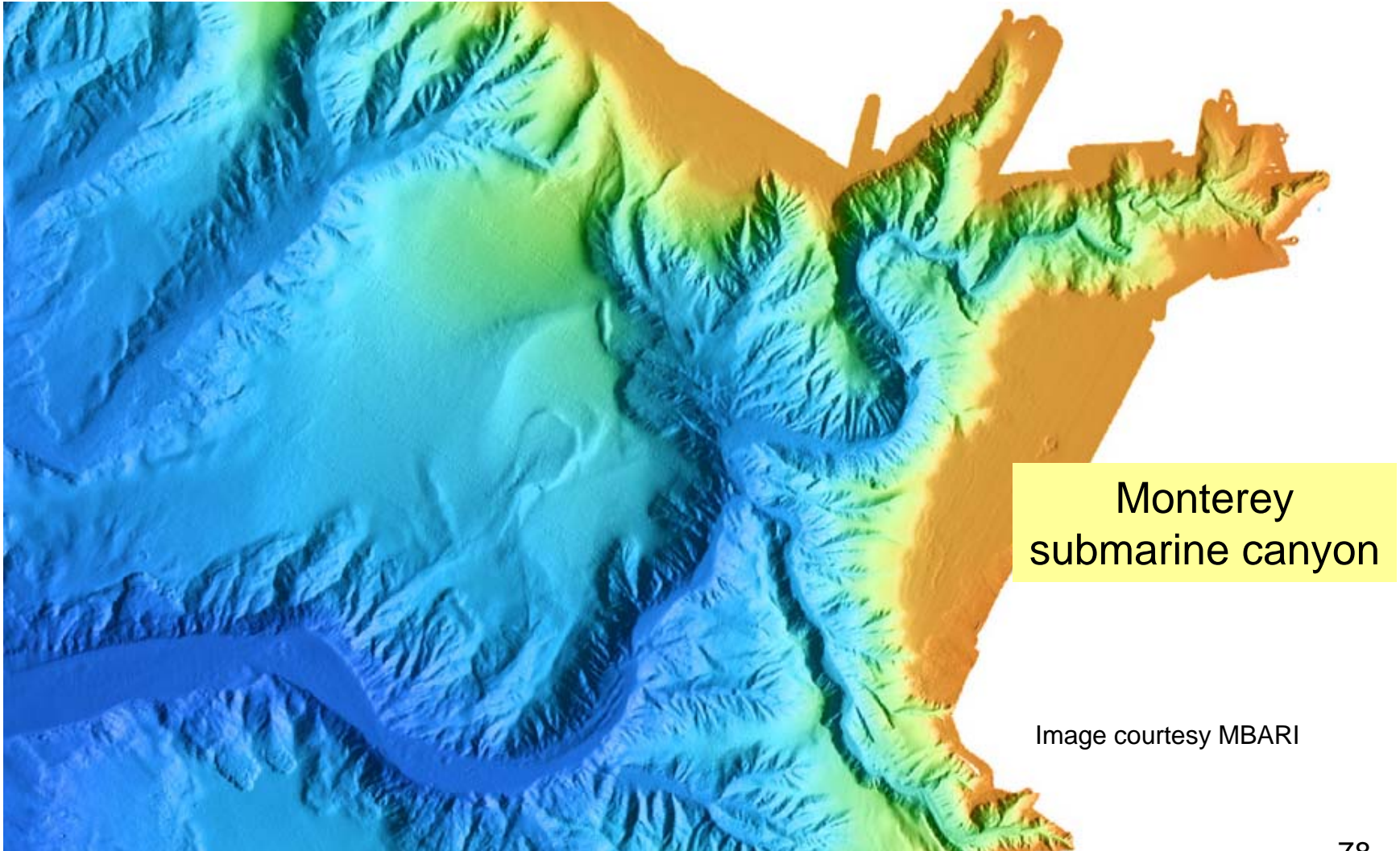
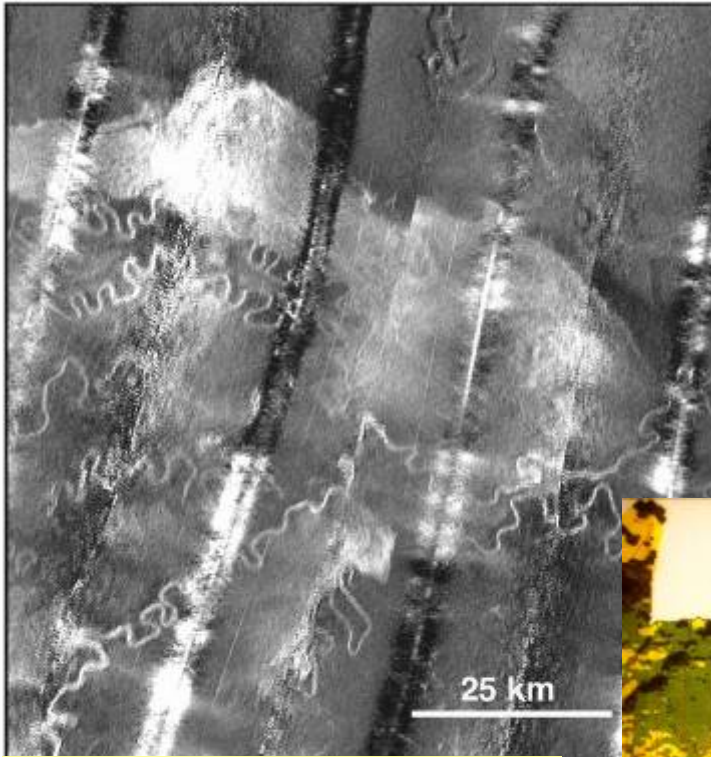
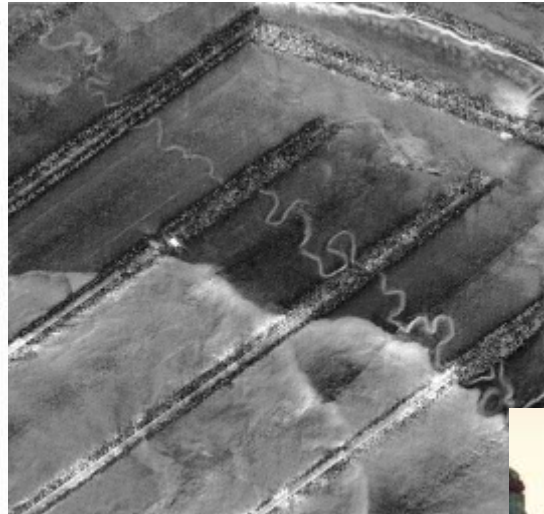


Image courtesy MBARI

MEANDERING OF SUBMARINE CHANNELS



Indus Submarine Fan
(Kenyon et al., 1995)

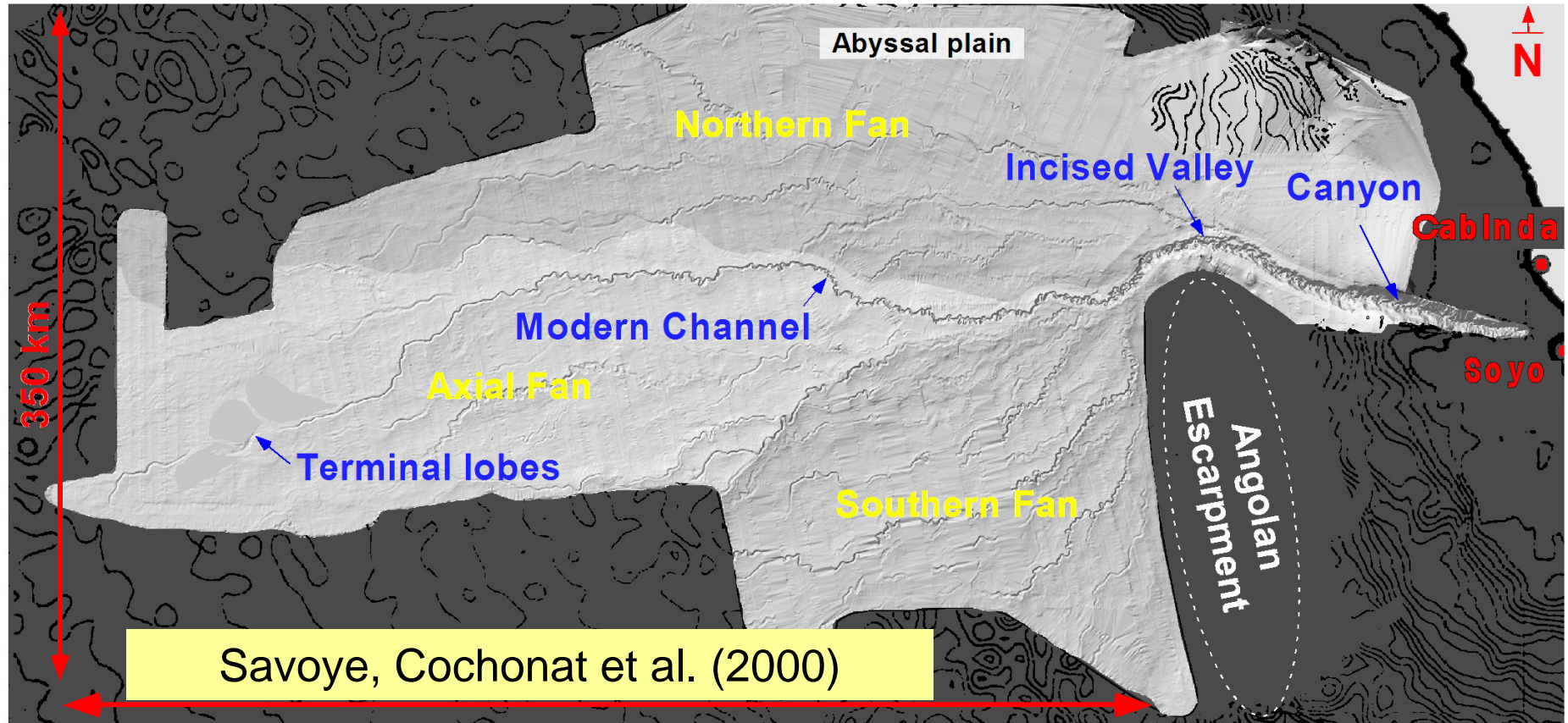


Mississippi
Submarine Fan
(Weimer, 1991).



Amazon Submarine
Fan (Pirmez, 1995)

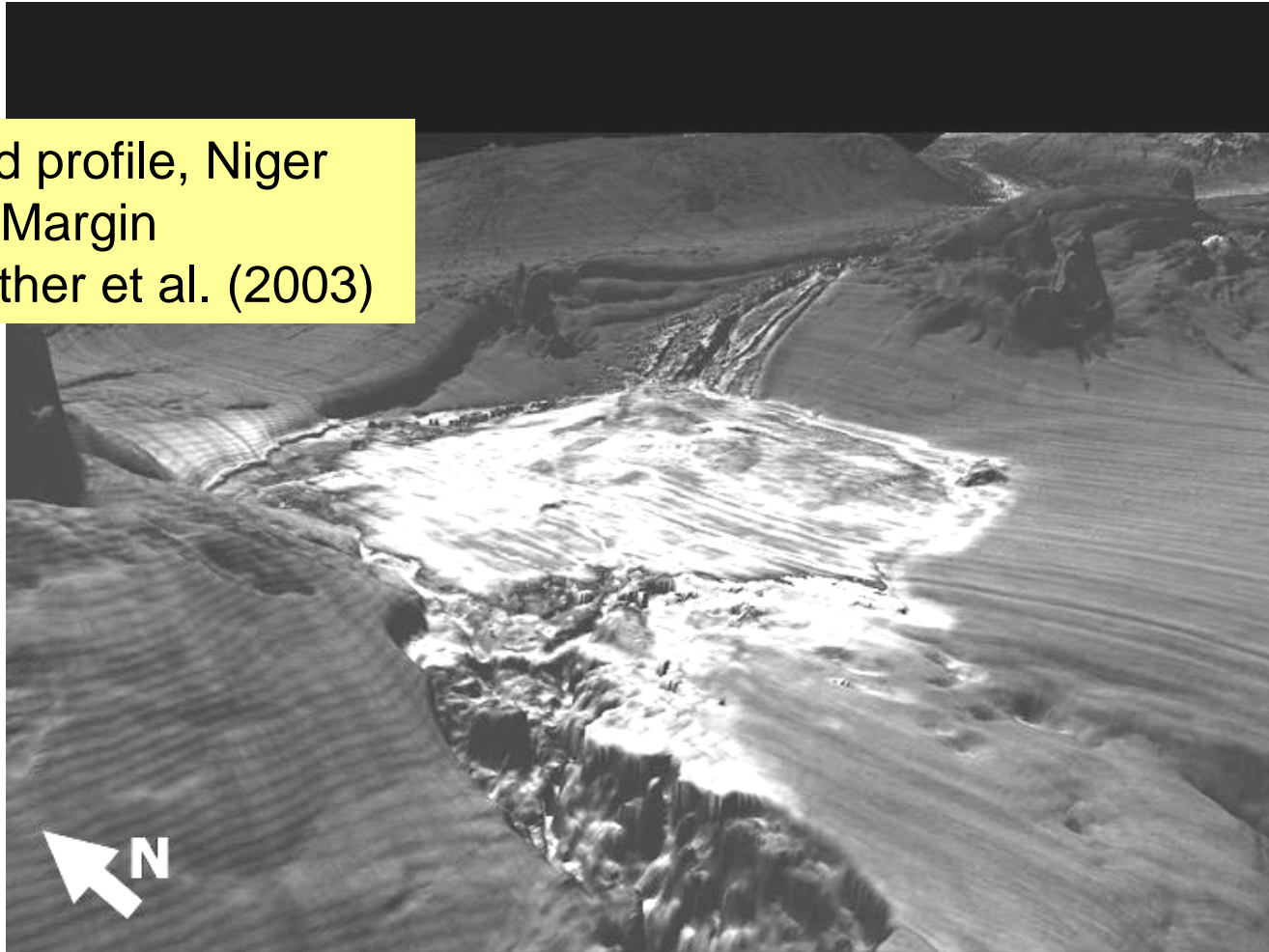
CONGO DEEP-SEA FAN



EM12 Bathymetry of the Pleistocene to Present Zaire Deep-Sea Fan.

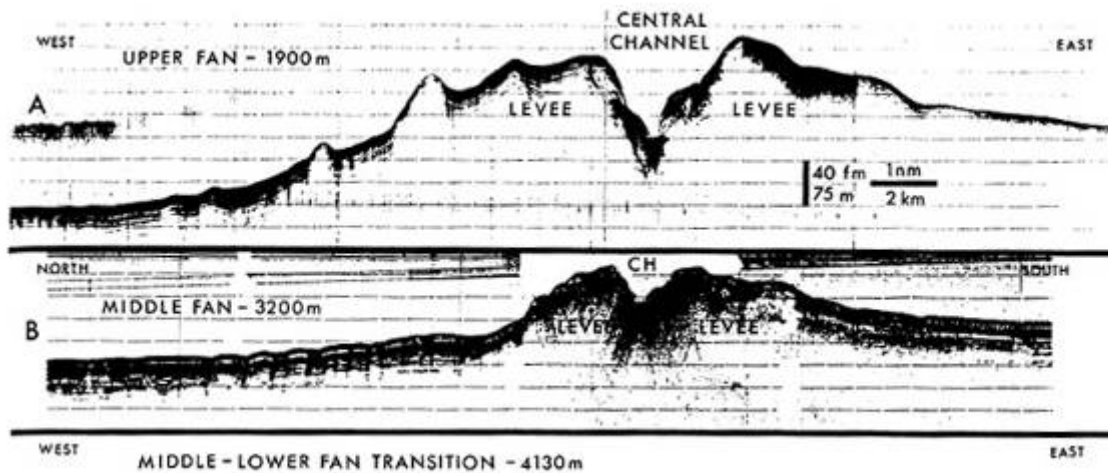
FANS AND CANYONS: STEPPED PROFILES

Stepped profile, Niger
Margin
From Prather et al. (2003)

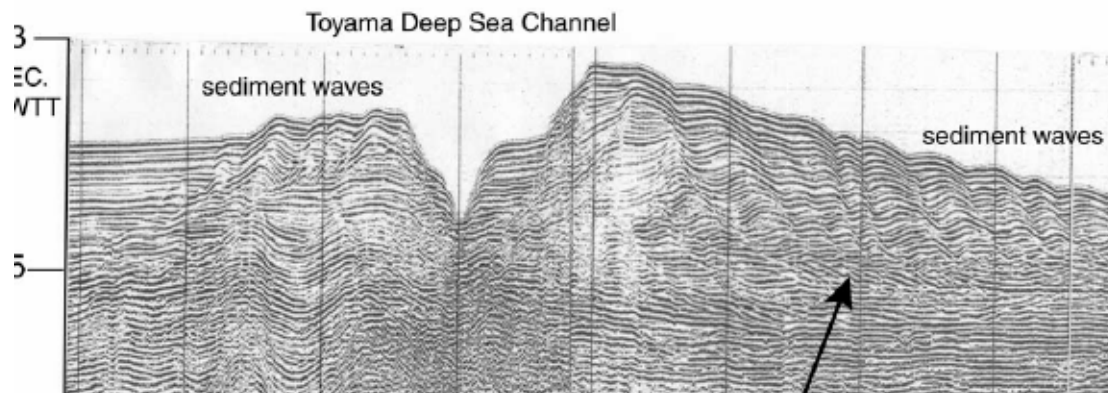


SELF-CONFINEMENT AND LEVEE CONSTRUCTION

Turbidity currents are adept at confining themselves between levees.



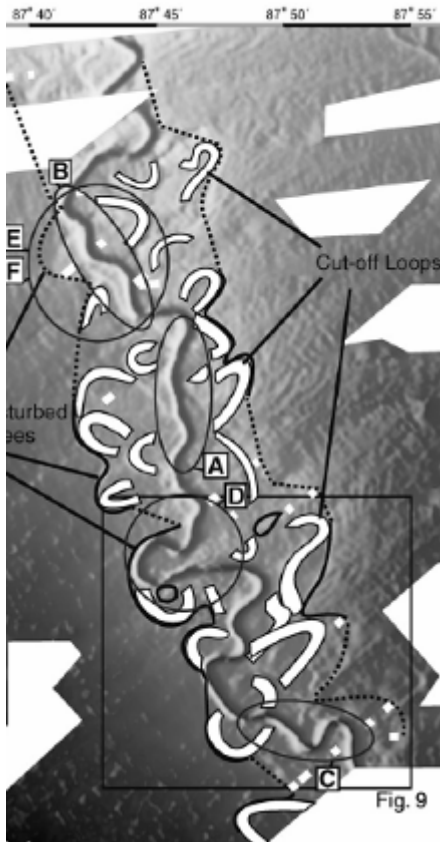
Channel on Amazon
Submarine Fan
Damuth and Flood (1985)



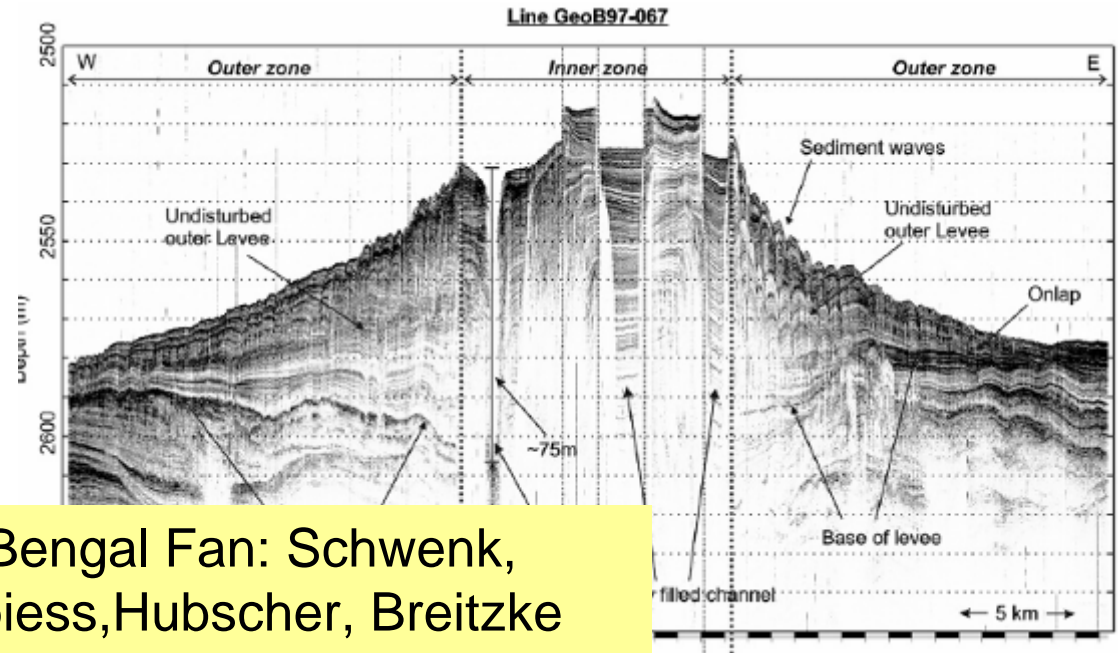
Toyama Submarine Channel
Kubo and Nakajima (2002)

SELF-CONTAINMENT

Submarine meandering channels contain themselves between levees over 100's ~ 1000's of km and scores ~ 100's of bends.



Bengal Fan: Schwenk, Spiess, Hubscher, Breitzke (2003)



Zaire Fan: Savoye, Cochonat et al. (2000)

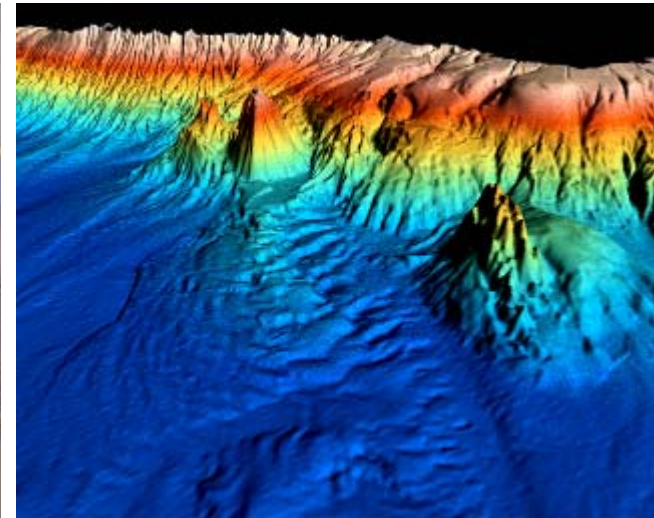
CYCLIC STEPS: **A UNIVERSAL BEDFORM OF** **FROUDE-SUPERCritical FLOW** **IN RIVERS AND TURBIDITY CURRENTS FLOWING** **OVER ERODIBLE BEDS**



Dry Meadow Creek, USA

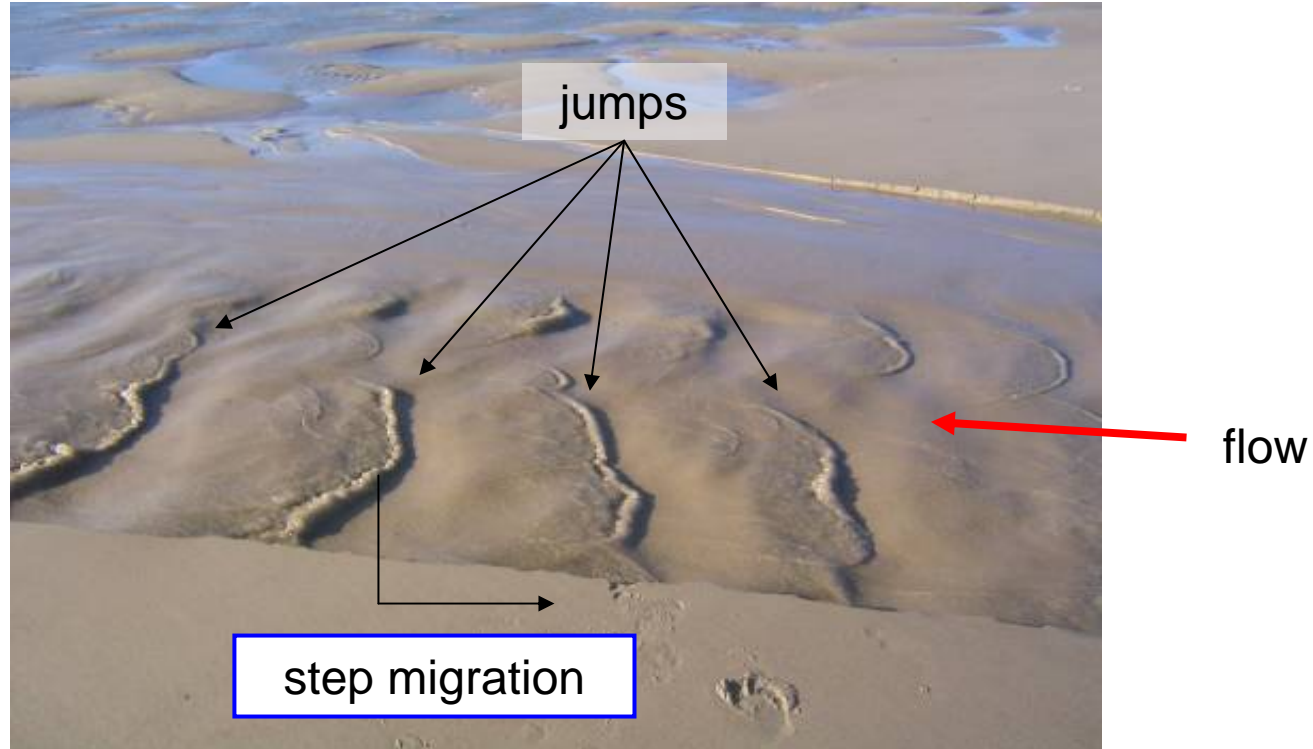


Small stream near
Calais, France



Deep sea offshore of
California, USA

TRAIN OF HYDRAULIC JUMPS



Trains of cyclic steps in a coastal outflow channel on a beach in Calais, France.
Image courtesy H. Capart.

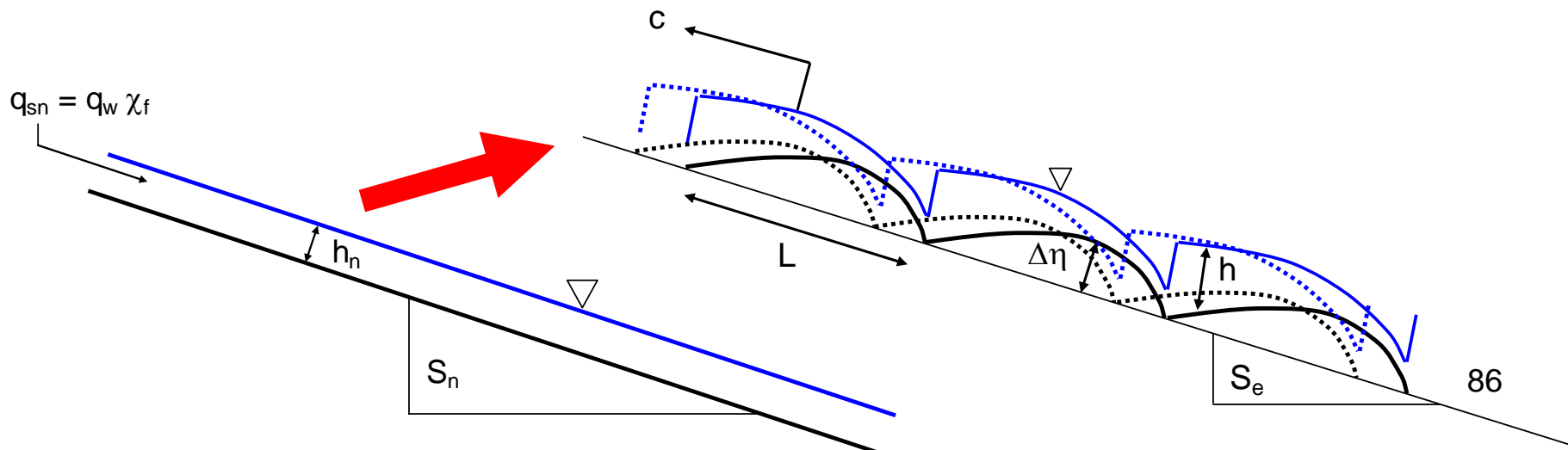
The steps move upstream

THE IDEA

Steady, uniform (normal) Froude-supercritical flow ($Fr_n > 1$) over a freely-erodible bed of sand might be unstable,

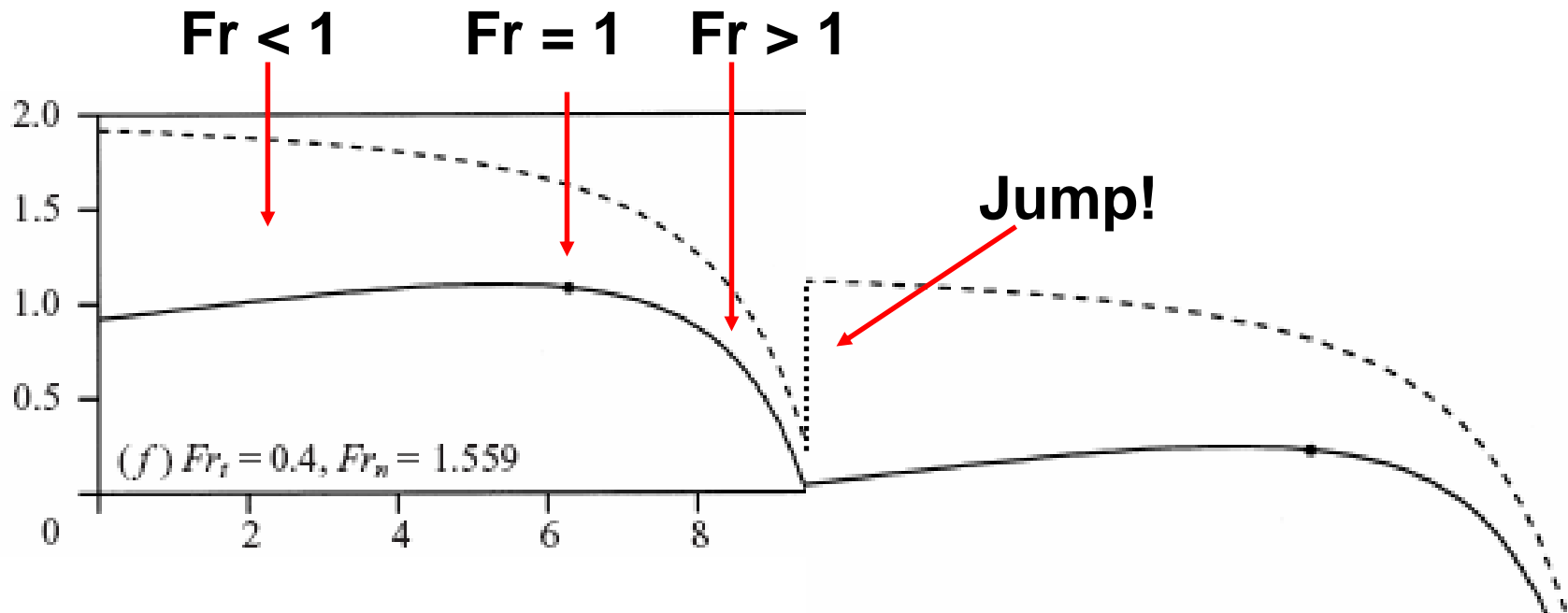
and within an appropriate range might not devolve to ephemeral, short-wave ($L/h \sim 1$) antidunes, but instead would devolve to

orderly, sustained trains of long-wave ($L/h \ll 1$) cyclic steps, with regions of subcritical and supercritical flow bounded by hydraulic jumps

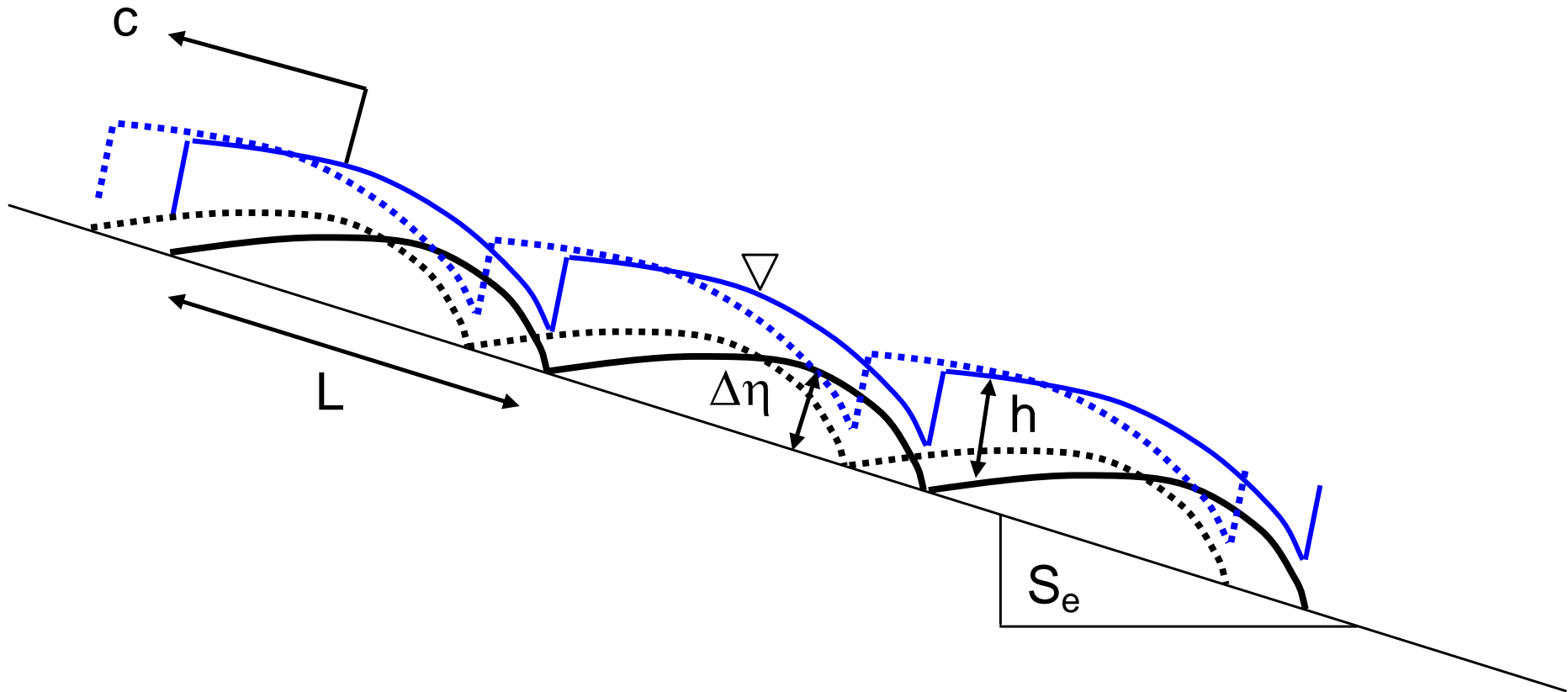


FULLY NONLINEAR PERIODIC SOLUTION OF PERMANENT FORM WITH CONSTANT UPSTREAM MIGRATION

Sufficiently supercritical flow over a plane bed is subject to a *long-wave* instability that devolves into upstream-migrating supercritical and subcritical regions bounded by hydraulic jumps.

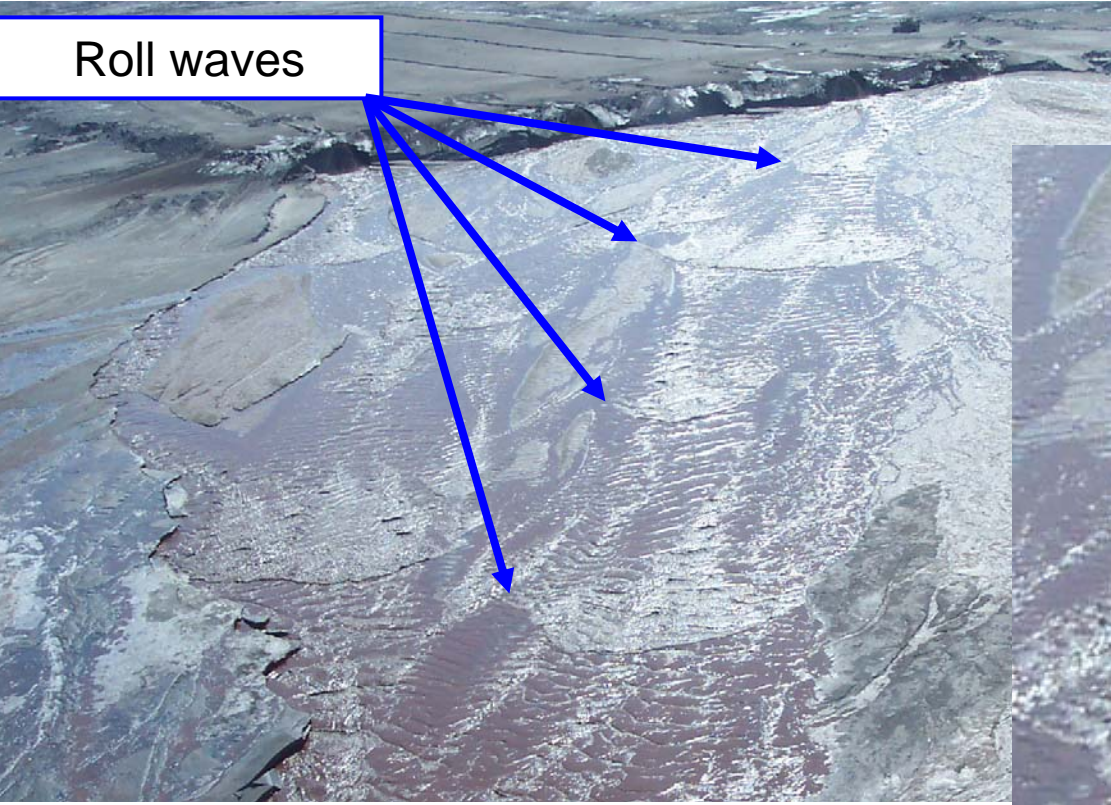


LIKE THIS



LET'S LOOK AT THIS IMAGE AGAIN

Roll waves



Cyclic steps!



Tailings Fan, Lake
Wabush, Canada

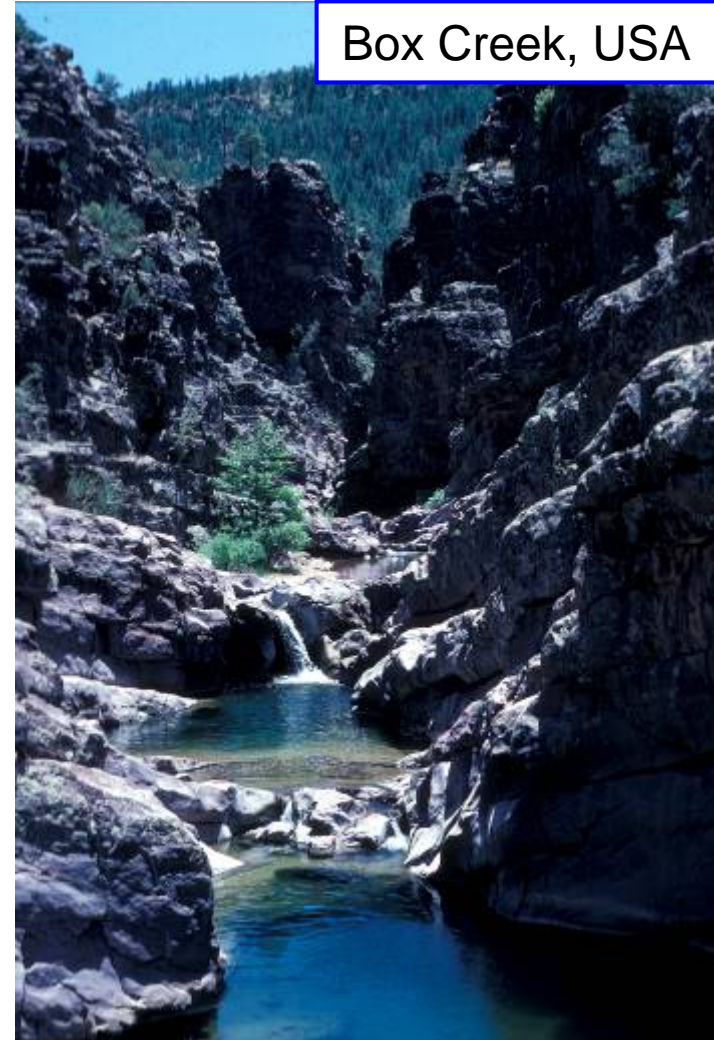
THE SAME CYCLIC STEP INSTABILITY IS FOUND IN INCISING BEDROCK STREAMS



Dry Meadow Creek, USA



Gough Island

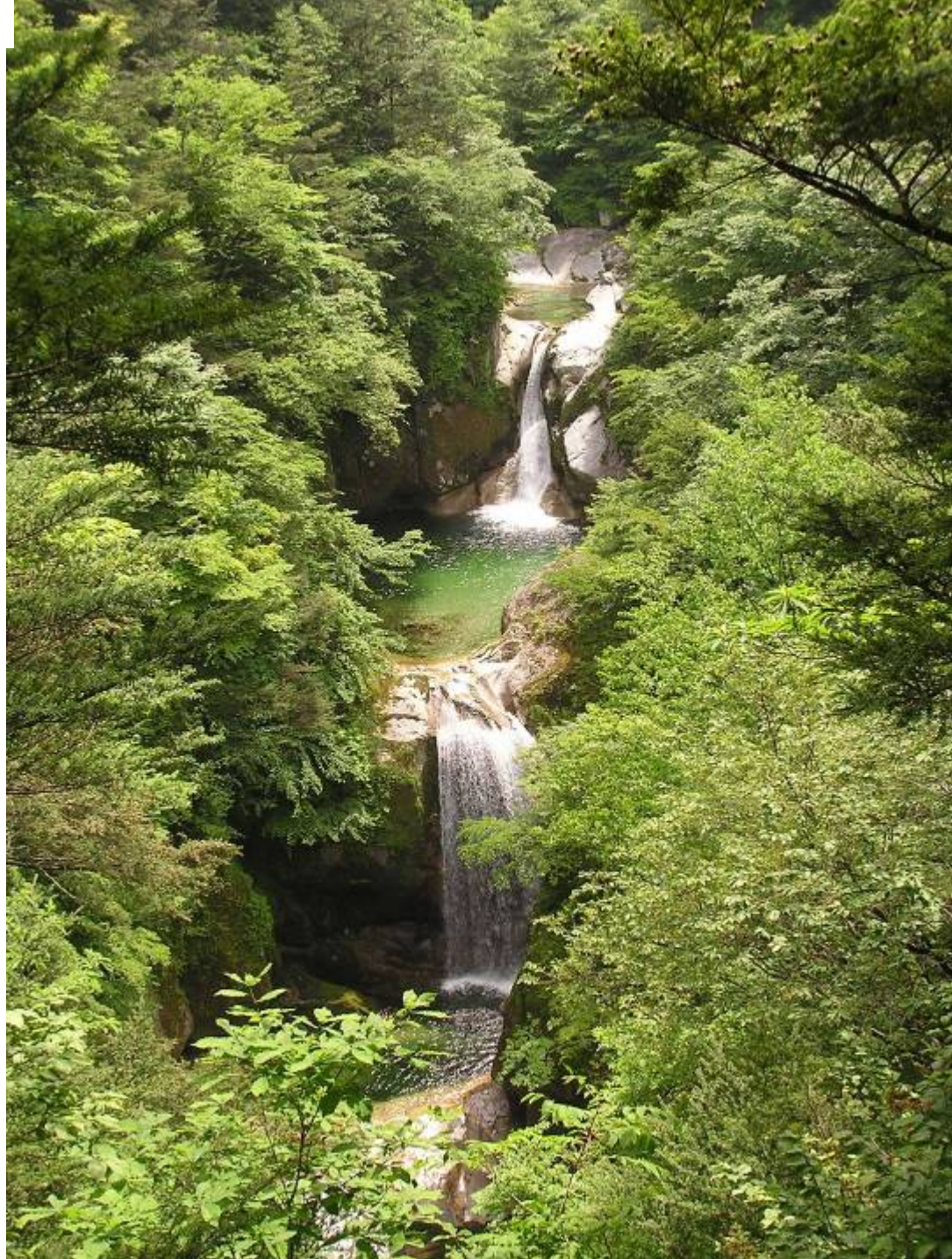


Box Creek, USA

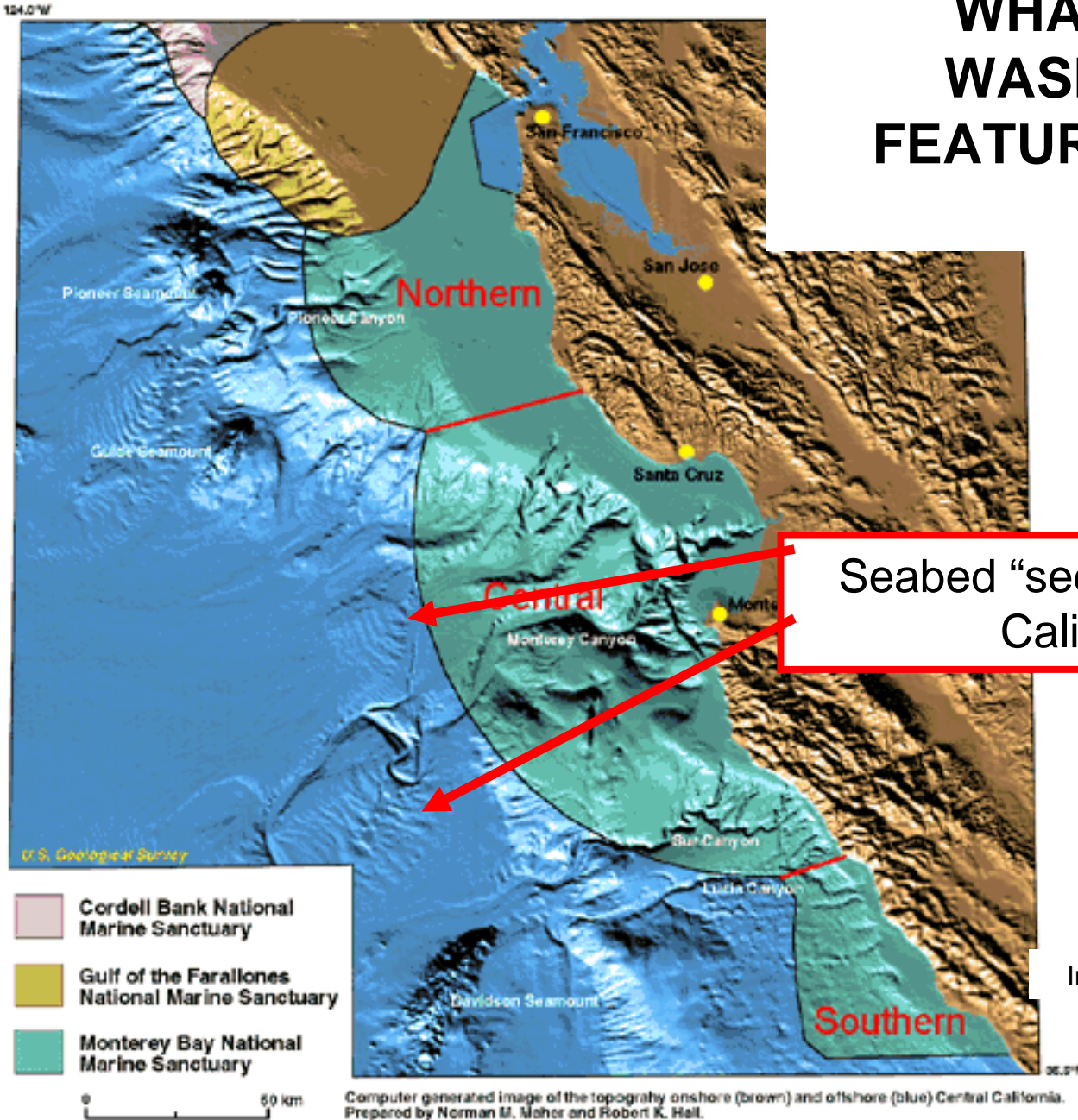
Images courtesy Michael
Neumann, Gough Island Weather
Station and Ellen Wohl

**CYCLIC STEPS IN
BEDROCK:**
*This one is too
beautiful not to show*

Ojiro River, Japan



WHAT ARE THESE WASHBOARD-LIKE FEATURES IN THE DEEP SEA?



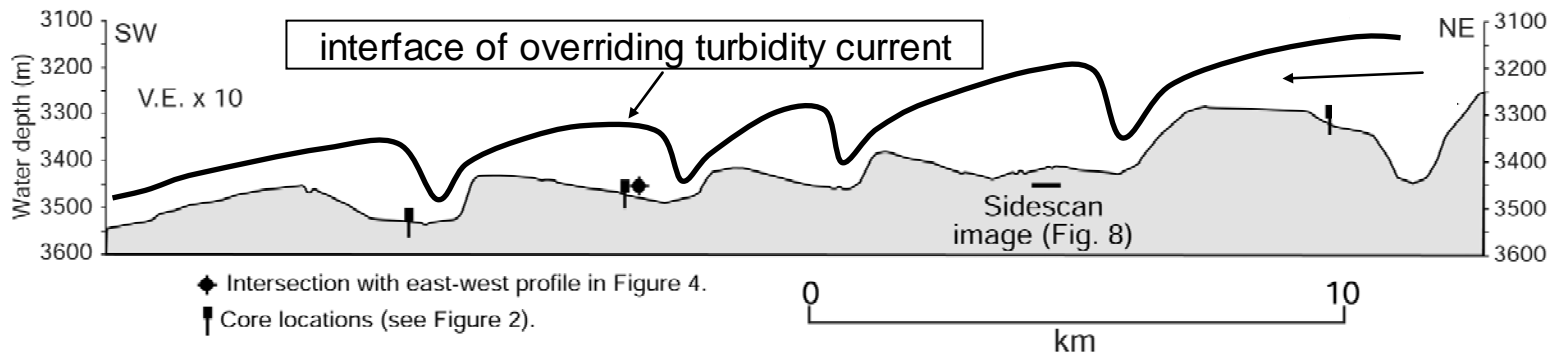
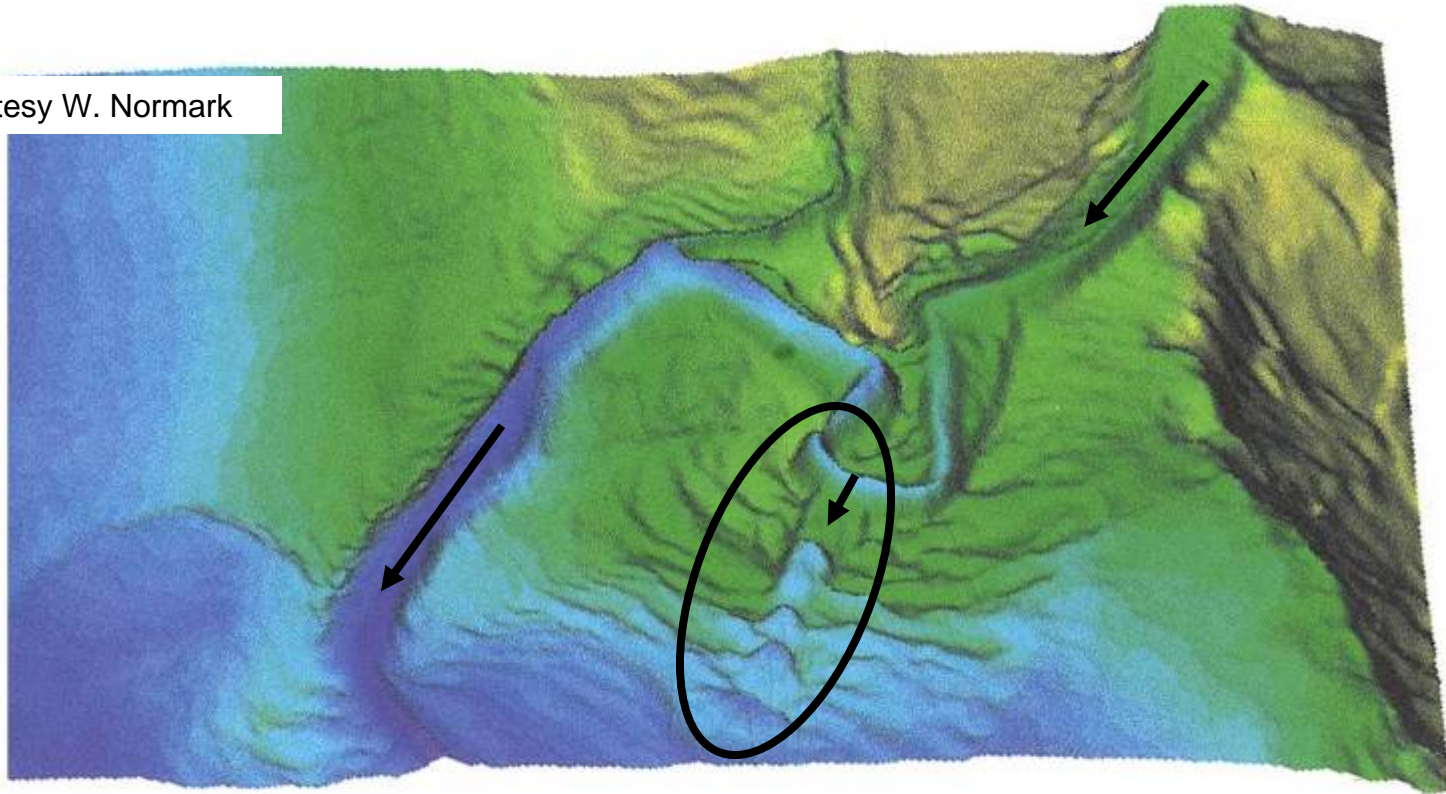
Seabed “sediment waves” off the California margin

Image courtesy MBARI

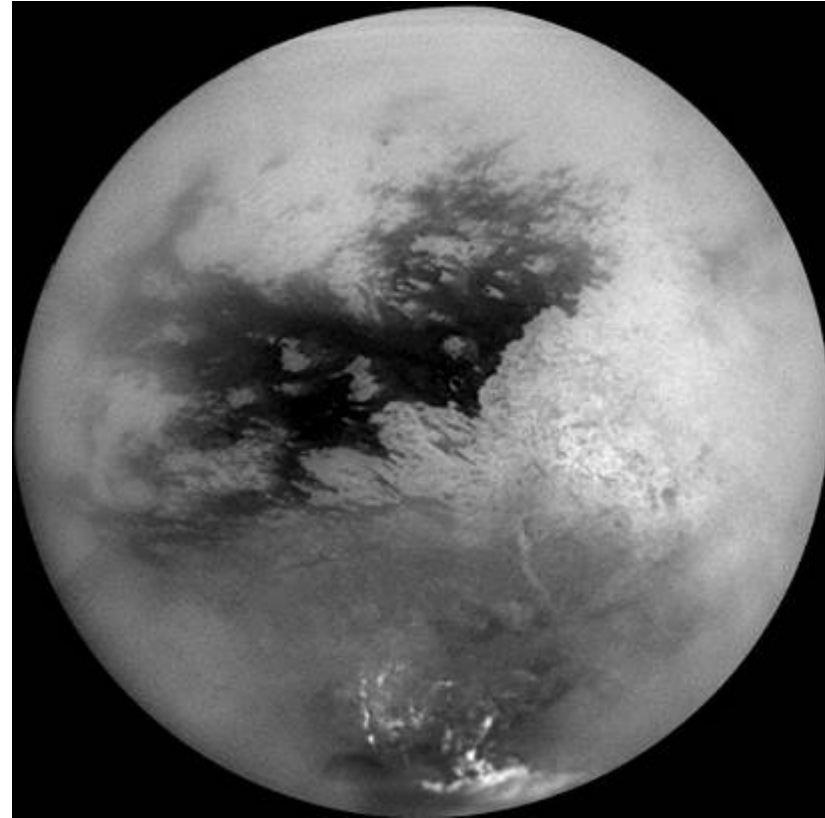
Computer generated image of the topography onshore (brown) and offshore (blue) Central California. Prepared by Norman M. Maher and Robert K. Hall.

SUBAQUEOUS DEPOSITIONAL AND EROSIONAL CYCLIC STEPS

Image courtesy W. Normark



EARTH → TITAN

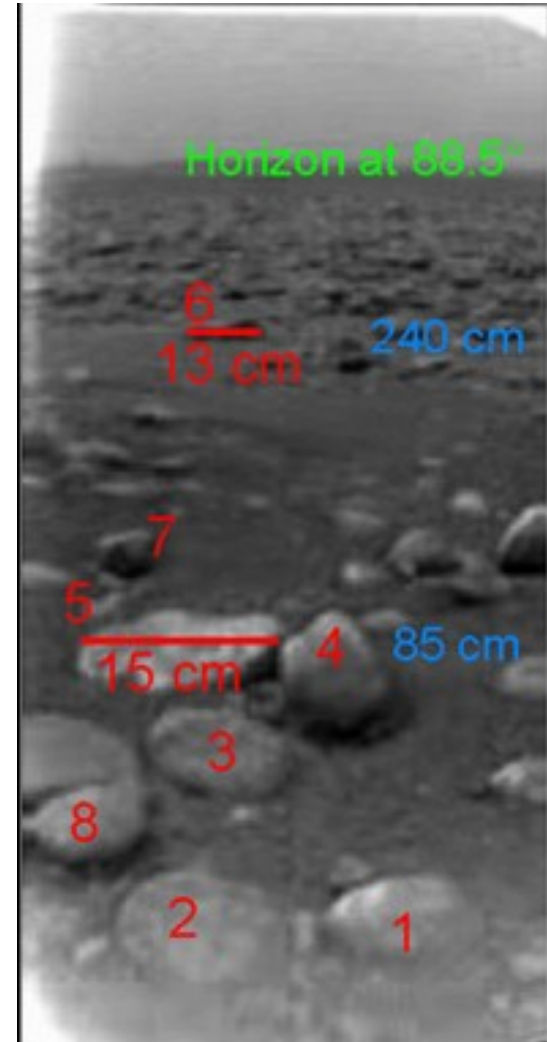
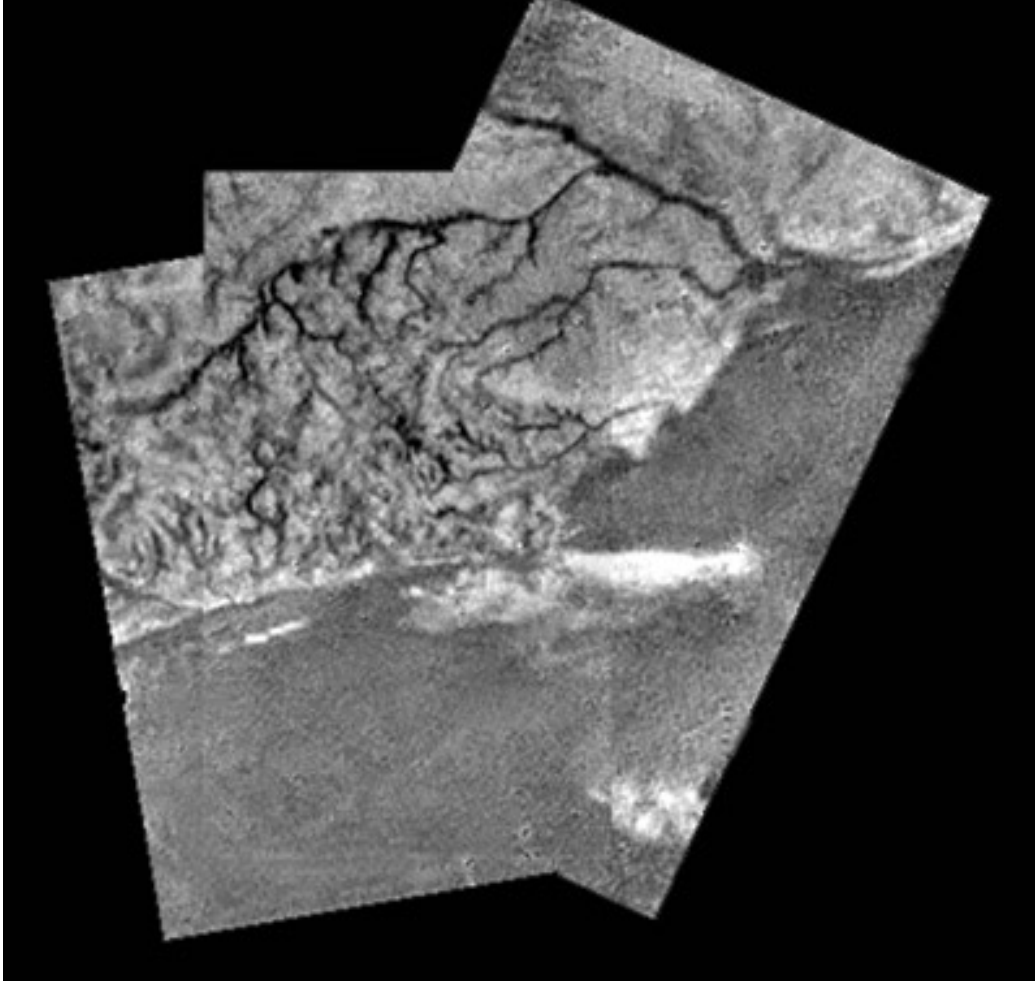


European Space Agency

Water → liquid methane
Granitic rock → ice as a “rock”

ALLUVIAL GRAVEL-BED RIVERS ON TITAN?

Images courtesy European Space Agency



The evidence suggests that at least near where Huygens touched down, there is a **plethora of alluvium in the gravel and sand sizes**. The **gravel** presumably consists of **water ice** and appears to be **fluvially rounded**.⁹⁵

AN EXCITING FUTURE WAITS

